

#### JOHN TEMPLETON FOUNDATION

SUPPORTING SCIENCE ~ INVESTING IN THE BIG QUESTIONS Project: *Quantum Causal Structures*, ID 60609



#### The Galileo principle for general dynamical systems: Lorentz transformations from Quantum Theory and a little bit more (homogeneity, isotropy, and locality)

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*Frascati Labs* Training school: *Are spin-statistics connection and quantum theory exact?* ... 19-21 December 2016

# New *algorithmic* paradigm: the physical law as an *algorithm*





# Old *mechanistic* paradigm: the physical law as a *mechanism*









### The information-theoretic paradigm

## The first opportunity of solving the problem of axiomatization of physics

The investigations on the foundations of geometry suggest the problem: To treat in the same manner by means of axioms, those physical sciences in which mathematics plays an important part; in the first rank are the theory of probabilities and mechanics.

David Hilbert



The VI Hilbert problem

## The information-theoretic paradigm

## The first opportunity of solving the problem of axiomatization of physics

Axiomatizing the theory of probabilities was a realistic goal: Kolmogorov accomplished this in 1933. The word 'mechanics' without a qualifier, however, is a Trojan horse."

Benjamin Yandell





The VI Hilbert problem

Ben H. Yandell

## The information-theoretic paradigm

The first opportunity of solving the problem of axiomatization of physics

PROGRAM

Derive Physics from "principles" stated in form of <u>purely mathematical</u> <u>axioms without physical primitives</u>, but having a thorough physical interpretation

physical primitives: mass, force, rods, clocks,...



The VI Hilbert problem





## High-energy/short-distance correspondence breakdown

- Colliding two particles at Planck energy (.54MWh) produces a black hole!
- A particle with a too large mass (2.18\*10<sup>-5</sup> g) becomes a black hole!
- GR-QFT patching: Planck scale





#### electron

### The Planck scale

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#### The information paradox

#### QT preserves information! Do black hole preserve it?



#### Solution: Holographic principle?



Causality paradox Pre-established (QT) vs dynamical (GR) causality



#### Localization issue in QFT

#### 10.8 Conclusion

Malament claims that his theorem justifies the belief that,

... in the attempt to reconcile quantum mechanics with relativity theory... one is driven to a field theory; all talk about "particles" has to be understood, at least in principle, as talk about the properties of, and interactions among, quantized fields. (Malament 1996, 1)

In order to buttress Malament's argument for this claim, we provided two further results (Theorems 3 and 5) which show that the conclusion continues We then went on to show that QFT does not permit an ontology of localizable particles; and so, strictly speaking, our talk about localizable particles is a fiction.

— albeit, if we understand this talk as really being about the properties of, and interactions among, quantized fields. Indeed, modulo the standard quantum measurement problem, RQFT has no trouble explaining the appearance of macroscopically well-localized objects, and shows that our talk of particles, though a *façon de parler*, has a legitimate role to play in empirically testing the theory.

#### No Place for Particles in Relativistic Quantum Theories?

Hans Halvorson Princeton University

Rob Clifton University of Pittsburgh

#### Localization issue in QFT

Physicists routinely describe the universe as being made of tiny subatomic particles that push and pull on one another by means of force fields. They call their subject "particle physics" and their instruments "particle accelerators." They hew to a Lego-like model of the world. But this view sweeps a little-known fact under the rug: the particle interpretation of quantum physics, as well as the field interpretation, stretches our conventional notions of "particle" and "field" to such an extent that ever more people think the world might be made of something else entirely.

The problem is not that physicists lack a valid theory of the subatomic realm. They do have one: it is called quantum field theory. Theorists developed it between the late 1920s and early 1950s by merging the earlier theory of quantum mechanics with Einstein's special theory of relativity. Quantum field theory provides the conceptual underpinnings of the Standard Model of particle physics, which describes the fundamental building blocks of matter and their interactions in one common framework. In terms of empirical precision, it is the most successful theory in the history of science. Physicists use it every day to calculate the aftermath of particle collisions, the synthesis of matter in the big bang, the ex-

**Meinard Kuhlmann,** a philosophy professor at Bielefeld University in Germany, received dual degrees in physics and in philosophy and has worked at the universities of Oxford, Chicago and Pittsburgh. As a student, he had an inquisitive reputation. "I would ask a lot of questions just for fun and because they produced an entertaining confusion," he says.

*ican* articles. However compelling it might appear, it is not at

there is an electron field as surely as there is an electron. At the

same time, the force fields are quantized rather than continu-

ous, which gives rise to particles such as the photon. So the distinction between particles and fields appears to be artificial, and

physicists often speak as if one or the other is more fundamen-

tal. Debate has swirled over this point-over whether quantum

field theory is ultimately about particles or about fields. It started as a battle of titans, with eminent physicists and philoso-

phers on both sides. Even today both concepts are still in use for

illustrative purposes, although most physicists would admit

that the classical conceptions do not match what the theory

says. If the mental images conjured up by the words "particle"

and "field" do not match what the theory says, physicists and

For starters, the two categories blur together. Quantum field theory assigns a field to each type of elementary particle, so

all satisfactory.



Physicists speak of the world as being made of particles and force fields, but it is not at all clear what particles and force fields actually are in the quantum realm. The world may instead consist of bundles of properties, such as color and shape **Bu Meinard Kuhlmann** 

philosophers must figure out what to put in their place. With the two standard, classical options gridlocked, some philosophers of physics have been formulating more radical alternatives. They suggest that the most basic constituents of the material world are intangible entities such as relations or properties. One particularly radical idea is that everything can be reduced to





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But now how to derive the "mechanics", namely: quantum field theory, special relativity, ... without using physical primitives, mechanics, kinematics, space, time,...?

For example: how would you formulate the principle of relativity in algorithmic terms?

## Algorithm -> discreteness!

discrete



continuum





"But you have correctly grasped the drawback that the continuum brings. If the molecular view of matter is the correct (appropriate) one, i.e., if a part of the universe is to be represented by a finite number of moving points, then the continuum of the present theory contains too great a manifold of possibilities. I also believe that this too great is responsible for the fact that our present means of description miscarry with the quantum theory. The problem seems to me how one can formulate statements about a discontinuum without calling upon a continuum (space-time) as an aid; the latter should be banned from the theory as a supplementary construction not justified by the essence of the problem, which corresponds to nothing "real". <u>But we still lack</u> the mathematical structure unfortunately. How much have I already plagued myself in this way!"

John Stachel in *From Quarks to Quasars: Philosophical Problems of Modern Physics,* University of Pittsburg Press, pag. 379

#### A new mathematics: geometric group theory

The geometrization of group theory



#### Mikhail Gromov



#### JOHN TEMPLETON FOUNDATION

#### SUPPORTING SCIENCE ~ INVESTING IN THE BIG QUESTIONS

Project: A Quantum-Digital Universe, Grant ID: 43796

• Mechanics (QFT) derived in terms of countably many quantum systems in interaction





Paolo Perinotti



Marco Erba



Alessandro Bisio



Alessandro Tosini



Nicola Mosco









#### The Weyl QW

Solution Minimal dimension for nontrivial unitary A: s=2

Unitarity + isotropy  $\Rightarrow$  for d=3 the only Cayley is the BCC!!

Unitary operator: 
$$A = \int_{B}^{\oplus} d\mathbf{k} A_{\mathbf{k}}$$

Two QWs connected by P

$$A_{\mathbf{k}}^{\pm} = -i\sigma_x(s_x c_y c_z \pm c_x s_y s_z)$$
  

$$\mp i\sigma_y(c_x s_y c_z \mp s_x c_y s_z)$$
  

$$-i\sigma_z(c_x c_y s_z \pm s_x s_y c_z)$$
  

$$+ I(c_x c_y c_z \mp s_x s_y s_z)$$

$$s_{\alpha} = \sin \frac{k_{\alpha}}{\sqrt{3}}$$
$$c_{\alpha} = \cos \frac{k_{\alpha}}{\sqrt{3}}$$

#### Physical interpretation: Weyl equation

$$i\partial_t \psi(t) \simeq \frac{i}{2} [\psi(t+1) - \psi(t-1)] = \frac{i}{2} (A - A^{\dagger}) \psi(t)$$

 $\frac{i}{2}(A_{\mathbf{k}}^{\pm} - A_{\mathbf{k}}^{\pm\dagger}) = + \sigma_x(s_x c_y c_z \pm c_x s_y s_z) \quad \text{"Hamiltonian"} \\ \pm \sigma_y(c_x s_y c_z \mp s_x c_y s_z) \\ + \sigma_z(c_x c_y s_z \pm s_x s_y c_z)$ 

$$k \ll 1$$
  $\square$   $i\partial_t \psi = \frac{1}{\sqrt{3}} \sigma^{\pm} \cdot \mathbf{k} \psi$  So Weyl equation!  $\sigma^{\pm} := (\sigma_x, \pm \sigma_y, \sigma_z)$ 

Two QCAs connected by P

$$A_{\mathbf{k}}^{\pm} = -i\sigma_x(s_x c_y c_z \pm c_x s_y s_z)$$
  

$$\mp i\sigma_y(c_x s_y c_z \mp s_x c_y s_z)$$
  

$$-i\sigma_z(c_x c_y s_z \pm s_x s_y c_z)$$
  

$$+ I(c_x c_y c_z \mp s_x s_y s_z)$$

$$s_{\alpha} = \sin \frac{k_{\alpha}}{\sqrt{3}}$$
$$c_{\alpha} = \cos \frac{k_{\alpha}}{\sqrt{3}}$$

#### D'Ariano, Perinotti, PRA 90 062106 (2014)

### Dirac QW



<u>Local</u> coupling:  $A_{\mathbf{k}}$  coupled with its inverse with off-diagonal identity block matrix

$$E_{\mathbf{k}}^{\pm} = \begin{pmatrix} nA_{\mathbf{k}}^{\pm} & imI\\ imI & nA_{\mathbf{k}}^{\pm\dagger} \end{pmatrix}$$
$$n^{2} + m^{2} = 1 \qquad n, m \in \mathbb{R}$$

$$\omega_{\pm}^{E}(\mathbf{k}) = \cos^{-1}[n(c_{x}c_{y}c_{z} \mp s_{x}s_{y}s_{z})]$$

Dirac in relativistic limit  $k \ll m \ll 1$ 

m: mass,  $m^2 \le 1$ n<sup>-1</sup>: refraction index





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Bisio, D'Ariano, Perinotti, Ann. Phys. 368 177 (2016)

### Maxwell QW



 $c^{\mp}(\mathbf{k}) = c \left( 1 \pm \frac{k_x k_y k_z}{|\mathbf{k}|^2} \right)$ 

 $k_z$ 

 $2\vec{n}_{\mathbf{k}}$ 

k

 $\vec{v}_g(\mathbf{k})$ 

 $k_y$ 

 $k_x$ 

$$M_{\mathbf{k}}^{\pm} = A_{\mathbf{k}}^{\pm} \otimes A_{\mathbf{k}}^{\pm *}$$
$$F^{\mu}(\mathbf{k}) = \int \frac{\mathrm{d}\,\mathbf{q}}{2\pi} f(\mathbf{q}) \tilde{\psi}(\frac{\mathbf{k}}{2} - \mathbf{q}) \sigma^{\mu} \varphi(\frac{\mathbf{k}}{2} + \mathbf{q})$$

Maxwell in relativistic limit  $k \ll 1$ Boson: emergent from convolution of fermions (De Broglie neutrino-theory of photon)







The theory contains its own LTM standards

$$x = \frac{x_{[m]}}{a_*} \in \mathbb{Z}, \quad t = \frac{t_{[sec]}}{t_*} \in \mathbb{N}, \quad m = \frac{m_{[kg]}}{m_*} \in [0, 1]$$
$$m_* \simeq \frac{1}{\sqrt{3\pi}} \frac{\hbar k}{c(k) - c(0)} \qquad \qquad \begin{cases} c \equiv c(0) = \frac{a_*}{t_*}\\ \hbar = m_* a_* c \end{cases}$$

Heuristic argument of the mini-black-hole:

 $m_{st}$  Planck mass

from the relativistic limit




## Fidelity of the quantum walk with QFT

Fidelity with Dirac for a narrowband packets in the relativistic limit  $k\simeq m\ll 1$ 

$$F = \left| \left\langle \exp\left[ -iN\Delta(\mathbf{k}) \right] \right\rangle \right|$$

$$\Delta(\mathbf{k}) := (m^2 + \frac{k^2}{3})^{\frac{1}{2}} - \omega^E(\mathbf{k})$$
  
=  $\frac{\sqrt{3}k_x k_y k_z}{(m^2 + \frac{k^2}{3})^{\frac{1}{2}}} - \frac{3(k_x k_y k_z)^2}{(m^2 + \frac{k^2}{3})^{\frac{3}{2}}} + \frac{1}{24}(m^2 + \frac{k^2}{3})^{\frac{3}{2}} + \mathcal{O}(k^4 + N^{-1}k^2)$ 

relativistic proton:  $N \simeq m^{-3} = 2.2 * 10^{57} \Rightarrow t = 1.2 * 10^{14} \text{s} = 3.7 * 10^{6} \text{ y}$ 

UHECRs: 
$$k = 10^{-8} \gg m \Rightarrow N \simeq k^{-2} = 10^{16} \Rightarrow 5 * 10^{-28}$$
 s

Special relativity translated into quantum-algorithmic terms



**Relativity Principle:** Invariance of the dynamical law with the inertial frame

*Inertial frame:* a reference frame where the Newton inertia law holds for a mechanically isolated system



**Poincaré group:** group of changes of inertial frame that leave the dynamical law invariant.

**Relativity Principle:** Invariance of the dynamical law with the inertial frame



*Inertial frame:* a reference frame where energy and momentum are conserved for a mechanically *isolated* system.

**Poincaré group:** group of changes of inertial frame that leave the dynamical law invariant.

**Relativity Principle:** Invariance of the dynamical law with the inertial frame

**Inertial frame:** Representation of the dynamical law for given *values* of the constants of motion for an *isolated* system.

**Dynamical law:** expressed in terms of the values of the constants of motion.

**Poincaré group:** group of changes of inertial frame that leave the dynamical law invariant.

**Relativity Principle:** Invariance of the dynamical law with the inertial frame

**Inertial frame:** Representation of the physical law in terms of eigenspaces of the constants of the dynamics  $k:=(\omega,\mathbf{k})$ 

Dynamical law: eigenvalue equation

$$A_{\mathbf{k}}\psi(\mathbf{k},\omega) = e^{i\omega}\psi(\mathbf{k},\omega)$$

**Poincaré group:** group of changes of representations in terms of eigenspaces of the constants of dynamics that leave the eigenvalue equation invariant.

## Special Relativity from Quantum theory

- Mathematical statement: invariance of eigenvalue equation under change of representation.
- Physical interpretation: invariance of the physical law under change of inertial reference frame.

#### •*m=0*

i.Deformed Lorentz group SO(1,3)

ii.Lorentz transformations are perfectly recovered for k,m«1

iii. For *k~1*: *Double Special Relativity* (Camelia-Smolin) [Relative locality]



Fig. 2 (Colors online) The red surfaces represents the orbit of a wavevector  $\mathbf{k} = (k_x, 0, 0)$ under the action of the deformed rotations  $\mathcal{R} = \mathcal{D}^{(f)^{-1}} \circ R \circ \mathcal{D}^{(f)}$  where f is the function defined in Eq. (28). Left:  $k_x = 0.07$ . Middle:  $k_x = 0.2$  Right:  $k_x = 0.4$ 

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FIG. 2: The distortion effects of the Lorentz group for the discrete Planck-scale theory represented by the quantum walk in Eq. (6). Left figure: the orbit of the wavevectors  $\mathbf{k} = (k_x, 0, 0)$ , with  $k_x \in \{.05, .2, .5, 1, 1.7\}$  under the rotation around the z axis. Right figure: the orbit of wavevectors with  $|\mathbf{k}| = 0.01$  for various directions in the  $(k_x, k_y)$  plane under the boosts with  $\boldsymbol{\beta}$  parallel to  $\mathbf{k}$  and  $|\boldsymbol{\beta}| \in [0, \tanh 4]$ .

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*m>0* Deformed <u>De Sitter</u> group SO(1,4)

## Particle notion without mechanics

 Mathematical statement: irreducible representation of the group of invariance of dynamics (deformed Poincaré group).



- The Brillouin zone separates into four Poincaré-invariant regions diffeomorphic to balls, corresponding to four different <u>particles</u>.

## Internal coherence of the theory

- Mathematical statement: topology of domain of the particle mass is a circle
- Physical interpretation: proper time is discrete!





## Besides, ...

Vladimir Ignatowski derived the Lorentz transformations from homogeneity, isotropy, reciprocity, ...

W. A. von Ignatowsky, Verh. Deutsch. Phys. Ges. 12 (1910), 788–796; Phys. Zeitsch. 11 (1910), 972–976;

see: S. Liberati, Annals of Physics 298, 167–185 (2002)



# Conventionality

### Homogeneity, simultaneity, ... unfalsifiable principles?





# γνῶθι σεαυτὸν

