

Quantum calibration

"New Frontiers in Physics", Istituto Lombardo (November 27 2003)

Giacomo Mauro D'Ariano

QUIT Group: <http://www.qubit.it>

Istituto Nazionale di Fisica della Materia

Università di Pavia

Center for Photonic Communication and Computing

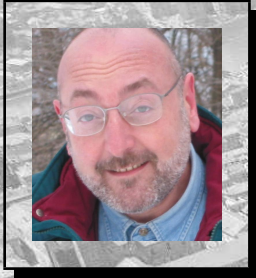
Northwestern University, Evanston IL 60208

Founded by: EC (ATESIT), INFN (PRA-CLON), MIUR (Cofin), US (MURI)

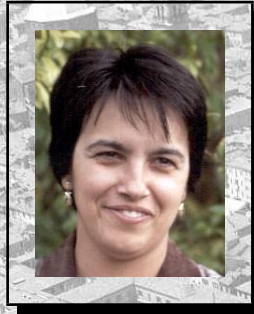
QUIT (Quantum Information Theory)



QUIT (Quantum Information Theory)



G. M. D'Ariano (full prof.)



C. Macchiavello (researcher)



M. F. Sacchi (postdoc)



P. Perinotti (postdoc)



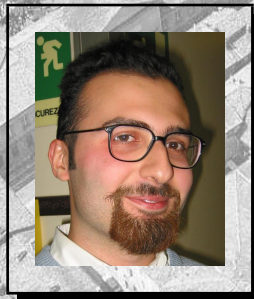
O. Rudolph (ATESIT)



G. Chiribella (PhD student)



P. Lo Presti (PhD student)



F. Buscemi (PhD student)



M. Medici (secretary)



http://www.qubit.it

QUit

quantum information theory group

Sunday, 23 Nov 2003

- people
- faq
- research
- educational
- for visitors
- sponsors



upcoming events

The fun with **Quantum Information** is that you can study the foundations of the enigmatic world of Quantum Mechanics, and, at the same time, you make something useful for practical applications

[Workshop on: Quantum Information Processing for Quantum Communications > November 2003 - July 2004](#)

Quantum Information

Information

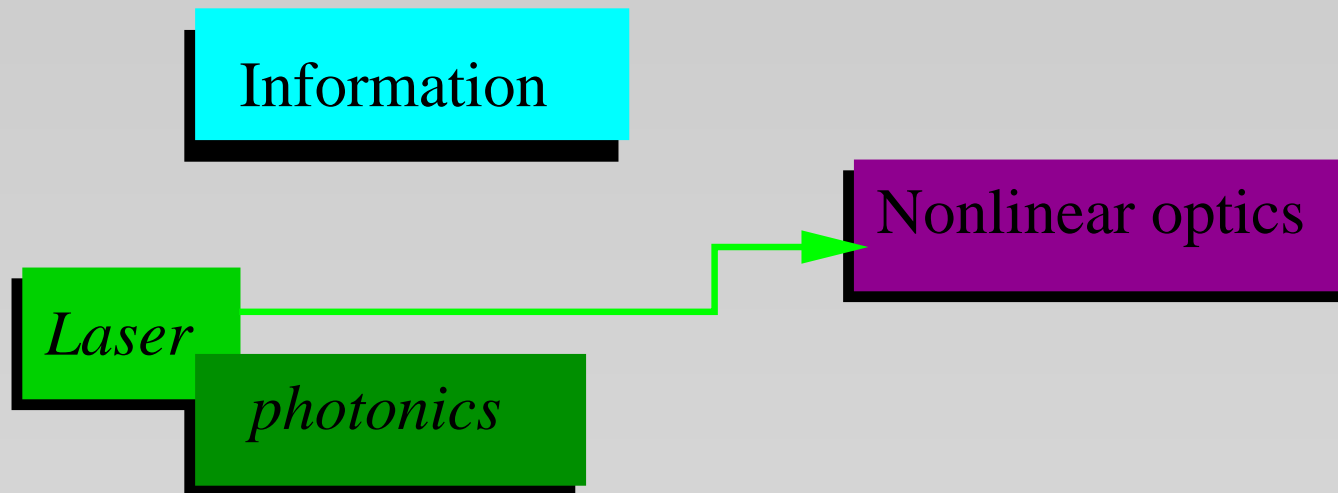
Quantum Information

Information

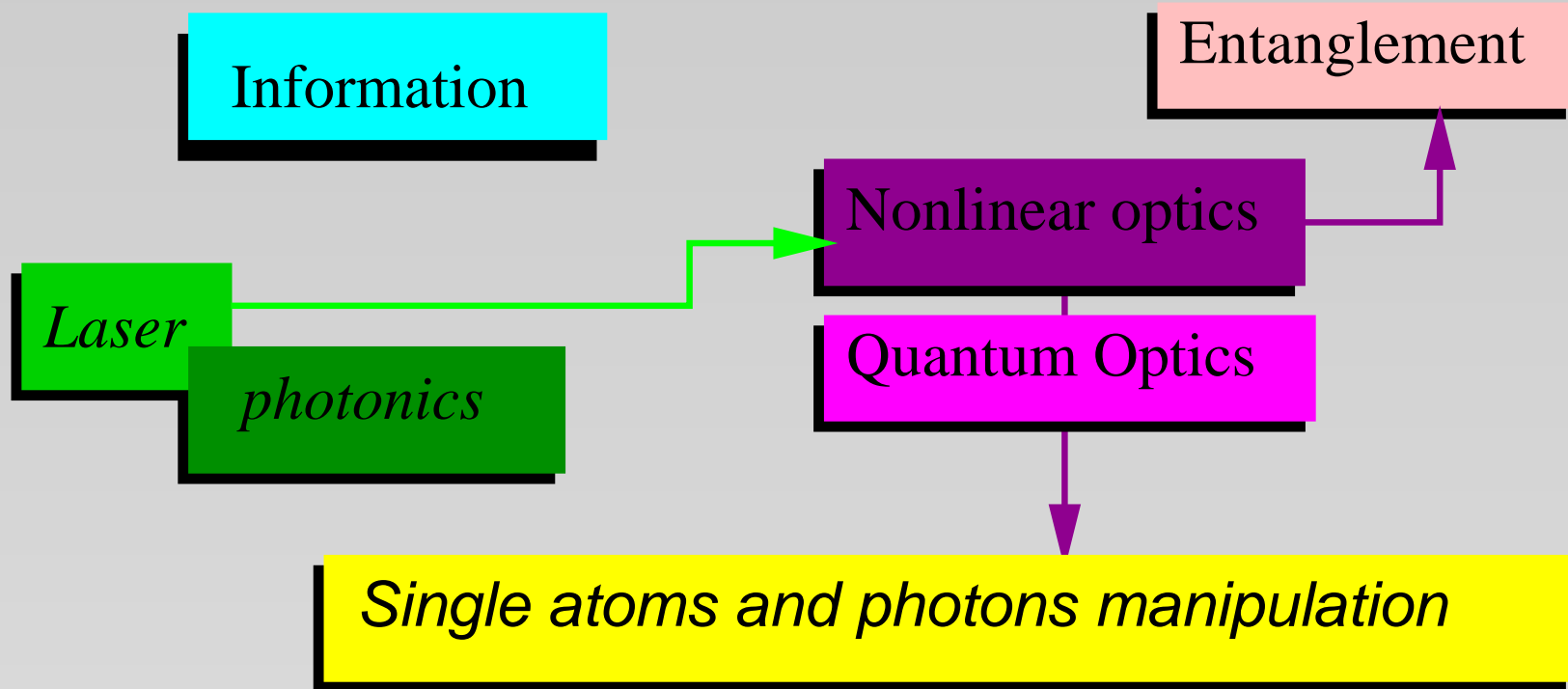
Laser

photonics

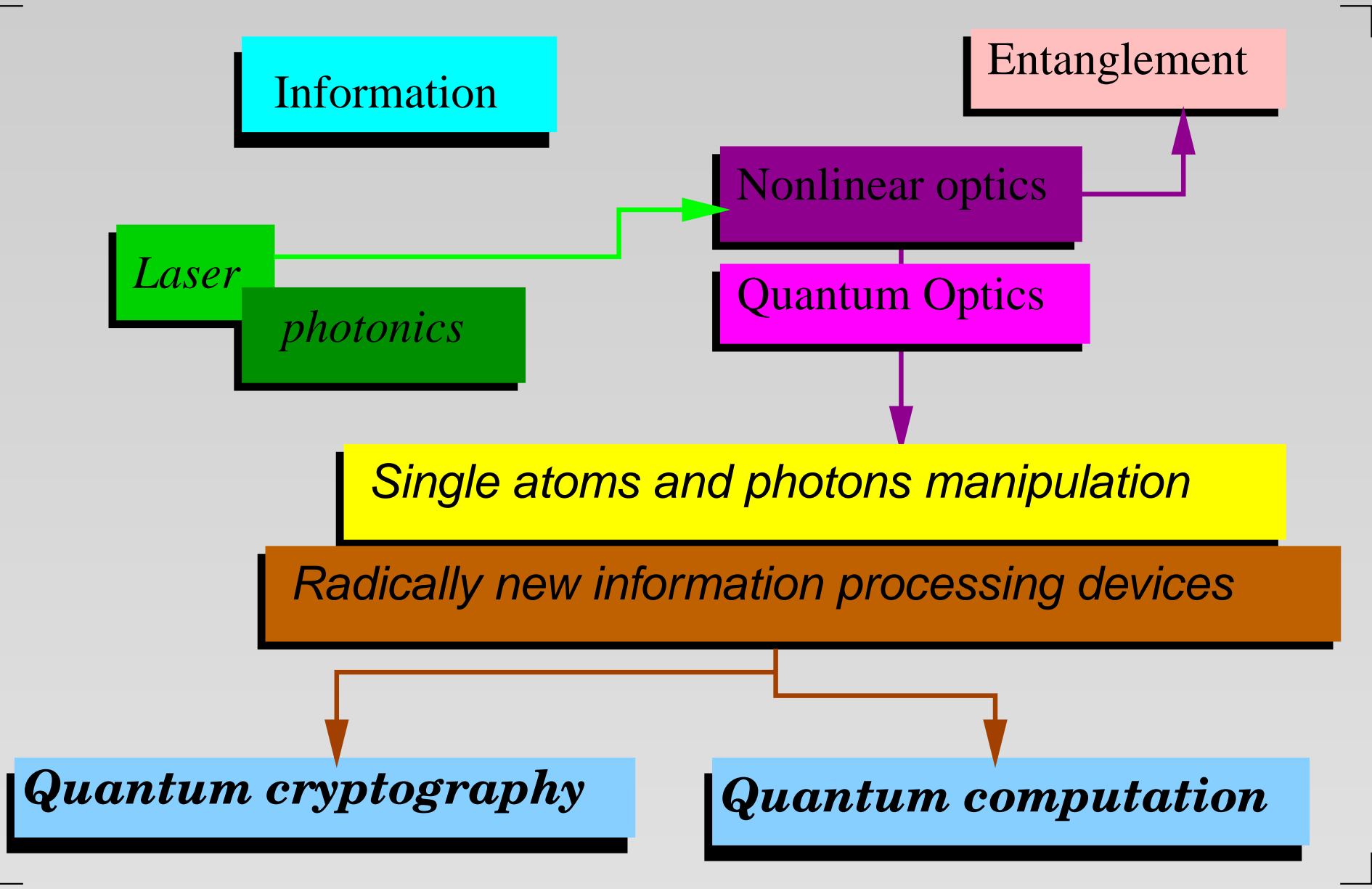
Quantum Information



Quantum Information



Quantum Information



Quantum Information

Applications

Quantum Information

Foundations

Quantum Information

Applications

Quantum Information

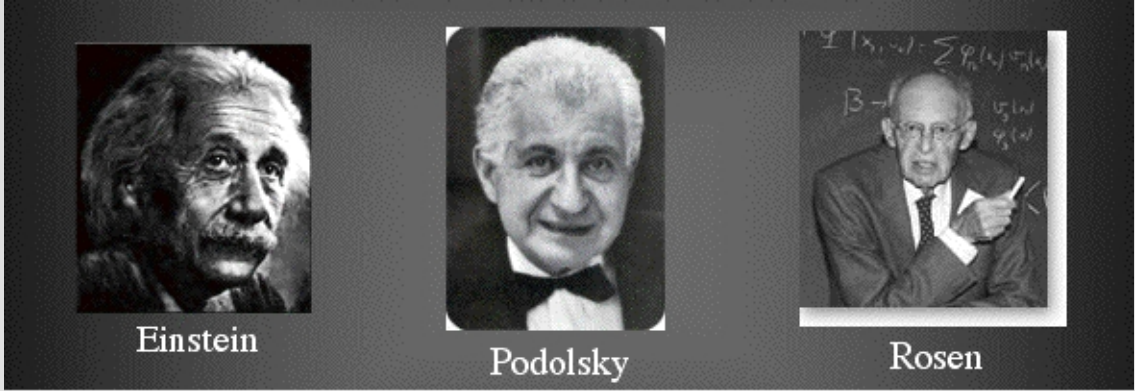
Foundations

Entanglement

Quantum Mechanics: foundations

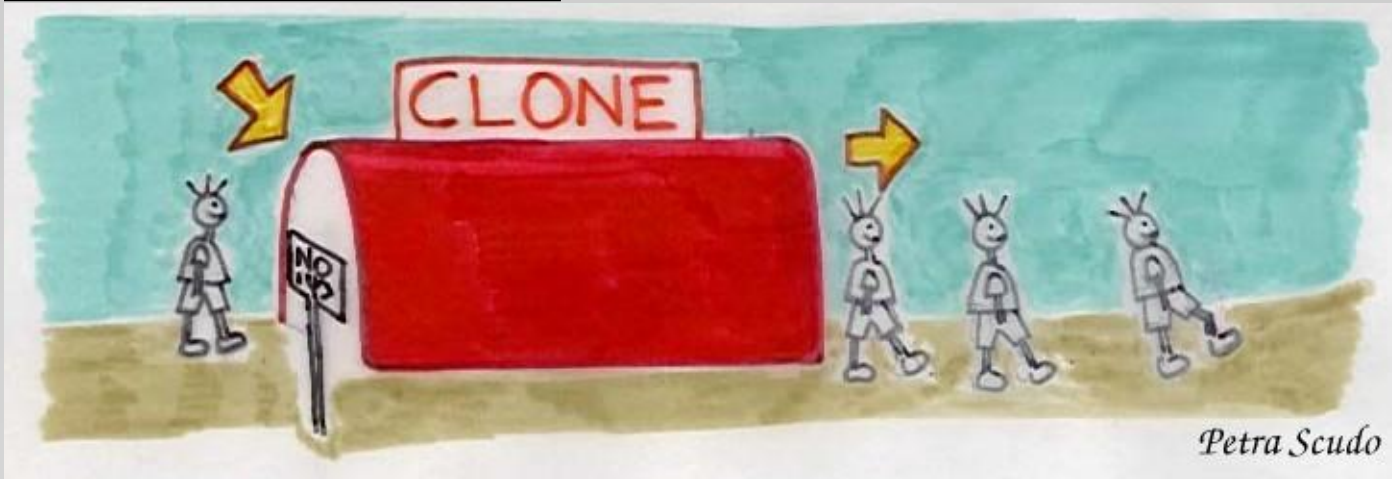


CONTEXTUALITY
NON LOCALITY



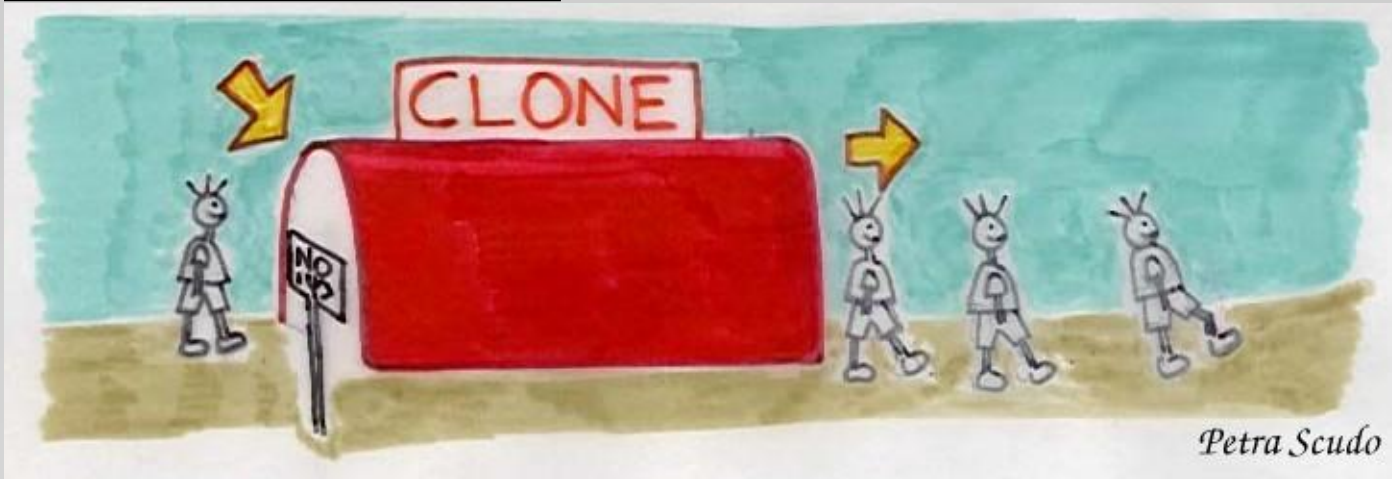
Quantum Mechanics: foundations

- **No cloning theorem** [Wootters and Zurek, Nature **299**, 802 (1982)].



Quantum Mechanics: foundations

- **No cloning theorem** [Wootters and Zurek, Nature **299**, 802 (1982)].



D'Ariano and Yuen, *On the Impossibility of Measuring the Wave Function of a Single Quantum System*, Phys. Rev. Lett. **76** 2832 (1996)

Quantum Mechanics: foundations

- **No cloning theorem** [Wootters and Zurek, *Nature* **299**, 802 (1982)].



D'Ariano and Yuen, *On the Impossibility of Measuring the Wave Function of a Single Quantum System*, *Phys. Rev. Lett.* **76** 2832 (1996)

Bruß, D'Ariano, Macchiavello, and Sacchi, *Approximate quantum cloning and the impossibility of superluminal information transfer*, *Phys. Rev. A* **62** 62302 (2000)

N. Herbert, *FLASH A Superluminal Communicator Based Upon a New Kind of Quantum Measurement*, *Found. Phys.* **12** 1171 (1982); O. H. Heberhard, *Bell's Theorem and the Different Concepts of Locality*, *Nuovo Cimento* **46B** 392 (1978)

Quantum Mechanics: foundations



Bennett, Brassard, Crepeau, Jozsa, Peres,
and Wootters,

*Teleporting an Unknown Quantum State
via Dual Classical and Einstein-Podolsky-
Rosen Channels,*

Phys. Rev. Lett. **70**, 1895 (1993).

Quantum Mechanics: foundations

- **Quantum Measurements** paradigm shift: from **uncontrollable disturbance of measurement** (Messiah) \Rightarrow **control of coherence and measurement engineering**

Quantum Mechanics: foundations

Quantum Measurements

- paradigm shift: from **uncontrollable disturbance of measurement** (Messiah) \Rightarrow **control of coherence and measurement engineering**

Breaching the **Standard Quantum Limit**



- Yuen, *Contractive States and the Standard Quantum Limit for Monitoring Free-Mass Positions*, Phys. Rev. Lett. **51**, 719 (1983).
- Ozawa, *Measurement Breacking the Standard Quantum Limit for Free-Mass Position*, Phys. Rev. Lett. **51**, 719 (1983).

Quantum Mechanics: foundations

Quantum Measurements

- paradigm shift: from **uncontrollable disturbance of measurement** (Messiah) \Rightarrow **control of coherence and measurement engineering**

Breaching the Standard Quantum Limit



- Yuen, *Contractive States and the Standard Quantum Limit for Monitoring Free-Mass Positions*, Phys. Rev. Lett. **51**, 719 (1983).
- Ozawa, *Measurement Breacking the Standard Quantum Limit for Free-Mass Position*, Phys. Rev. Lett. **51**, 719 (1983).

Heisenberg principle (γ -ray microscope gedanken experiment)

- Ozawa, *Universally valid reformulation of the Heisenberg uncertainty principle on noise and disturbance in measurement*, Phys. Rev. Lett. **67** 042105 (2003)
- D'Ariano, *On the Heisenberg principle, namely on the information-disturbance trade-off in a quantum measurement*, Fortschr. Phys. **51** 318 (2003)

Quantum Mechanics: foundations

Quantum Measurements

- paradigm shift: from **uncontrollable disturbance of measurement** (Messiah) \Rightarrow **control of coherence and measurement engineering**

Breaching the **Standard Quantum Limit**



- Yuen, *Contractive States and the Standard Quantum Limit for Monitoring Free-Mass Positions*, Phys. Rev. Lett. **51**, 719 (1983).
- Ozawa, *Measurement Breacking the Standard Quantum Limit for Free-Mass Position*, Phys. Rev. Lett. **51**, 719 (1983).

Heisenberg principle (γ -ray microscope gedanken experiment)

- Ozawa, *Universally valid reformulation of the Heisenberg uncertainty principle on noise and disturbance in measurement*, Phys. Rev. Lett. **67** 042105 (2003)
- D'Ariano, *On the Heisenberg principle, namely on the information-disturbance trade-off in a quantum measurement*, Fortschr. Phys. **51** 318 (2003)
- Buscemi, D'Ariano and Perinotti, *Non orthogonal perfectly repeatable quantum measurements*, Phys. Rev. Lett. (quant-ph/0310041)

Quantum Information: applications

Applications

Quantum Information

Foundations

Quantum Information: applications

Optimization



Classification

= *Purely mathematical*

Quantum Information: applications

Optimization

Classification

= *Purely mathematical*

Feasibility analysis

= *Computer simulation, theory*

Quantum Information: applications

Optimization

Classification

= *Purely mathematical*

Feasibility analysis

= *Computer simulation, theory*

Experiment

Our general research program

To construct a **radically new generation of quantum devices** for quantum information technology

Our general research program

To construct a **radically new generation of quantum devices** for quantum information technology

Programmable quantum detectors [*]

Our general research program

To construct a **radically new generation of quantum devices** for quantum information technology

Programmable quantum detectors [*]



Our general research program

To construct a **radically new generation of quantum devices** for quantum information technology

Programmable quantum detectors [*]



Universal quantum detectors



Our general research program

To construct a **radically new generation of quantum devices** for quantum information technology

Programmable quantum detectors [*]



Universal quantum detectors



Quantum tomographers



Our general research program

... and to design devices which **emulate optimally impossible machines**

Our general research program

... and to design devices which **emulate optimally impossible machines**

Cloners



Our general research program

... and to design devices which **emulate optimally impossible machines**

Cloners



Transpositors



Our general research program

... and to design devices which **emulate optimally impossible machines**

Cloners



Transpositors



Superpositors



Our general research program

... and to design devices which **emulate optimally impossible machines**

Cloners



Transpositors



Superpositors

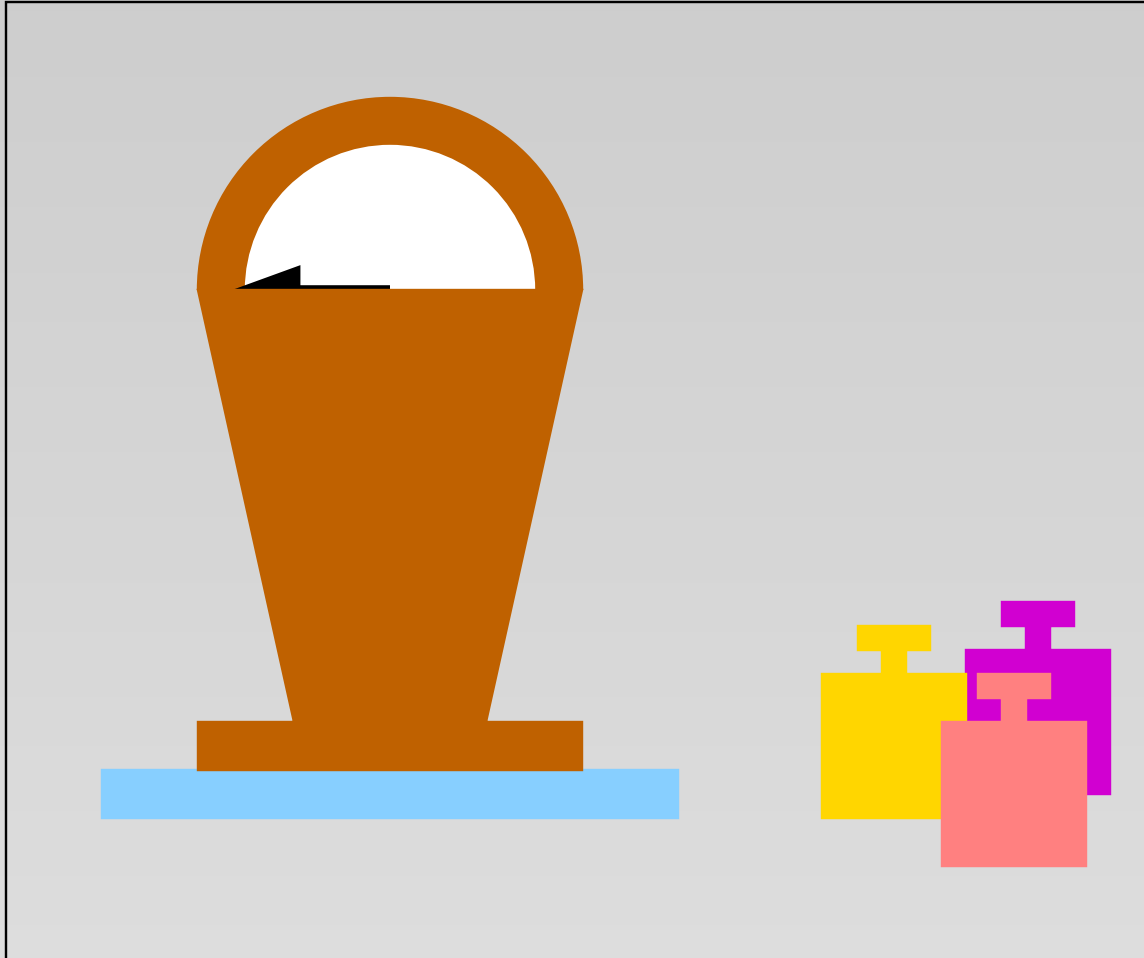


.... and their

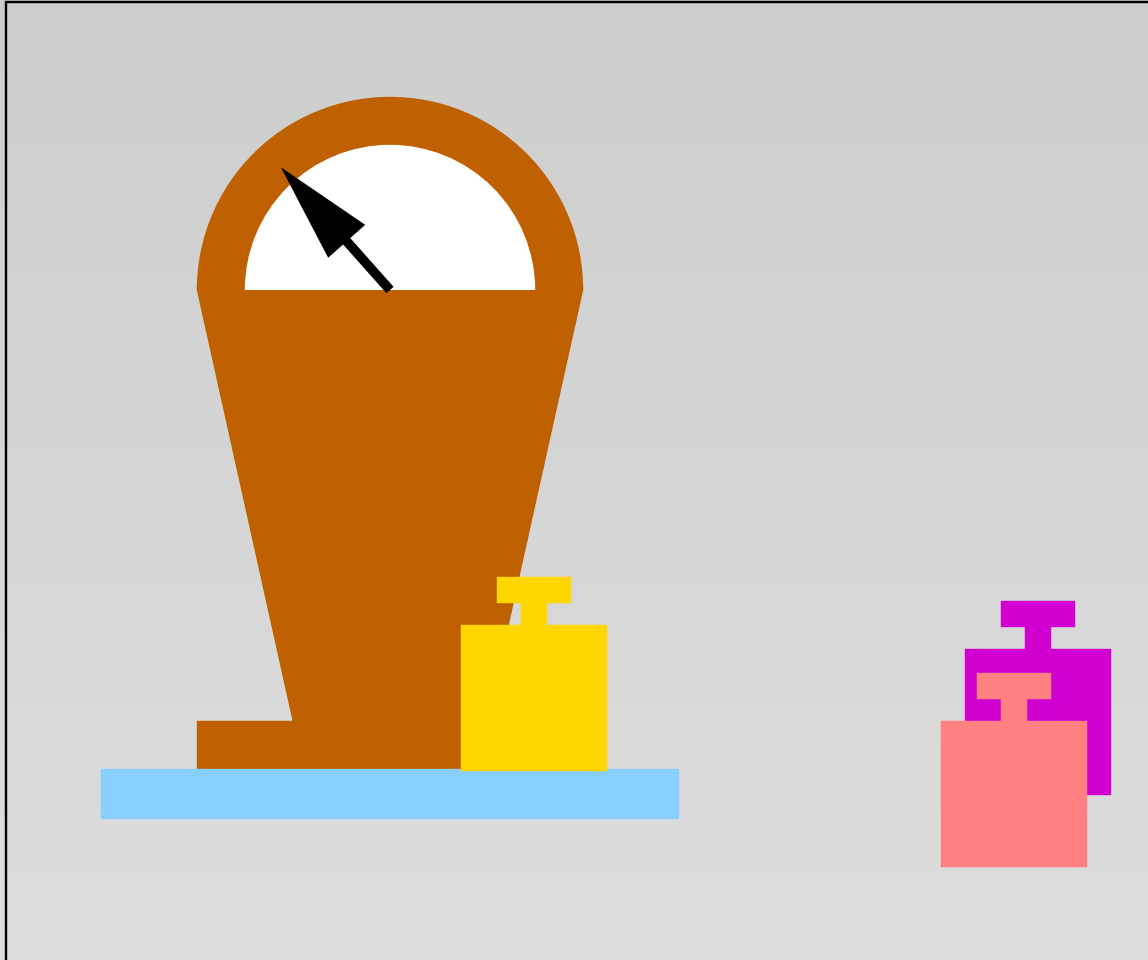
Quantum calibration



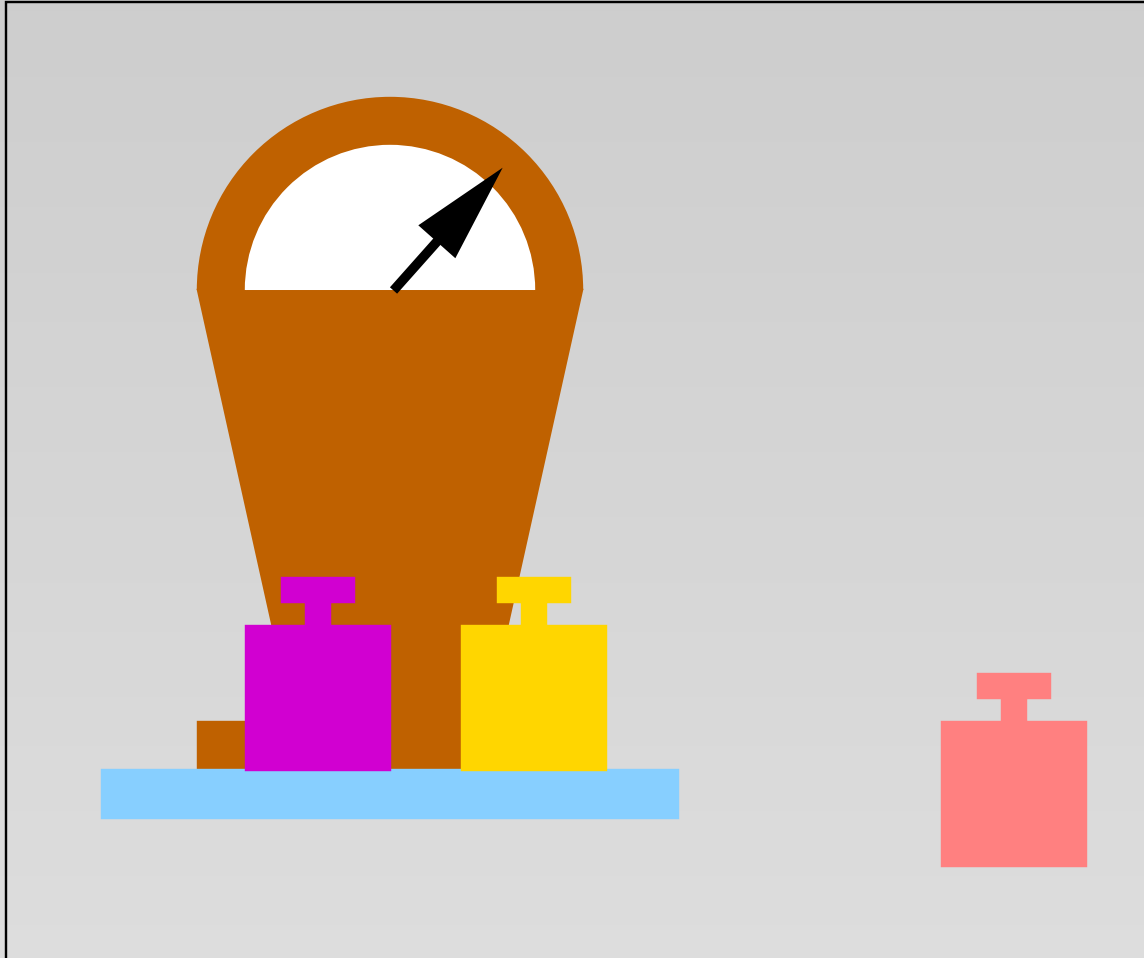
Calibration of a scale



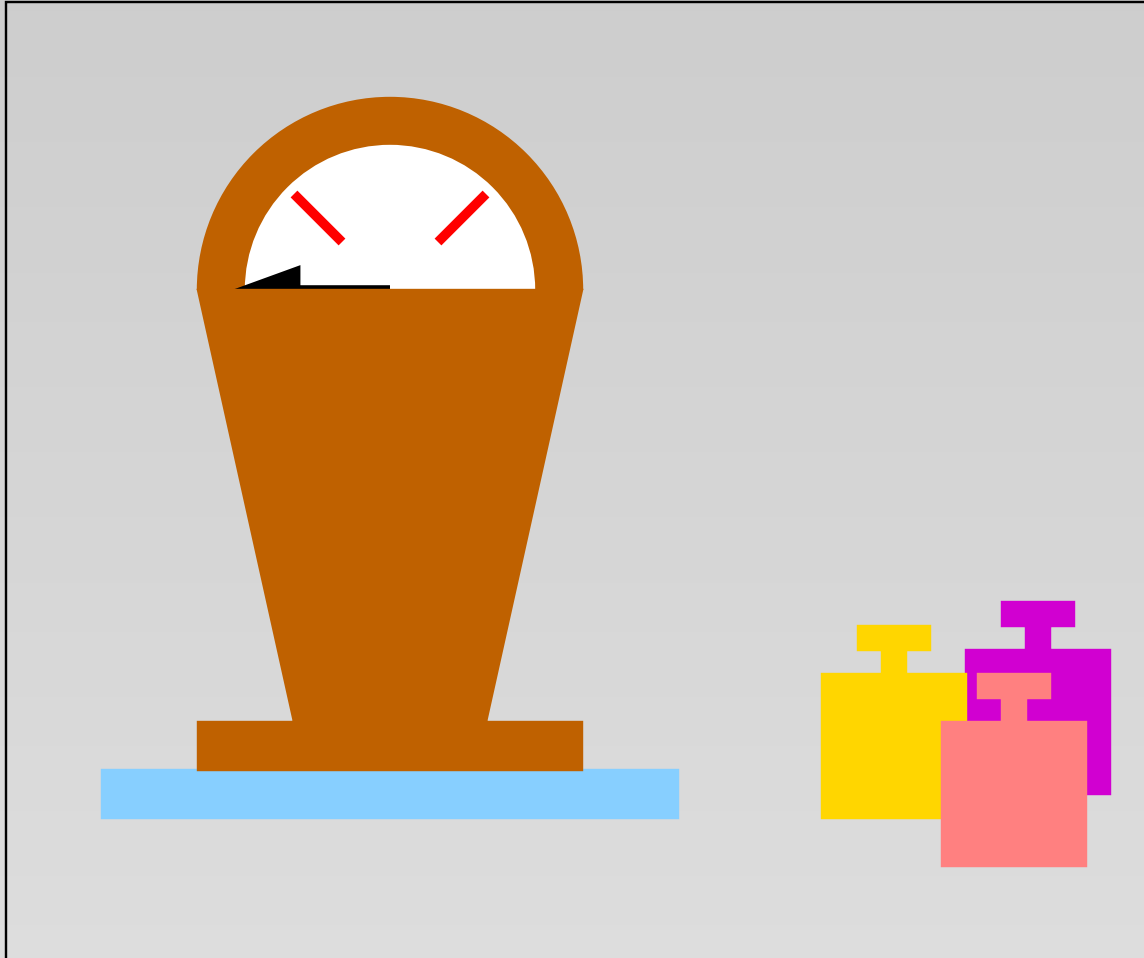
Calibration of a scale



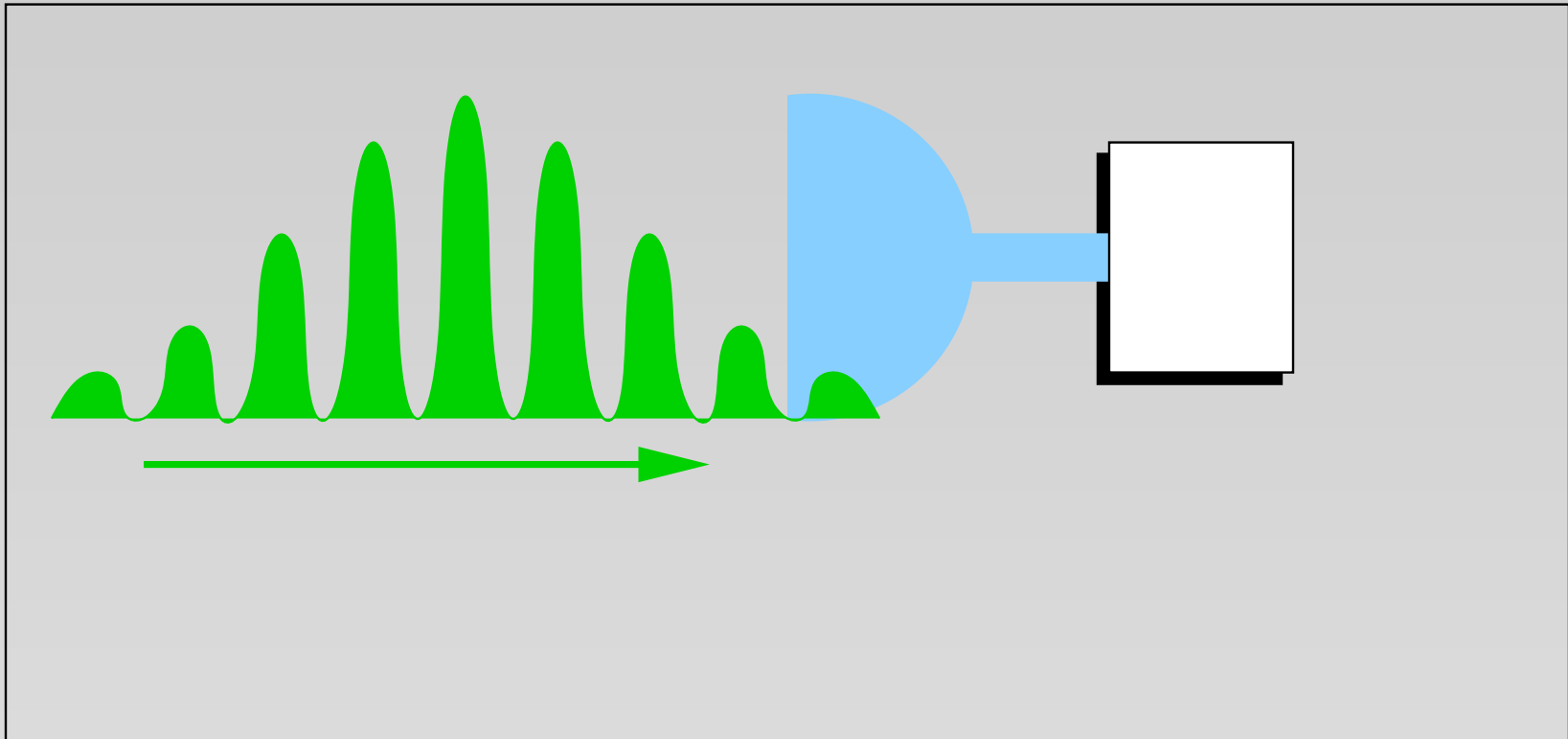
Calibration of a scale



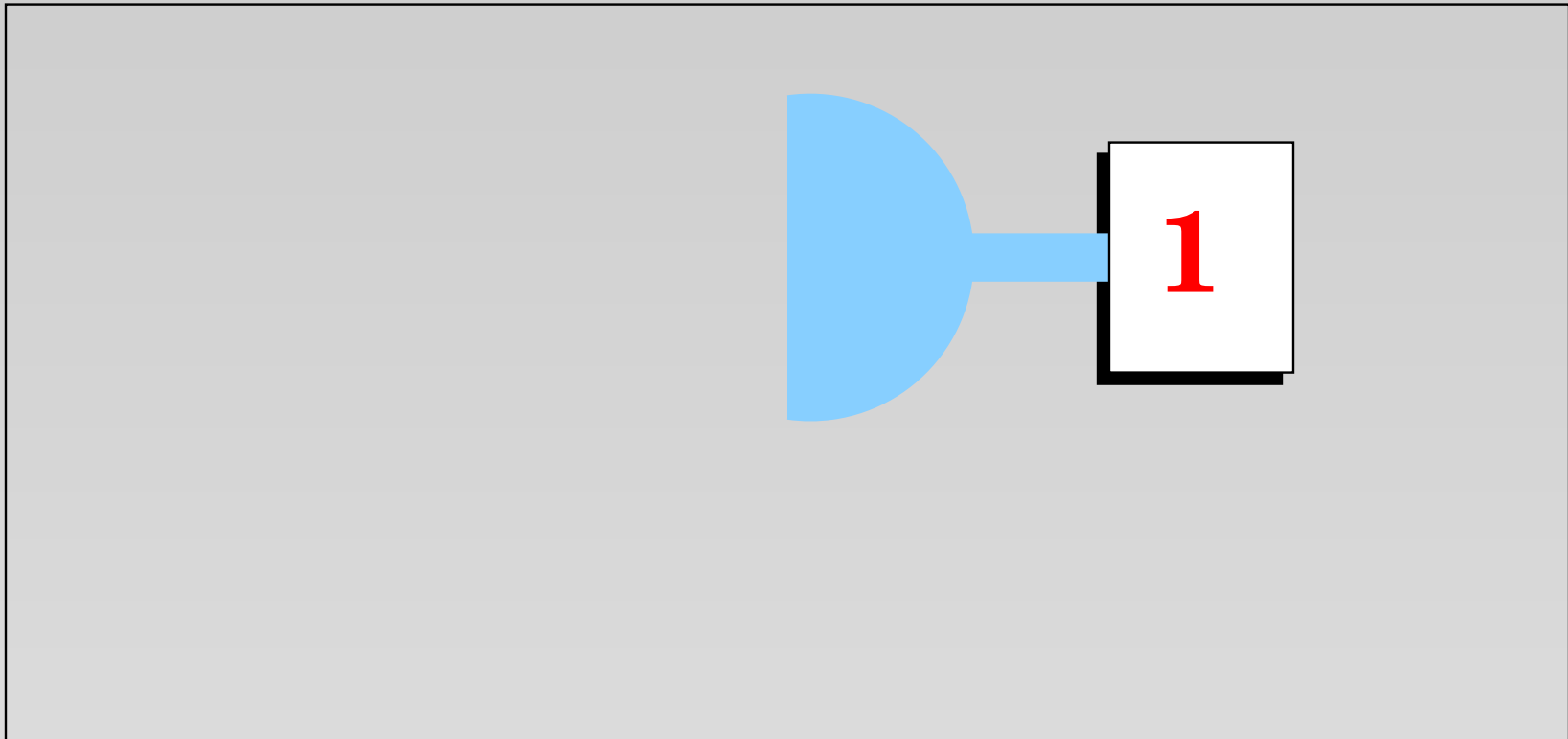
Calibration of a scale



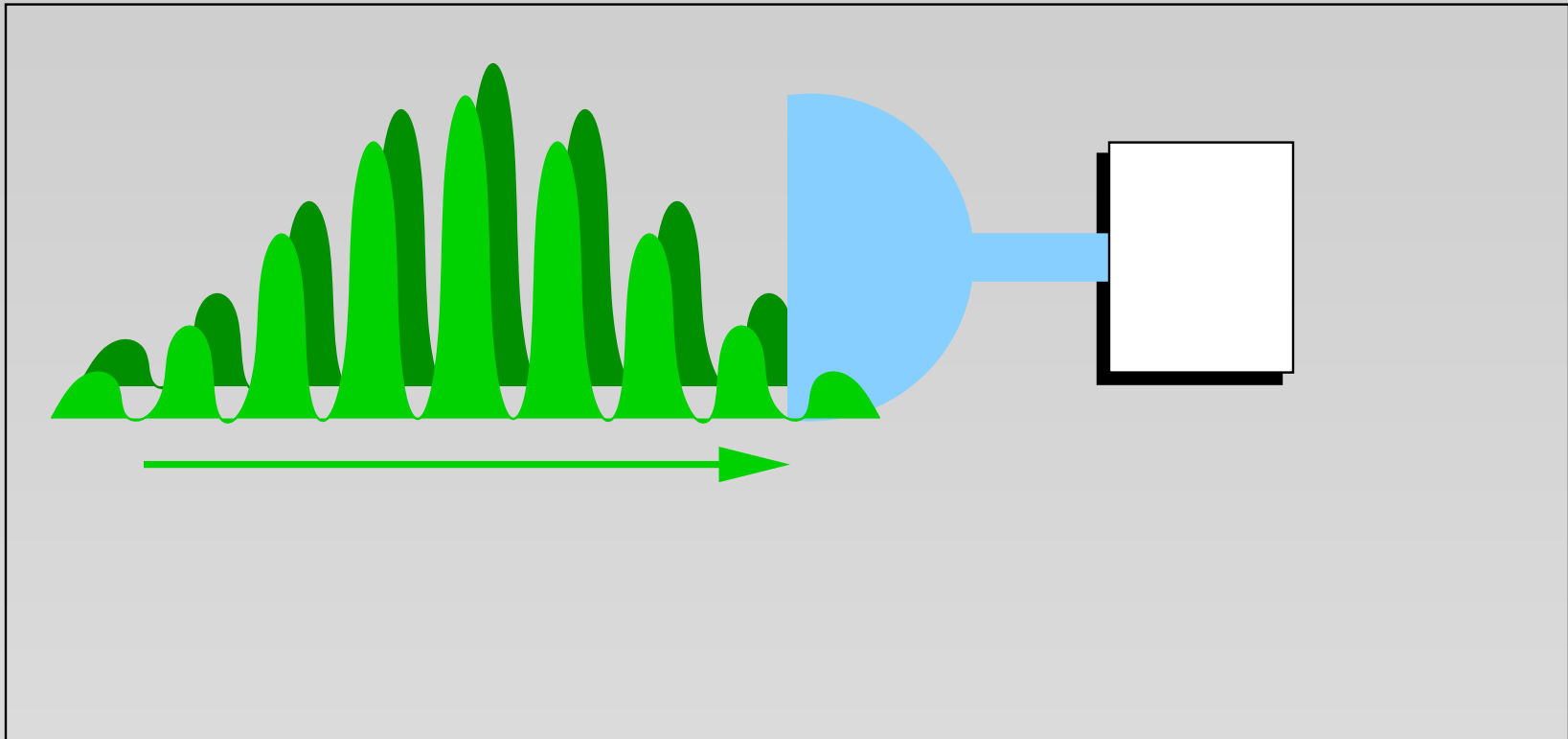
Calibration of a photocounter



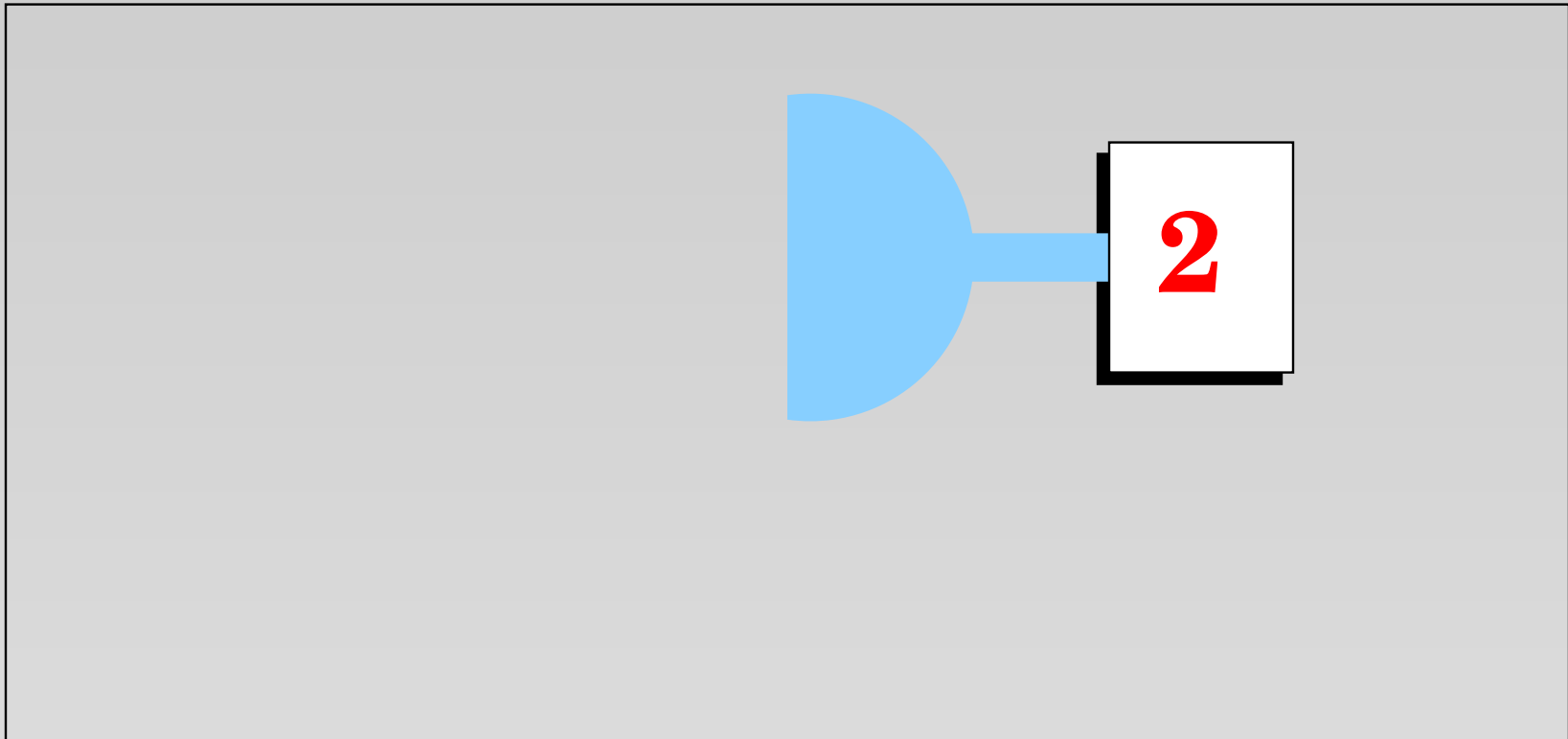
Calibration of a photocounter



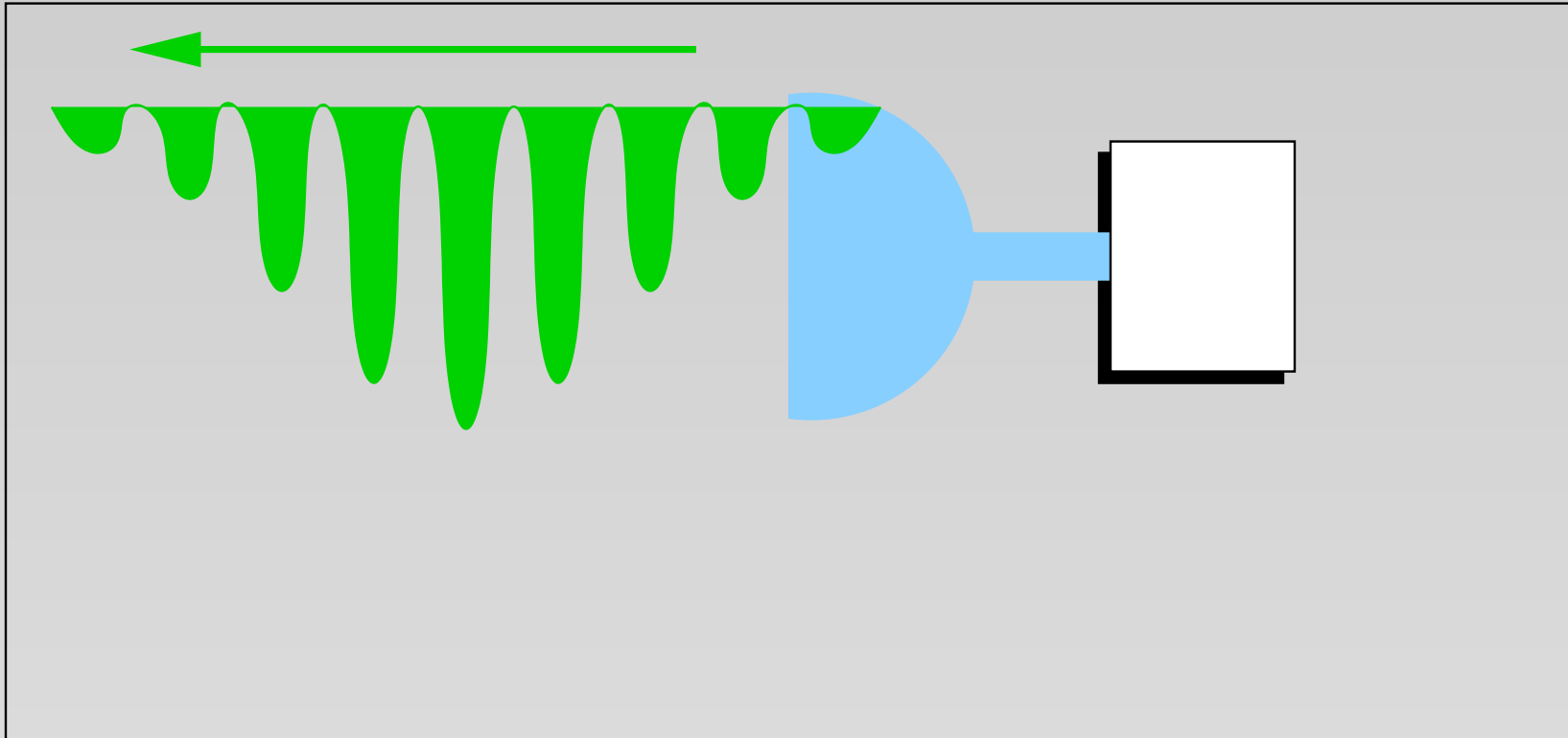
Calibration of a photocounter



Calibration of a photocounter



Calibration of a photocounter



What does an apparatus measure?

What does an apparatus measure?

If we have an apparatus which performs a quantum measurement, how can we know what and how much it measures?

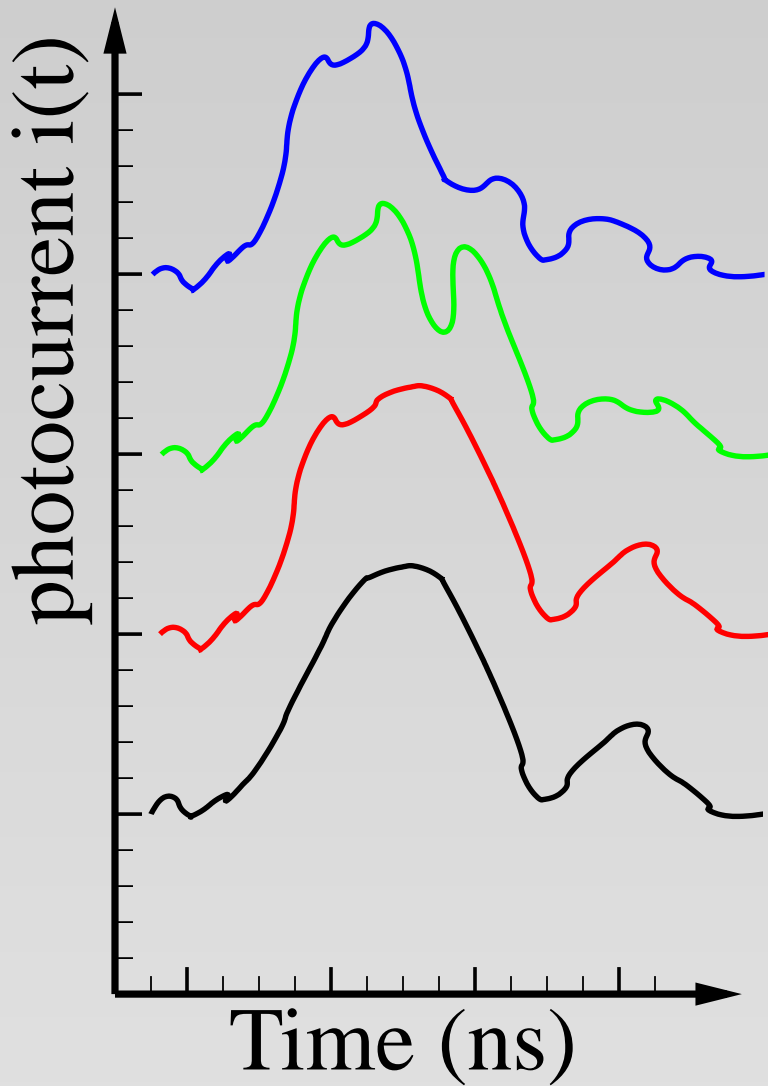
What does an apparatus measure?

If we have an apparatus which performs a quantum measurement, how can we know what and how much it measures?

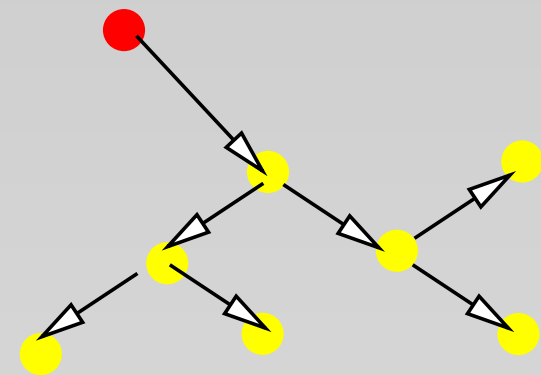
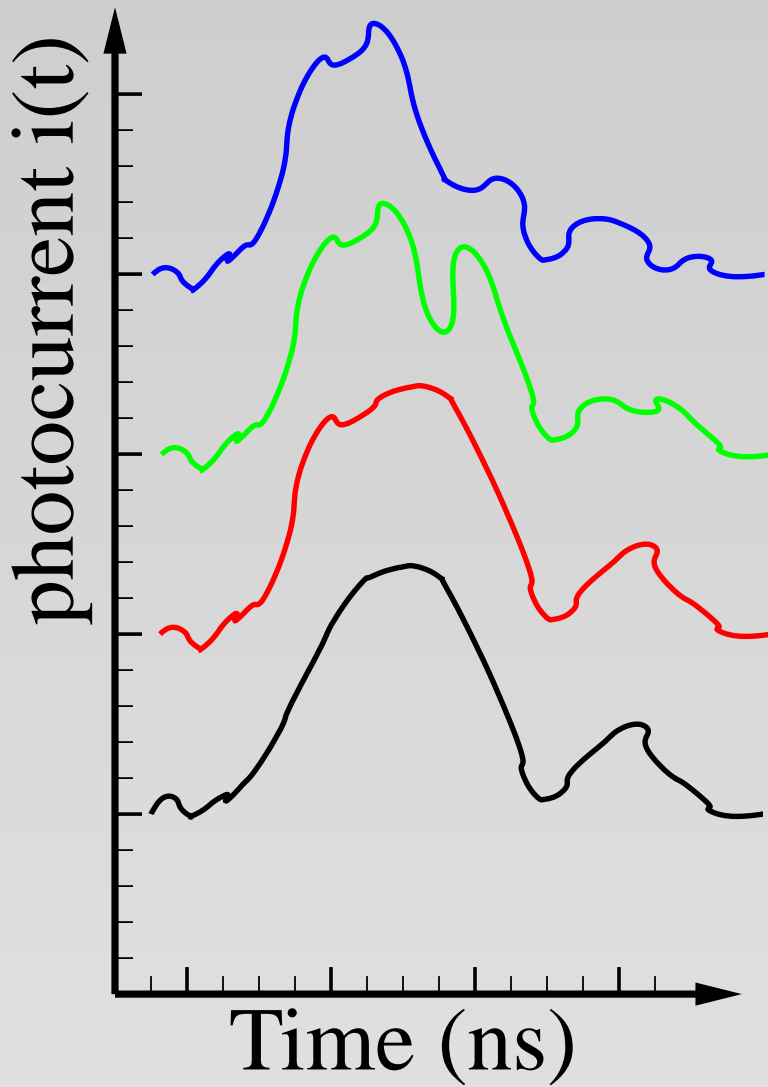
It is the theory which decides what we can observe!

—Einstein to Heisenberg

How many photons are detected?



How many photons are detected?



Problem

Is it possible to calibrate a photo-detector, and more generally any quantum measuring apparatus, without using the theoretical statistical mechanics description of its functioning?

How we describe a measuring apparatus

How we describe a measuring apparatus

A measuring apparatus with possible "outcomes" $\{n = 1, 2, \dots\}$ is described by a set of operators (called POVM)

$$\mathbf{P} = \{P_n\},$$

How we describe a measuring apparatus

A measuring apparatus with possible "outcomes" $\{n = 1, 2, \dots\}$ is described by a set of operators (called POVM)

$$\mathbf{P} = \{P_n\},$$

which provide the probability $p(n)$ of each n for all possible states ρ via

$$p(n) = \text{Tr}[P_n \rho] \quad \text{Born rule}$$

How we describe a measuring apparatus

A measuring apparatus with possible "outcomes" $\{n = 1, 2, \dots\}$ is described by a set of operators (called POVM)

$$\mathbf{P} = \{P_n\},$$

which provide the probability $p(n)$ of each n for all possible states ρ via

$$p(n) = \text{Tr}[P_n \rho] \quad \text{Born rule}$$

In order to have $p(n)$ a probability the operators P_n must satisfy the constraints

$$P_n \geq 0, \quad \sum_n P_n = I.$$

How can we calibrate a measuring apparatus

How can we calibrate a measuring apparatus

In principle, we can calibrate a quantum measuring apparatus without knowing its functioning by determining experimentally its POVM $\{P_n\}$.

Quantum tomography

- Quantum tomography is a method to estimate the ensemble average $\langle H \rangle$ of any arbitrary operator H by measuring a **quorum of observables** $\{O_l\}$.

Quantum tomography

- Quantum tomography is a method to estimate the ensemble average $\langle H \rangle$ of any arbitrary operator H by measuring a **quorum of observables** $\{O_l\}$.
- The operator H is expanded on the quorum:

$$H = \sum_l c_l(H) O_l.$$

Quantum tomography

- Quantum tomography is a method to estimate the ensemble average $\langle H \rangle$ of any arbitrary operator H by measuring a **quorum of observables** $\{O_l\}$.
- The operator H is expanded on the quorum:

$$H = \sum_l c_l(H) O_l.$$

- The tomographic estimation of the ensemble average $\langle H \rangle$ is obtained by averaging over both the ensemble and the quorum.

Quantum tomography

- Quantum tomography is a method to estimate the ensemble average $\langle H \rangle$ of any arbitrary operator H by measuring a **quorum of observables** $\{O_l\}$.
- The operator H is expanded on the quorum:

$$H = \sum_l c_l(H) O_l.$$

- The tomographic estimation of the ensemble average $\langle H \rangle$ is obtained by averaging over both the ensemble and the quorum.
- The estimation of the density matrix element ρ_{ij} corresponds to $H = |i\rangle\langle j|$.

Quantum tomography

- Quantum tomography is a method to estimate the ensemble average $\langle H \rangle$ of any arbitrary operator H by measuring a **quorum of observables** $\{O_l\}$.
- The operator H is expanded on the quorum:

$$H = \sum_l c_l(H) O_l.$$

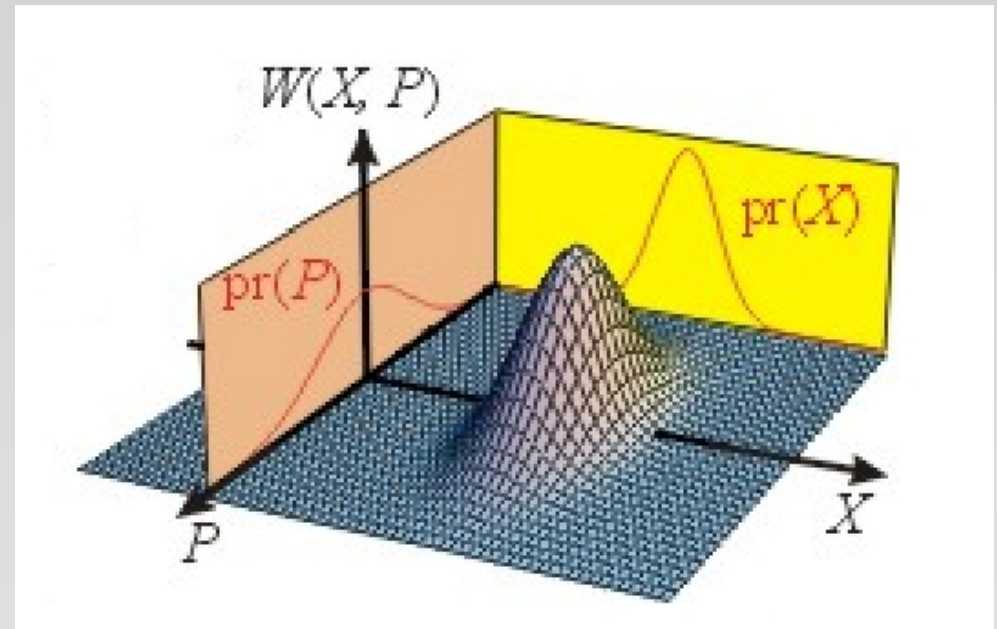
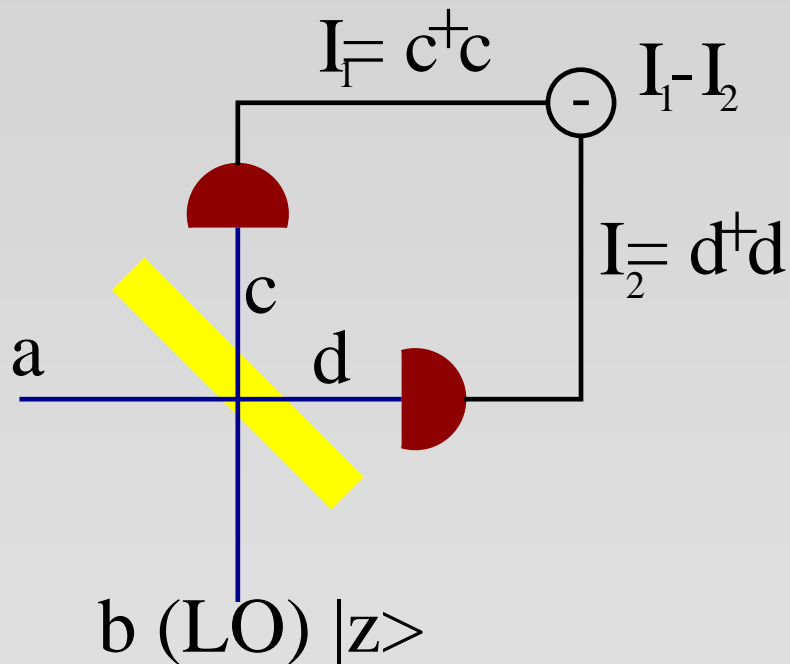
- The tomographic estimation of the ensemble average $\langle H \rangle$ is obtained by averaging over both the ensemble and the quorum.
- The estimation of the density matrix element ρ_{ij} corresponds to $H = |i\rangle\langle j|$.
- There are general method for unbiasing instrumental noise, adaptive techniques, maximum-likelihood strategies, etc.

Homodyne tomography

- In quantum optics for each field mode a **quorum** \equiv {quadratures}

$$X_\phi = \frac{1}{2} (a^\dagger e^{i\phi} + a e^{-i\phi}) \equiv Q \cos \phi + P \sin \phi.$$

$$\langle H \rangle = \int_0^\pi \frac{d\phi}{\pi} \langle E_H(X_\phi; \phi) \rangle, \quad E_H(x; \phi) = \frac{1}{4} \int_{-\infty}^{+\infty} dk |k| \text{Tr}[H e^{ikX_\phi}] e^{-ikx}.$$



Pauli tomography

Pauli matrices I , σ_x , σ_y , σ_z orthonormal basis for the qubit operator space:

$$H = \frac{1}{2}[\boldsymbol{\sigma} \cdot \text{Tr}(\boldsymbol{\sigma} H) + I \text{Tr}(H)] .$$

Pauli tomography

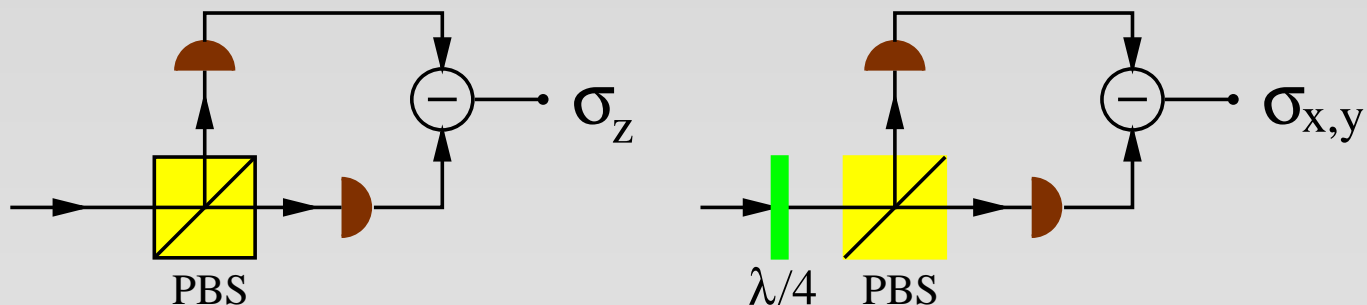
Pauli matrices $I, \sigma_x, \sigma_y, \sigma_z$ orthonormal basis for the qubit operator space:

$$H = \frac{1}{2}[\boldsymbol{\sigma} \cdot \text{Tr}(\boldsymbol{\sigma} H) + I \text{Tr}(H)] .$$

- In Quantum Optics the qubits are encoded on polarization of single photons:

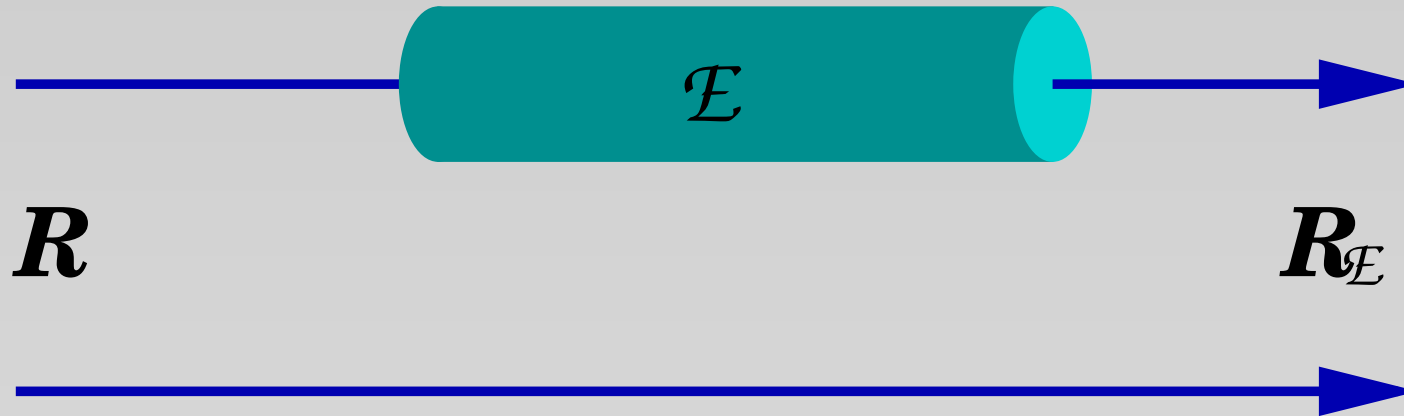
$$\sigma_z = h^\dagger h - v^\dagger v,$$

$$|\uparrow\rangle \equiv |1\rangle_h |0\rangle_v, \quad |\downarrow\rangle \equiv |0\rangle_h |1\rangle_v,$$



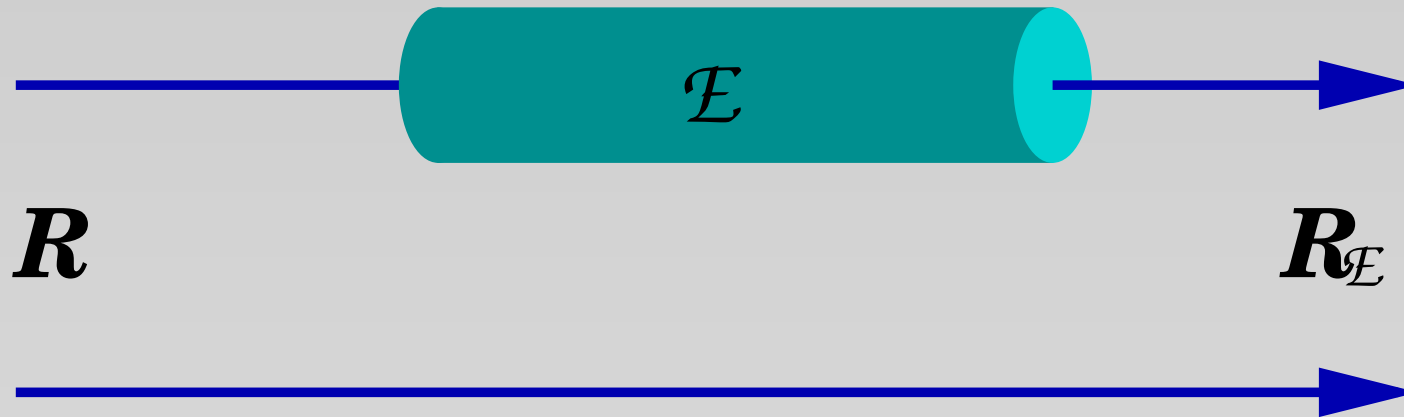
Faithful states

- A bipartite state R is **faithful** when acting with a device on R as in figure the output $R_{\mathcal{E}}$ carries a complete information about the operation \mathcal{E} of the device



Faithful states

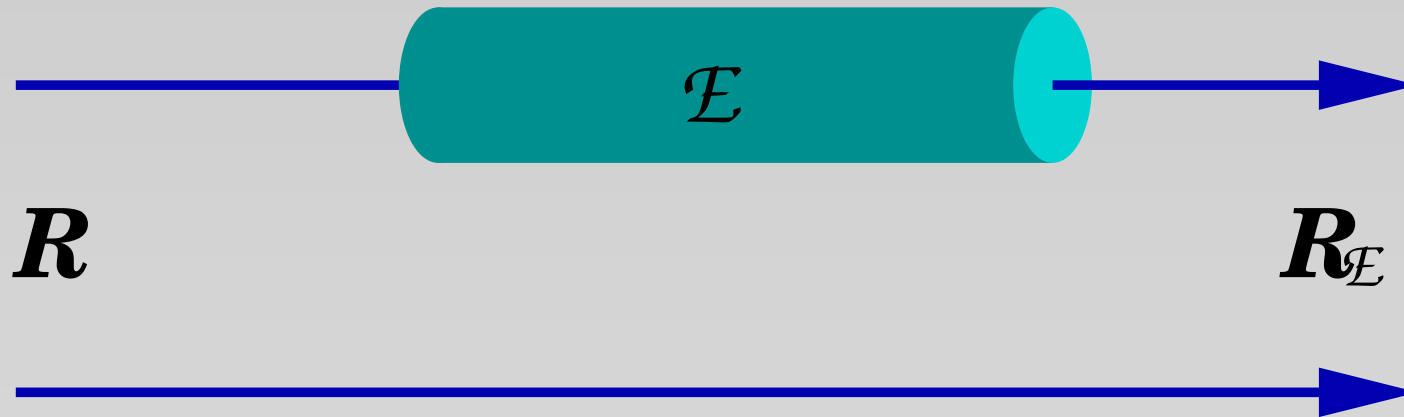
- A bipartite state R is **faithful** when acting with a device on R as in figure the output $R_{\mathcal{E}}$ carries a complete information about the operation \mathcal{E} of the device



correspondence $R_{\mathcal{E}} \Leftrightarrow \mathcal{E}$ one-to-one

Faithful states

- A bipartite state R is **faithful** when acting with a device on R as in figure the output $R_{\mathcal{E}}$ carries a complete information about the operation \mathcal{E} of the device

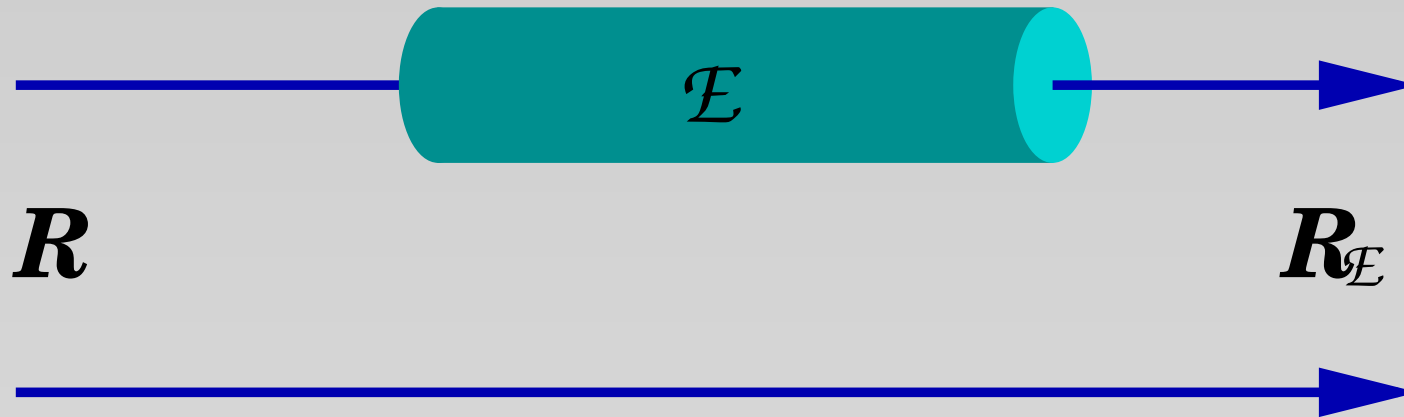


correspondence $R_{\mathcal{E}} \Leftrightarrow \mathcal{E}$ one-to-one

- R is faithful when its **associated map** $\mathcal{R}(\rho) = \text{Tr}_1[(\rho^T \otimes I)R]$ is invertible.

Faithful states

- A bipartite state R is **faithful** when acting with a device on R as in figure the output $R_{\mathcal{E}}$ carries a complete information about the operation \mathcal{E} of the device



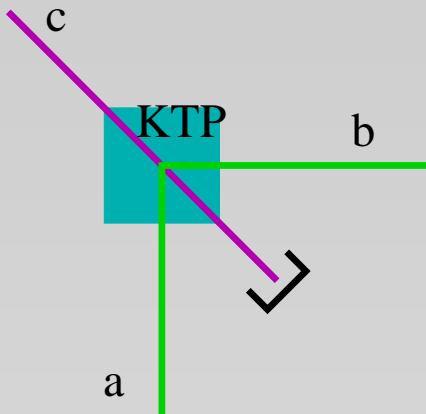
correspondence $R_{\mathcal{E}} \Leftrightarrow \mathcal{E}$ one-to-one

- R is faithful when its **associated map** $\mathcal{R}(\rho) = \text{Tr}_1[(\rho^T \otimes I)R]$ is invertible.

There are **pure faithful states: the entangled states** (quantum parallelism of entanglement!)

Entangled states in quantum optics

- Nonlinear Quantum Optics: **parametric downconversion of vacuum**

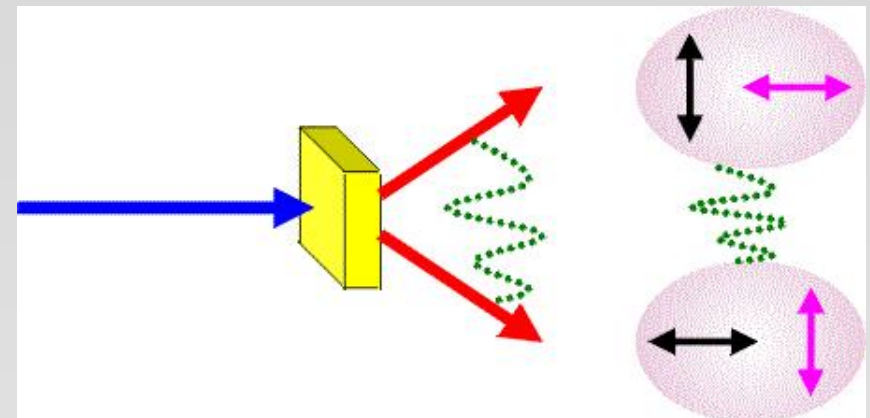


Hamiltonian $H \propto ca^\dagger b^\dagger + h.c.$
 where $\omega_c = \omega_a + \omega_b$

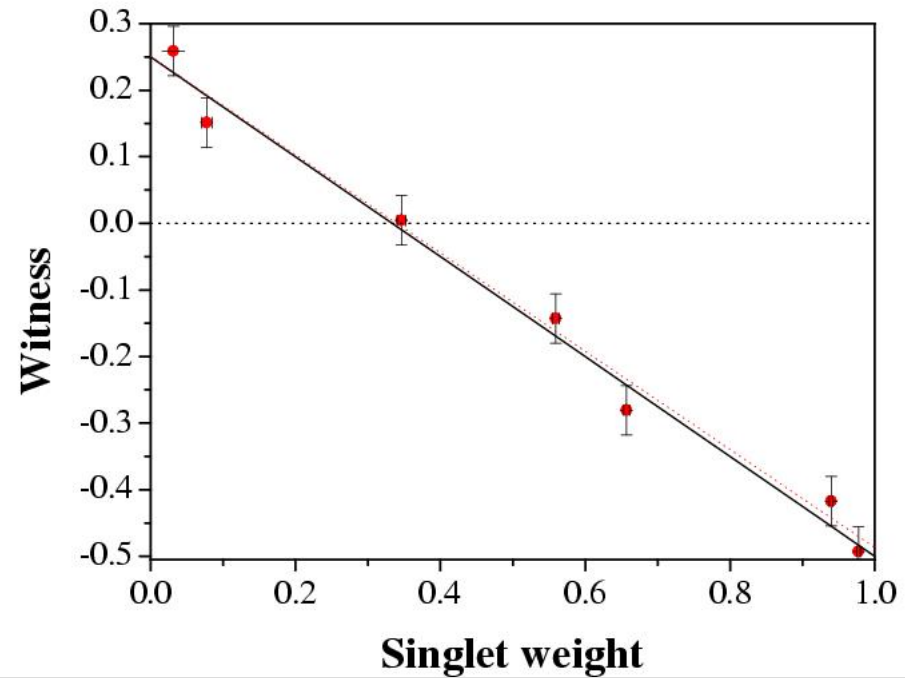
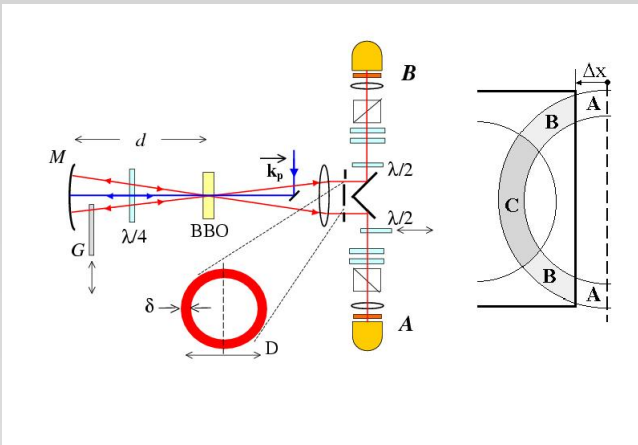
- From input vacuum in a and b and classical pump c produces the

twin-beam

$$|\Psi\rangle\rangle = (1 - |\xi|^2)^{\frac{1}{2}} \sum_{n=0}^{\infty} \xi^n |n\rangle \otimes |n\rangle$$



Production of faithful states

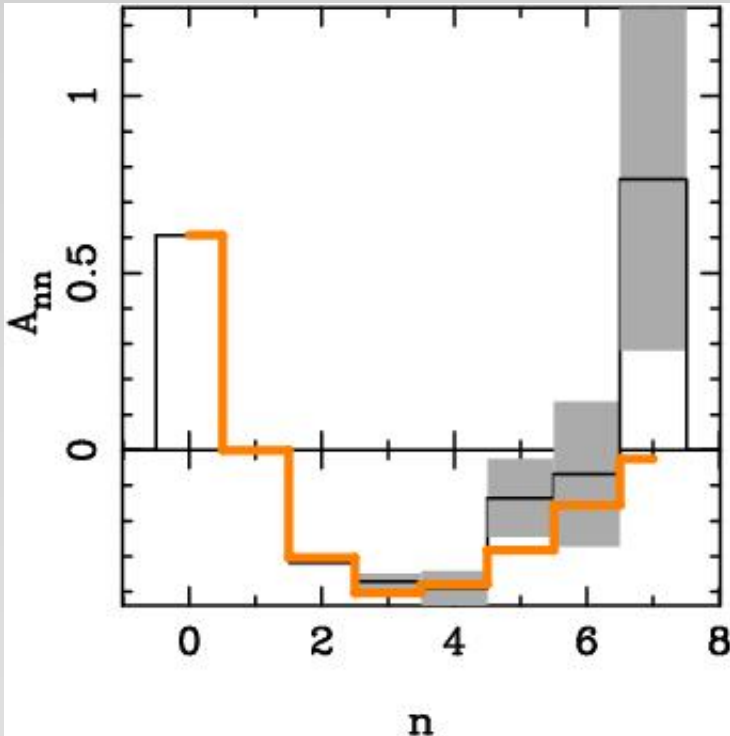


Faithful states

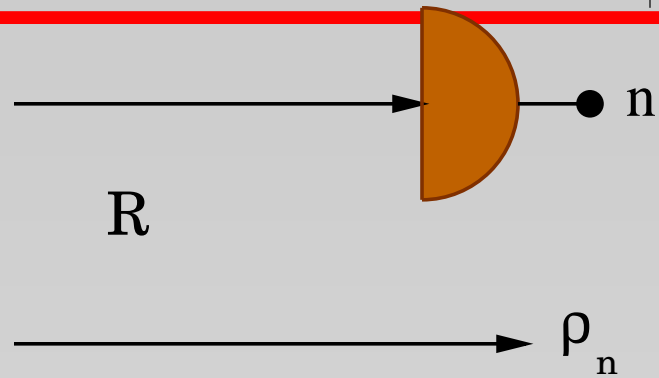
- Essentially **any garbage state is faithful** (invertibility is a *dense* condition).

Faithful states

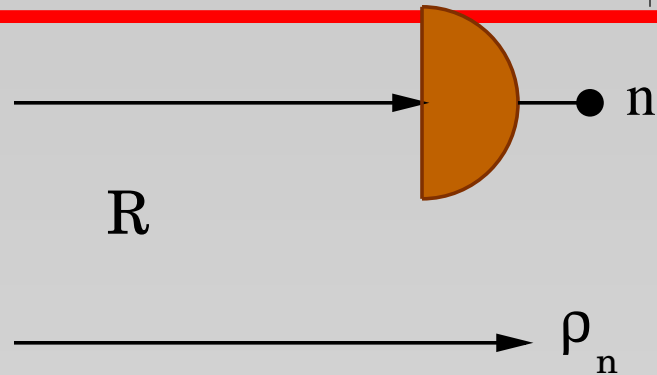
- Essentially **any garbage state is faithful** (invertibility is a *dense* condition).
- However, the knowledge of the map \mathcal{E} from the measured output state $R_{\mathcal{E}}$ will be affected by increasingly large statistical errors for input state R approaching a non-faithful one.



Quantum Calibration of a measuring apparatus



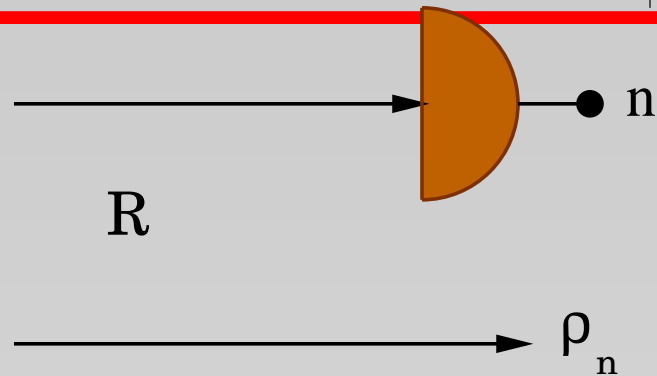
Quantum Calibration of a measuring apparatus



- The calibration is achieved by determining experimentally the POVM using a faithful state as in figure.

$$P_n = p(n)[\mathcal{R}^{-1}(\rho_n)]^\top.$$

Quantum Calibration of a measuring apparatus



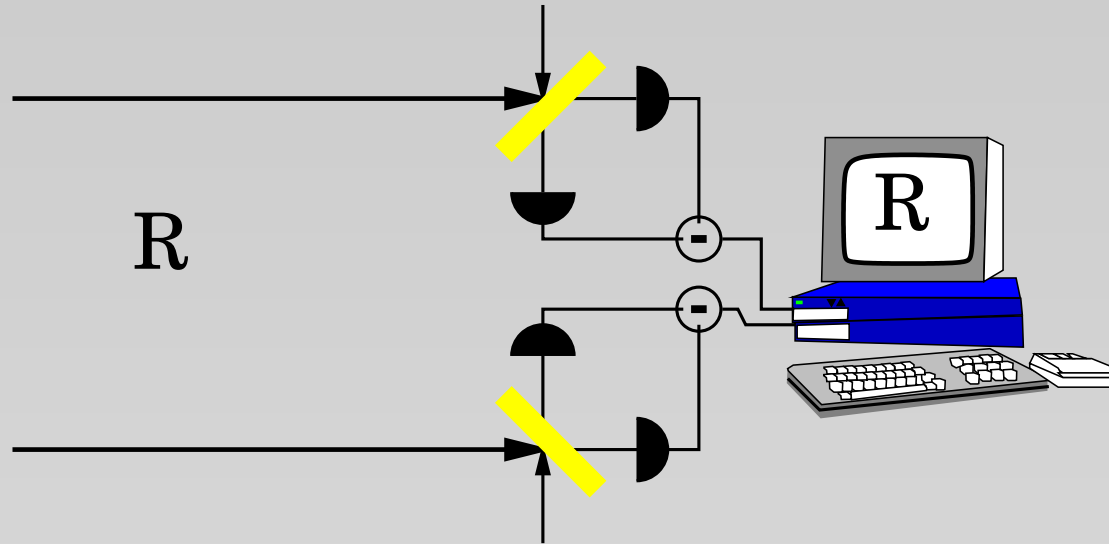
- The calibration is achieved by determining experimentally the POVM using a faithful state as in figure.

$$P_n = p(n)[\mathcal{R}^{-1}(\rho_n)]^\top.$$

- $p(n)$ probability of the outcome n ,
- ρ_n conditioned state, to be determined by quantum tomography,
- \mathcal{R} associated map of the faithful state R .

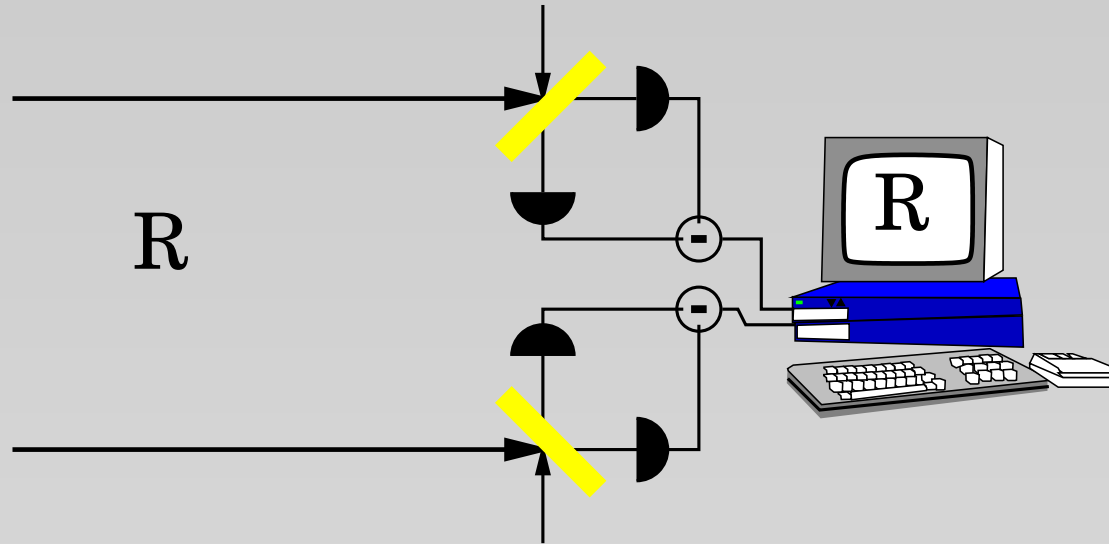
Quantum Calibration of a measuring apparatus

- *Precalibration*

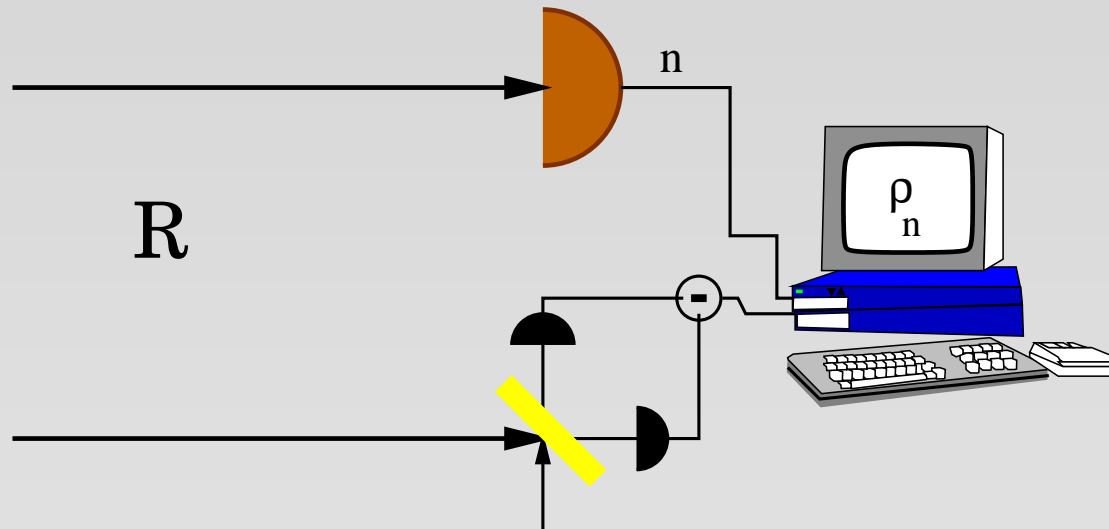


Quantum Calibration of a measuring apparatus

- *Precalibration*



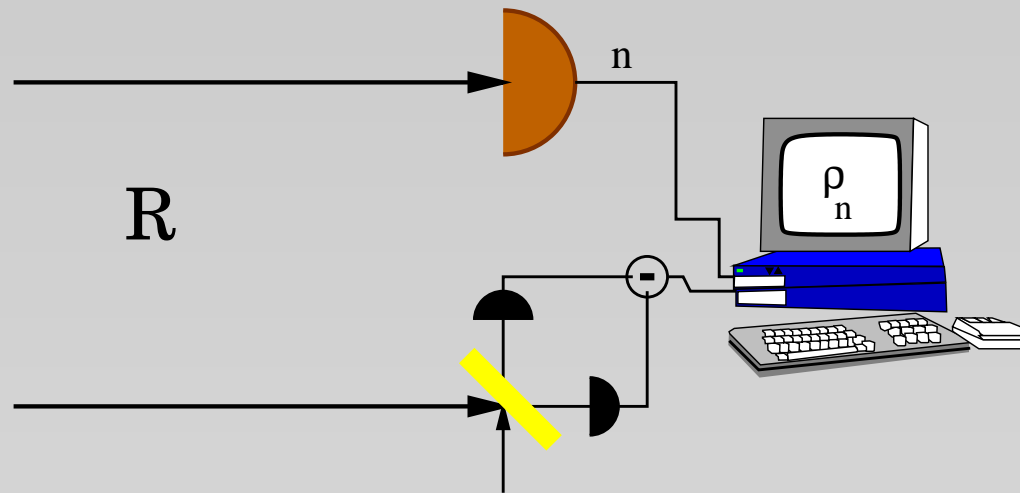
- *Calibration*



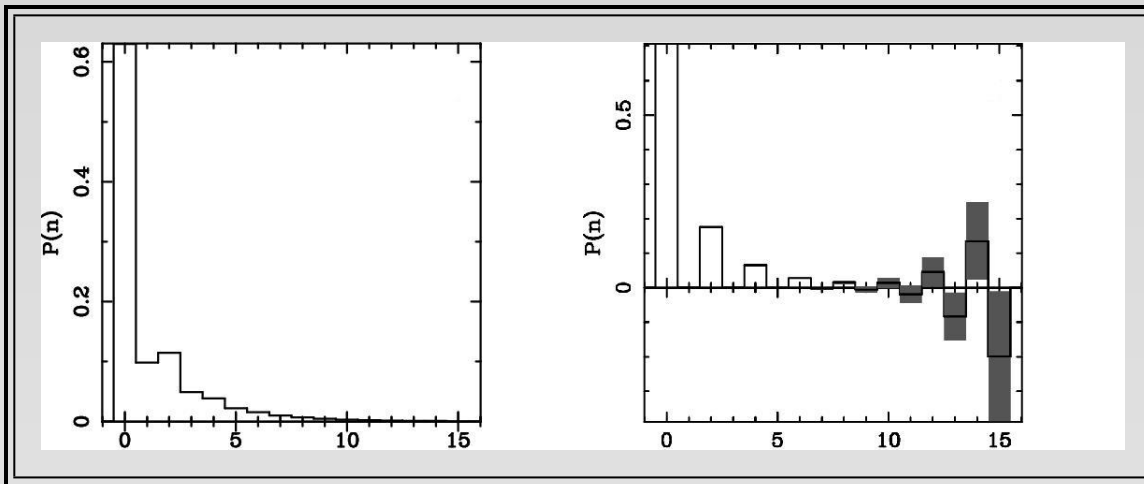
Quantum Calibration of a measuring apparatus

In principle we need only two tomographers and a single faithful state to calibrate any measuring apparatus.

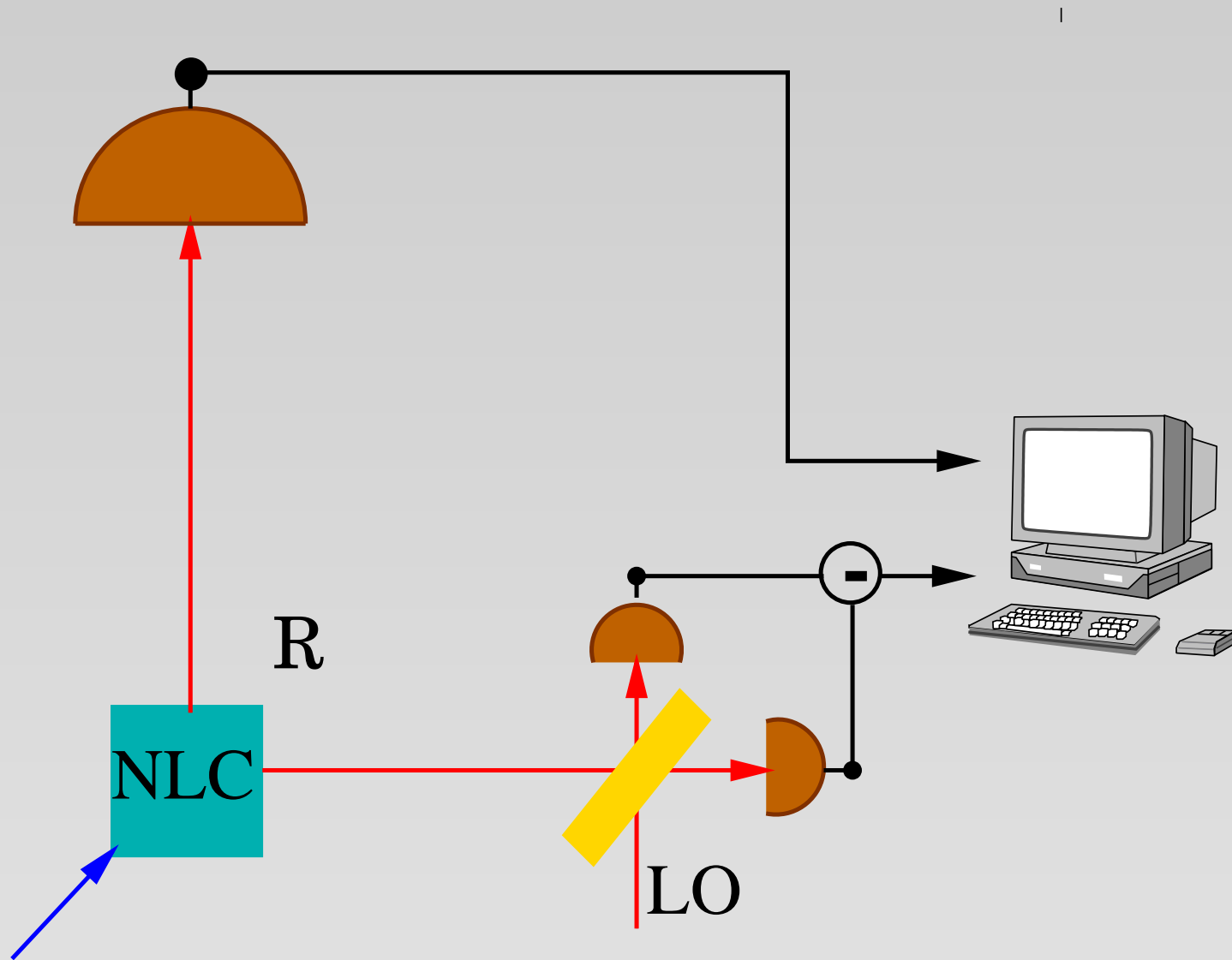
Quantum Calibration of a measuring apparatus



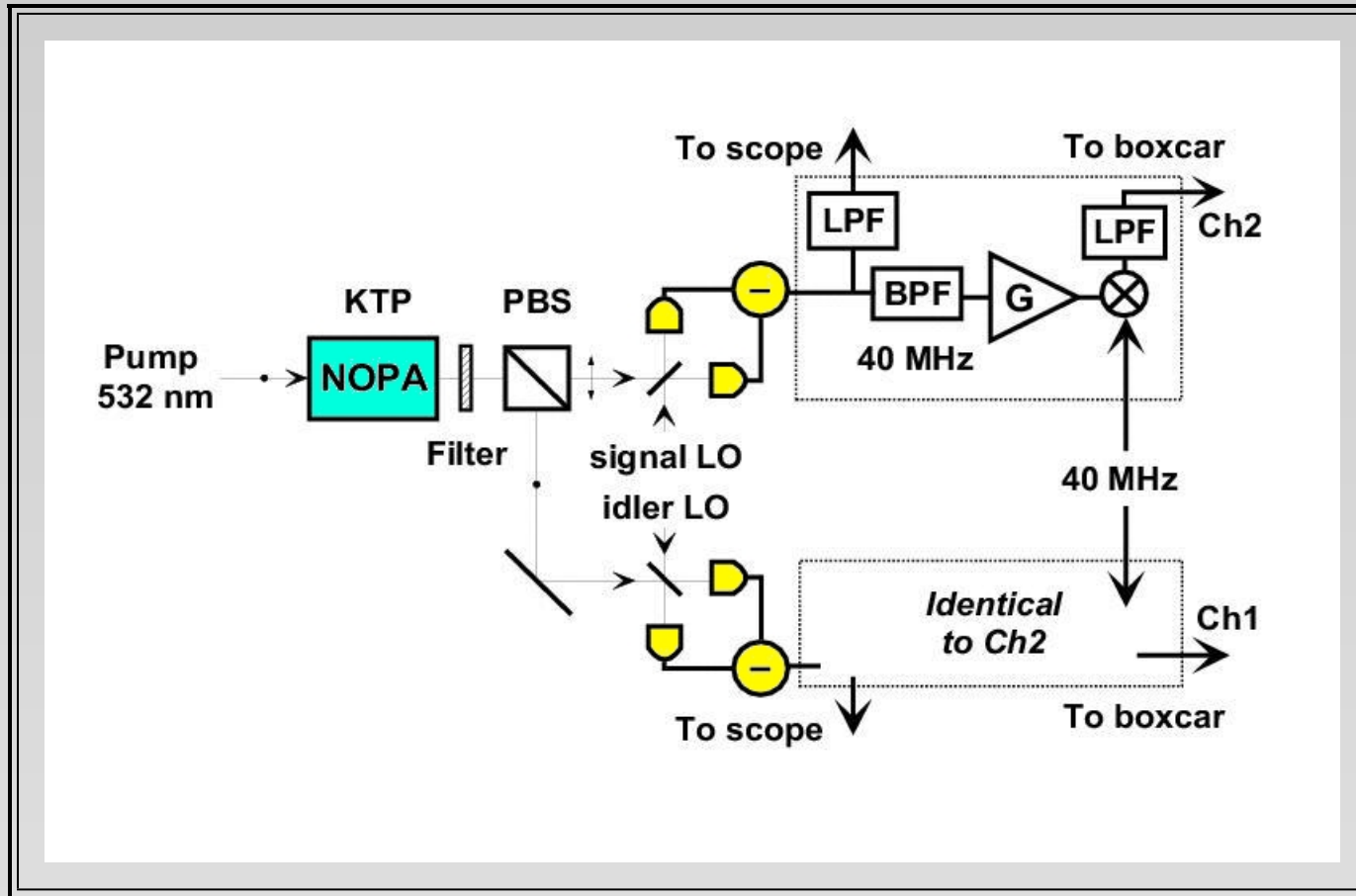
- Using a "calibrated observable" the measurement is "unbiased" (at expense of some increasing statistical error).



Quantum calibration of a photodetector



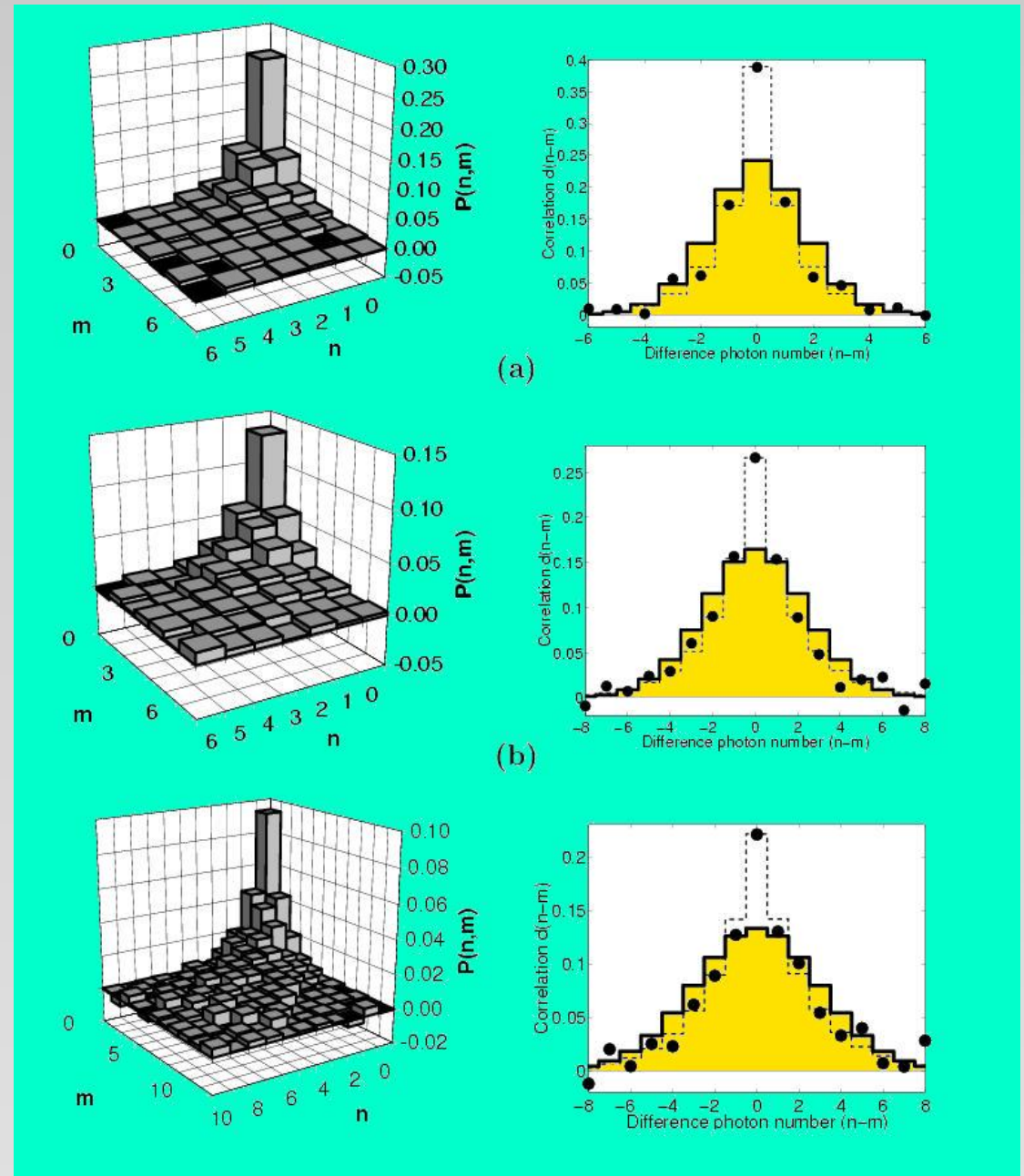
Tomography of a twin-beam



Measurement of the joint photon-number probability distribution of a twin-beam: schematic of the experimental setup. NOPA, non-degenerate optical parametric amplifier; LOs, local oscillators; PBS, polarizing beam splitter; LPFs, low-pass filters; BPF, band-pass filter; G, electronic amplifier. Electronics in the two channels are identical.

Results

Left: Measured joint photon-number probability distributions for the twin-beam state. Right: Difference photon number distributions corresponding to the left graphs (filled circles, experimental data; solid lines, theoretical predictions; dashed lines, difference photon number distributions for two independent coherent states with the same total mean number of photons and $\bar{n} = \bar{m}$.) (a) 400000 samples, $\bar{n} = \bar{m} = 1.5$, $N = 10$; (b) 240000 samples, $\bar{n} = 3.2$, $\bar{m} = 3.0$, $N = 18$; (c) 640000 samples, $\bar{n} = 4.7$, $\bar{m} = 4.6$, $N = 16$. The measured distributions exhibit up to 1.9 dB of quantum correlation between the signal and idler photon numbers.



Homodyne calibration of a photodetector

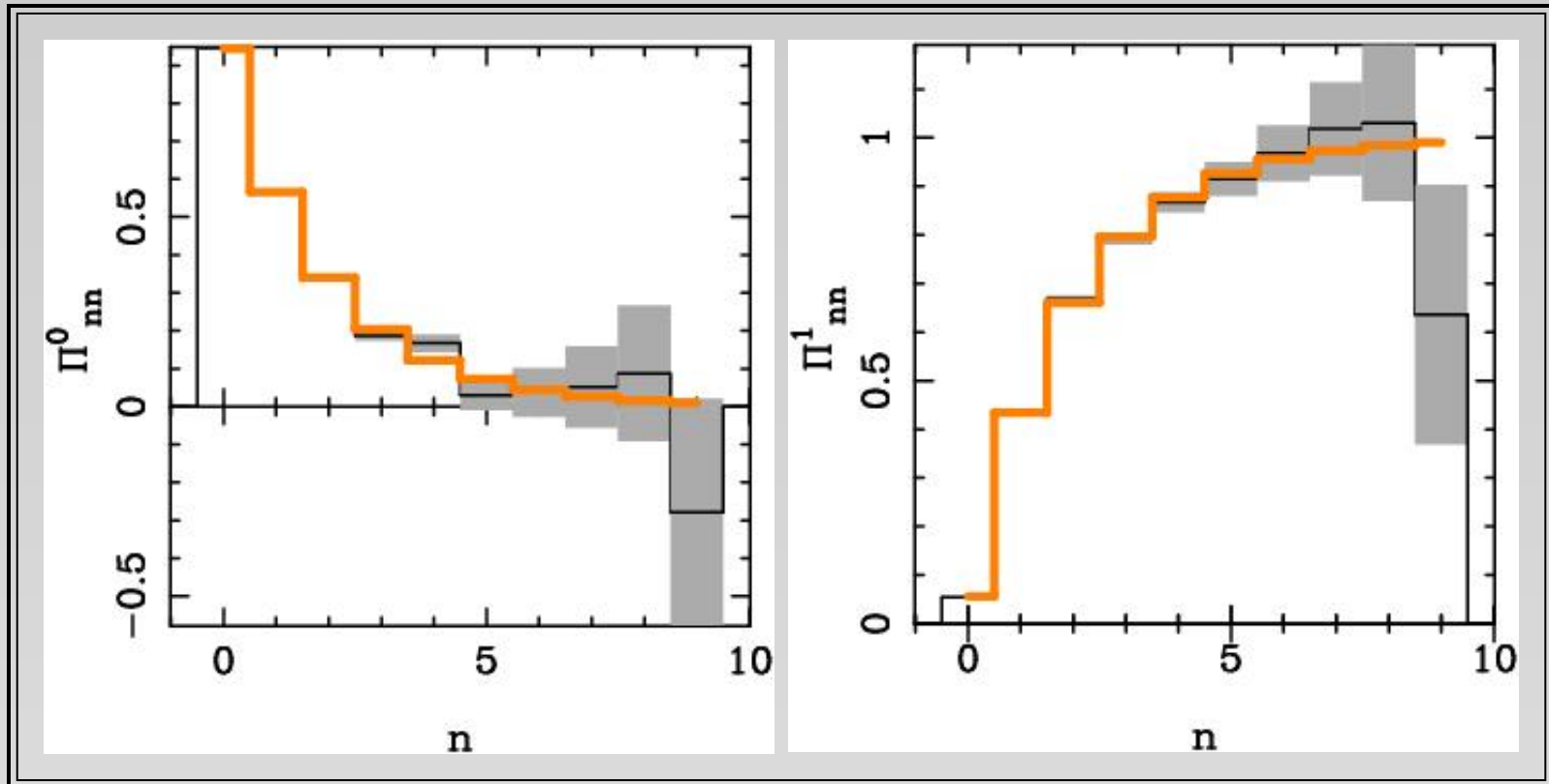


Figure 1: Homodyne tomography of an On/Off photo-detector with quantum efficiency $\eta = 0.4$ and thermal noise photon number $\nu = 0.1$. The reconstruction is obtained by pattern-function averaging of $1.5 \cdot 10^6$ data, for homodyne quantum efficiency $\eta = 0.9$ and twin beam thermal photon $\bar{n} = 3$.

Homodyne calibration of a photodetector

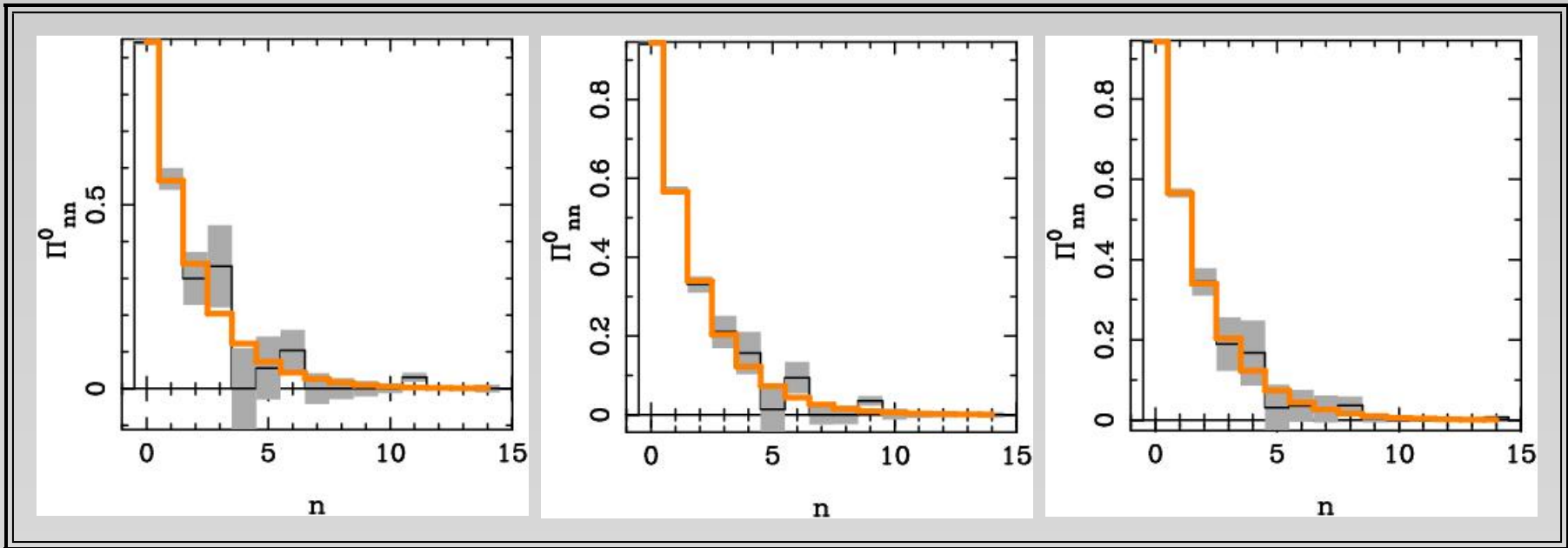
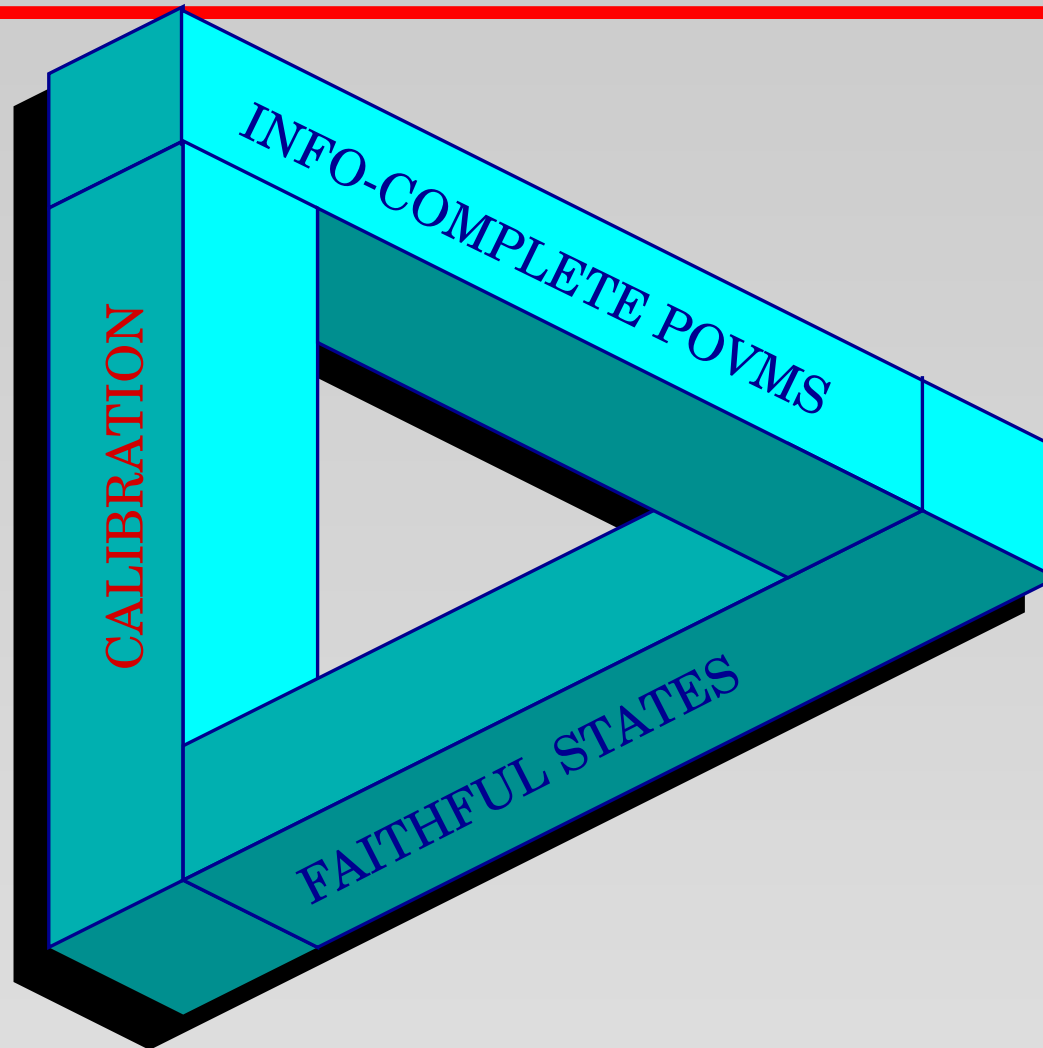


Figure 2: Homodyne tomography of an On/Off photodetector with quantum efficiency $\eta = 0.4$ and thermal noise photon number $\nu = 0.1$, with $\bar{n} = 3$ photons in the twin-beam. The ML estimation of the diagonal of the only Off POVM element are reported for different values of sample size N and homodyne quantum efficiency η_H . Left: $N = 10^5$, $\eta_H = 0.7$; Middle: $N = 10^4$, $\eta_H = 0.9$; Right: $N = 10^6$, $\eta_H = 0.7$.

Quantum Mechanics: physical axioms?



Informationally complete POVM's = **calibrators**: *"the quantum standards of the International Bureau of Weights and Measures à Paris"* — Chris Fuchs.



http://www.qubit.it

QUit

quantum information theory group

Sunday, 23 Nov 2003

- people
- faq
- research
- educational
- for visitors
- sponsors

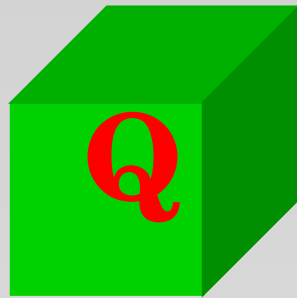


upcoming events

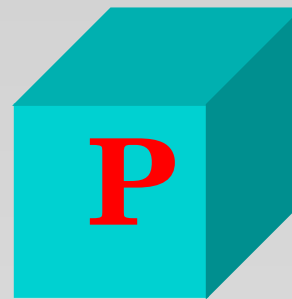
The fun with **Quantum Information** is that you can study the foundations of the enigmatic world of Quantum Mechanics, and, at the same time, you make something useful for practical applications

[Workshop on: Quantum Information Processing for Quantum Communications > November 2003 - July 2004](#)

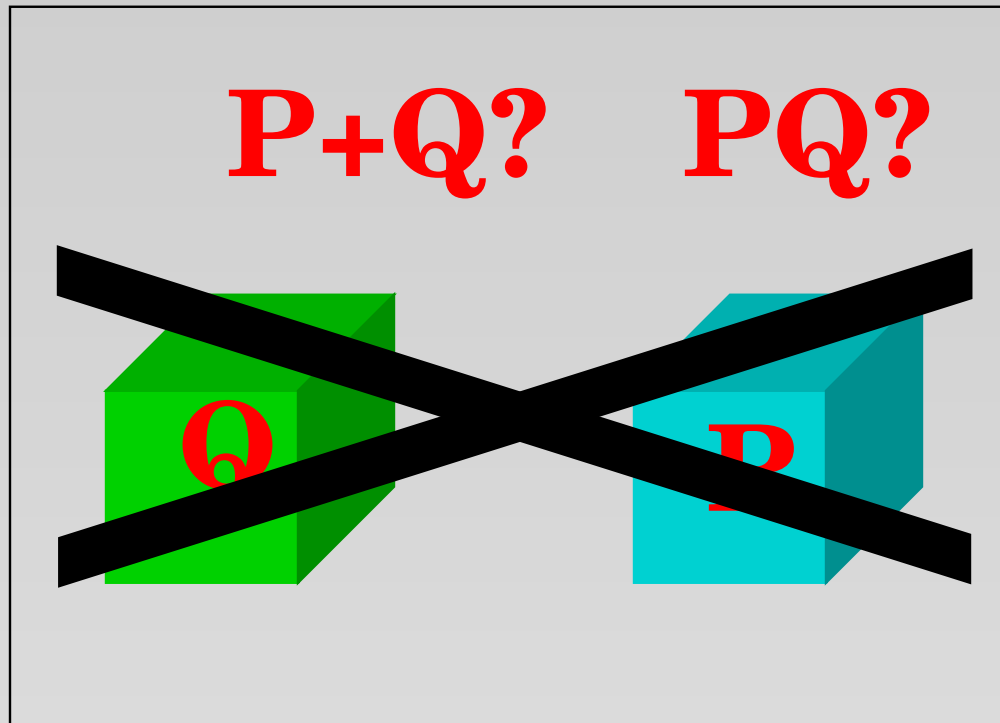
Q position



P momentum



Quantum observables and measurement devices



[*]