

Physics without physics

Giacomo Mauro D'Ariano
Università degli Studi di Pavia

Quantum Theory: from foundations to technologies – QTFT
Linnaeus University, Växjö, June 12-17 2016

Program

Deriving the whole Physics axiomatically

from “principles” stated in form of purely mathematical axioms without physical primitives,

but having a thorough physical interpretation.

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from “principles” stated in form of purely mathematical axioms without physical primitives,

but having a thorough physical interpretation.

Examples of physical primitives:


mass, force, clocks, rods, ...

Physical interpretation from where?

from experience and from un-axiomatized physics ...

Principles for Quantum Theory



 Selected for a [Viewpoint](#) in *Physics*
PHYSICAL REVIEW A **84**, 012311 (2011)

Informational derivation of quantum theory

Giulio Chiribella*

Perimeter Institute for Theoretical Physics, 31 Caroline Street North, Ontario, Canada N2L 2Y5[†]

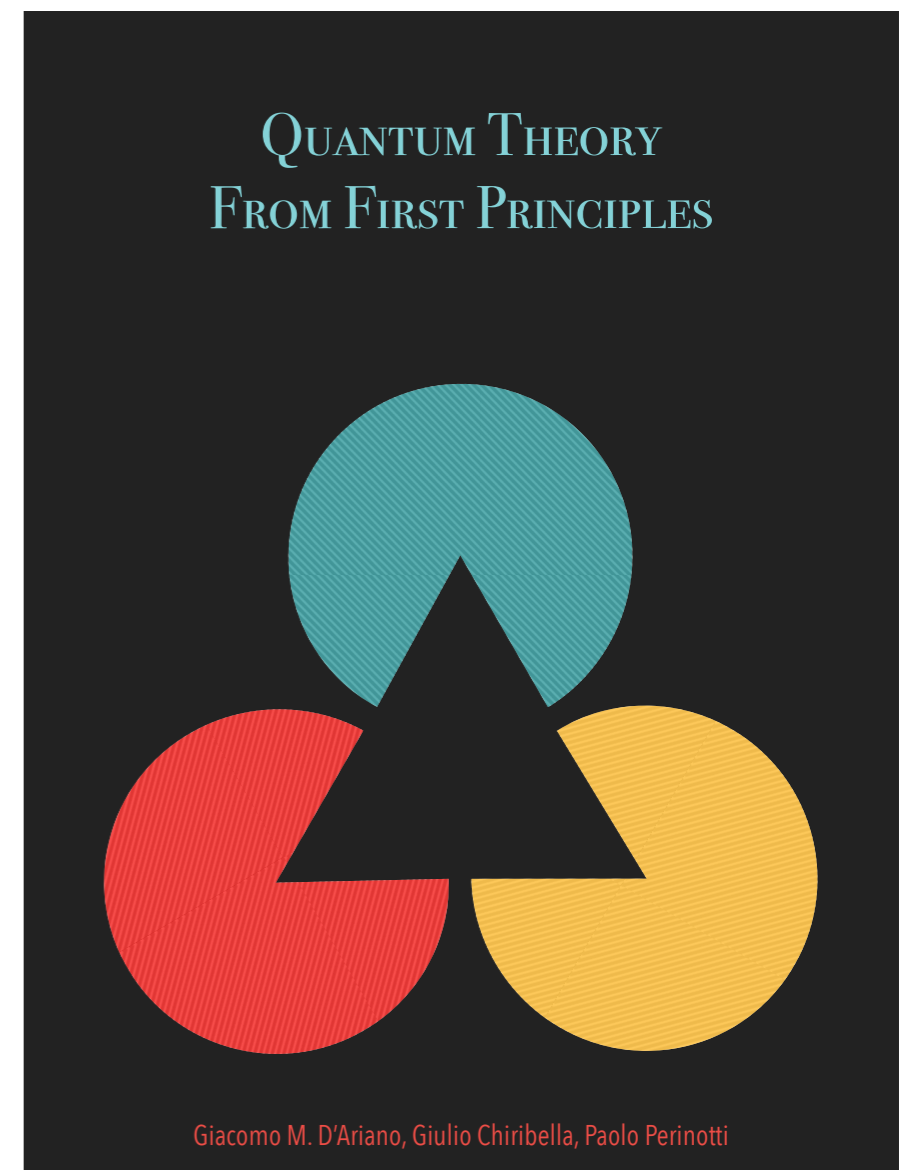
Giacomo Mauro D'Ariano[‡] and Paolo Perinotti[§]

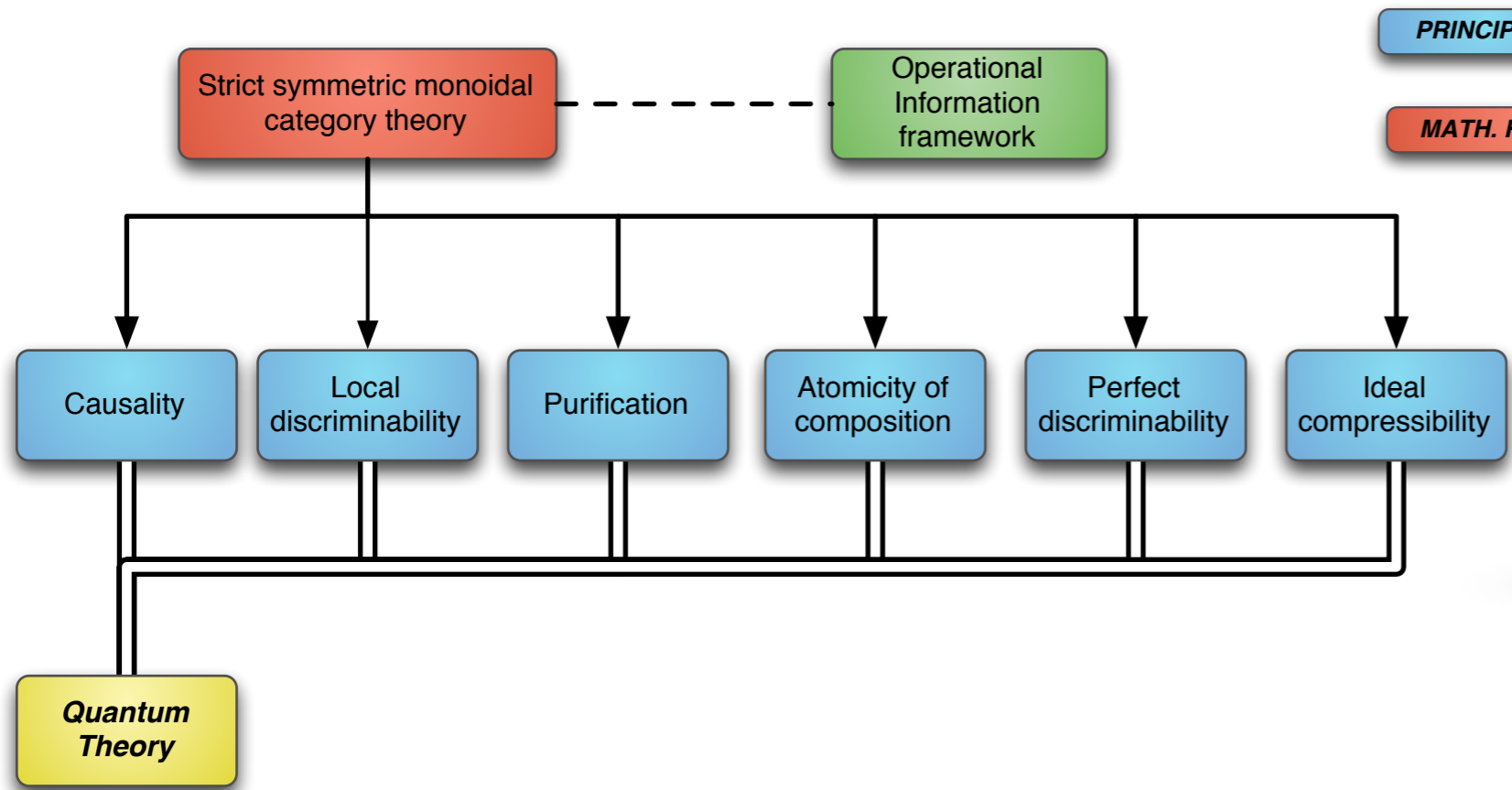
QUIT Group, Dipartimento di Fisica "A. Volta" and INFN Sezione di Pavia, via Bassi 6, I-27100 Pavia, Italy^{||}
(Received 29 November 2010; published 11 July 2011)

We derive quantum theory from purely informational principles. Five elementary axioms—causality, perfect distinguishability, ideal compression, local distinguishability, and pure conditioning—define a broad class of theories of information processing that can be regarded as standard. One postulate—purification—singles out quantum theory within this class.

DOI: [10.1103/PhysRevA.84.012311](https://doi.org/10.1103/PhysRevA.84.012311)

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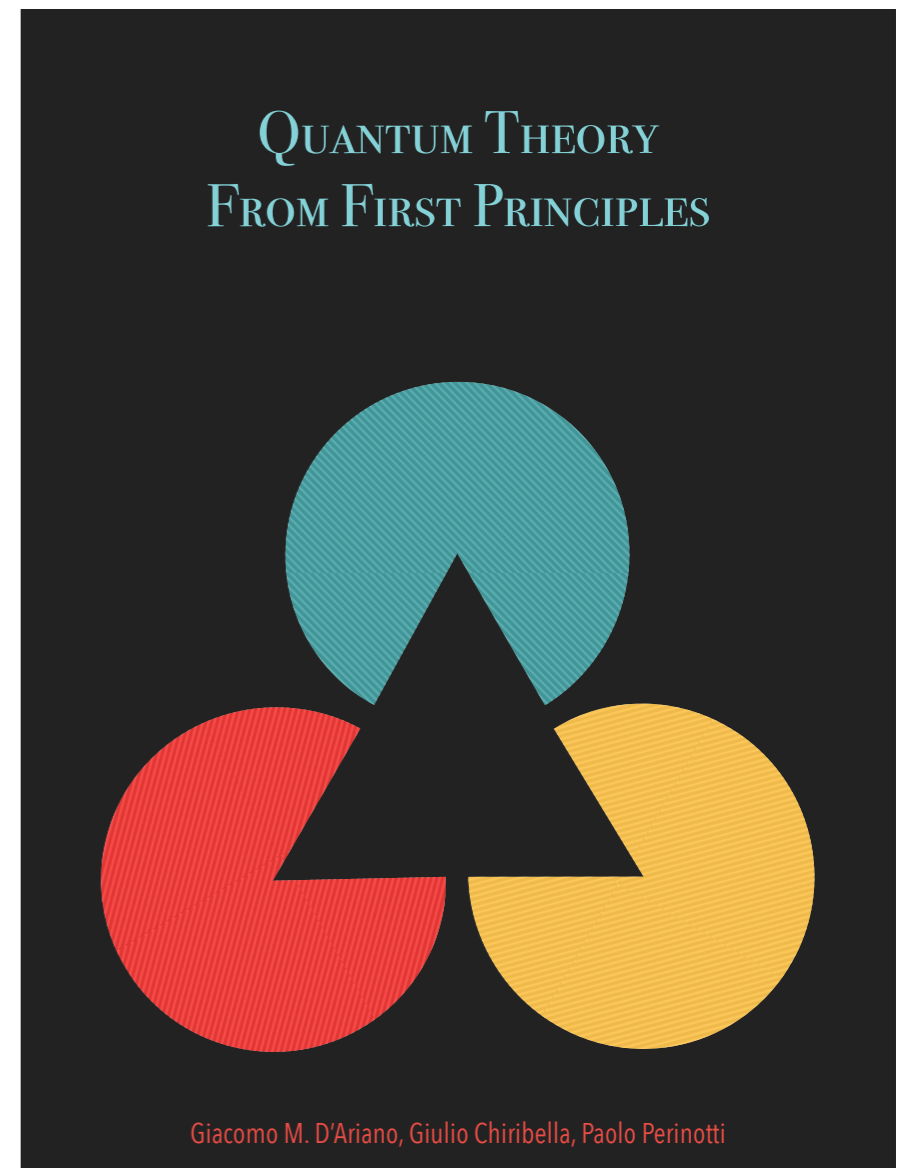
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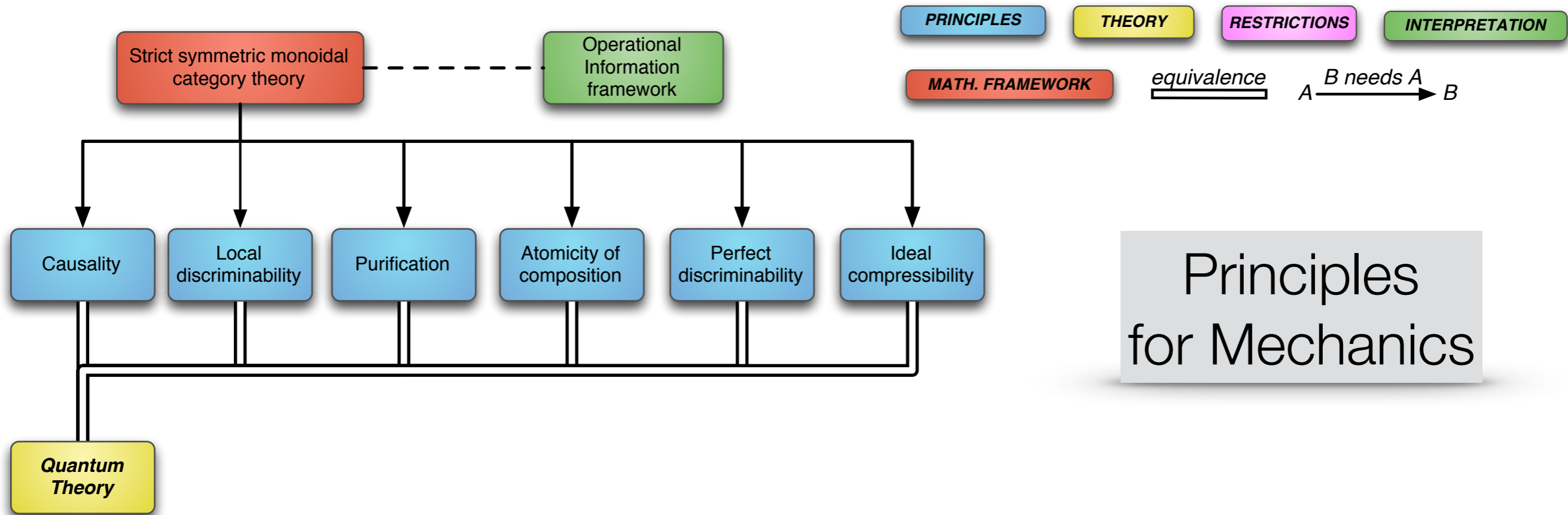
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Project: *A Quantum-Digital Universe*, Grant ID: 43796



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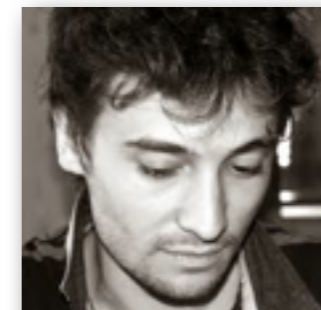
Alessandro Bisio



Alessandro Tosini



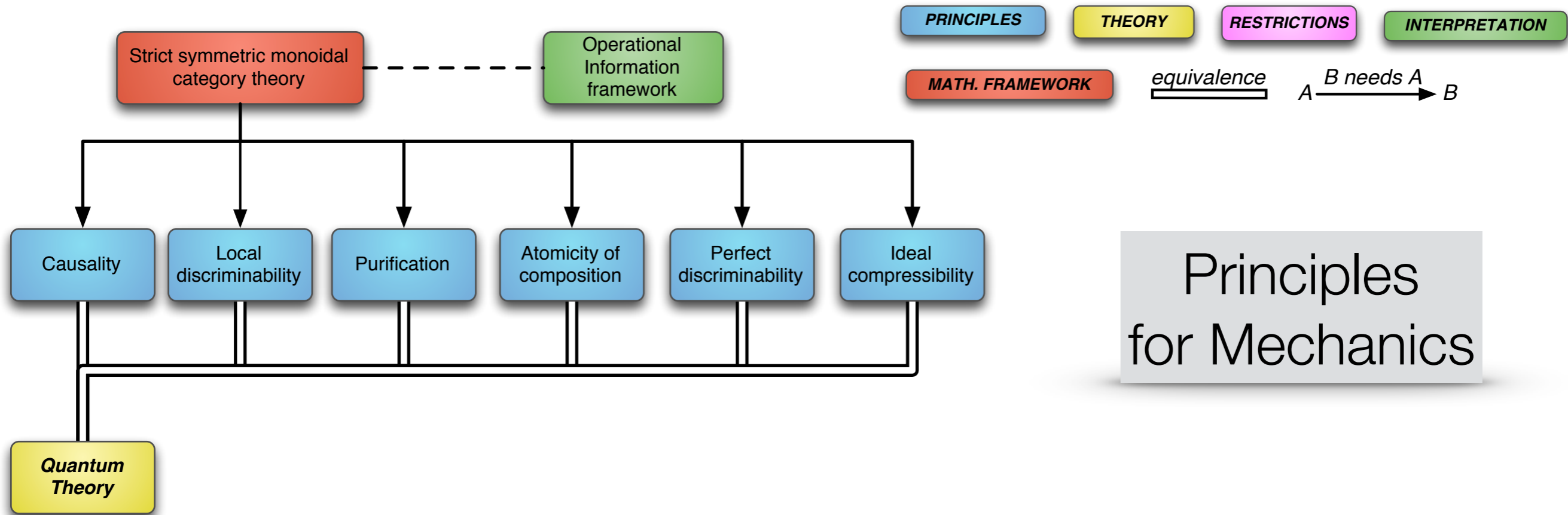
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- *Mechanics (QFT) derived in terms of countably many quantum systems in interaction*



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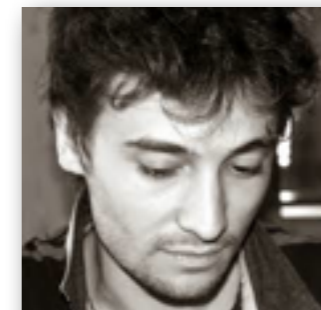
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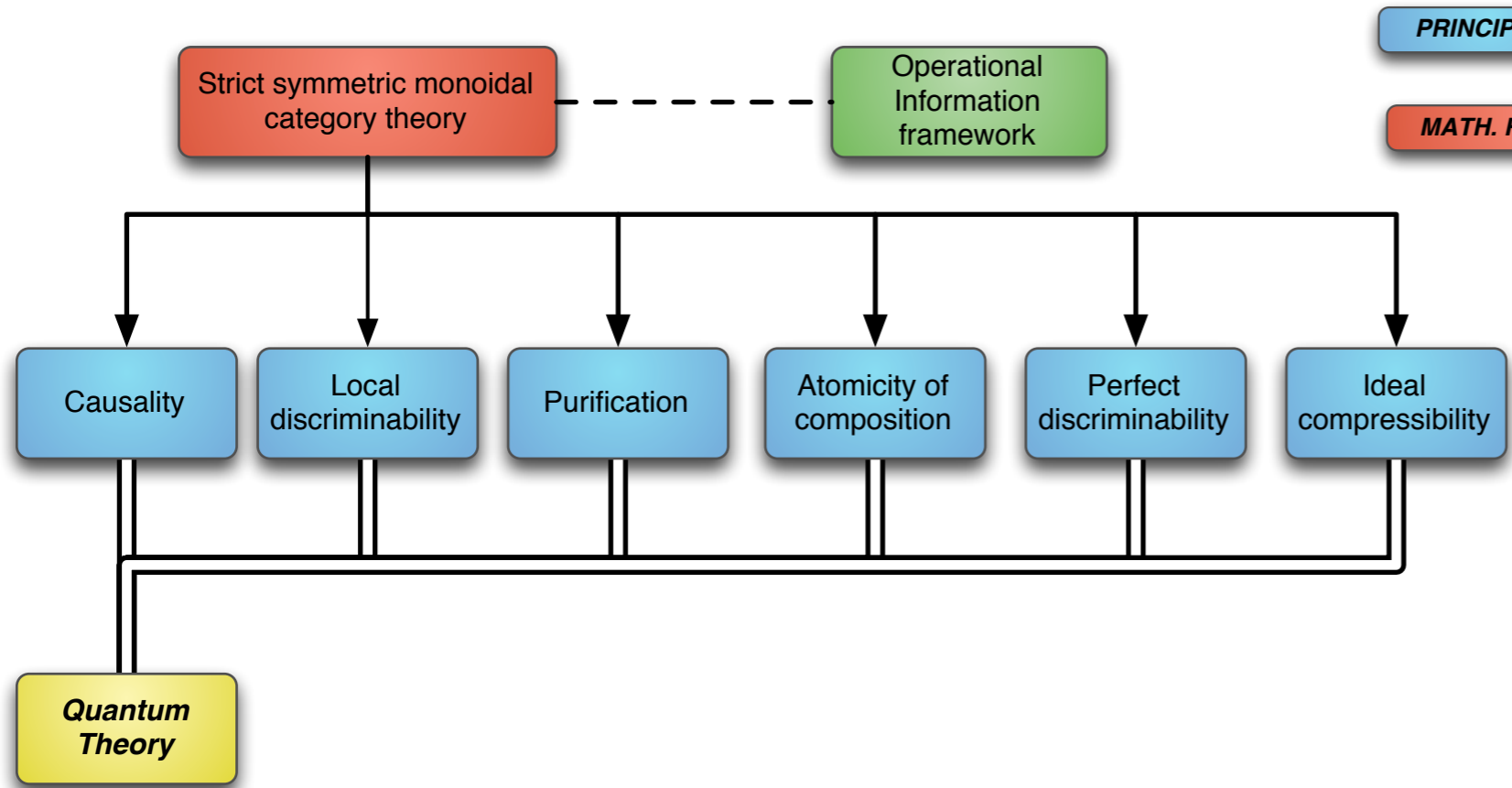
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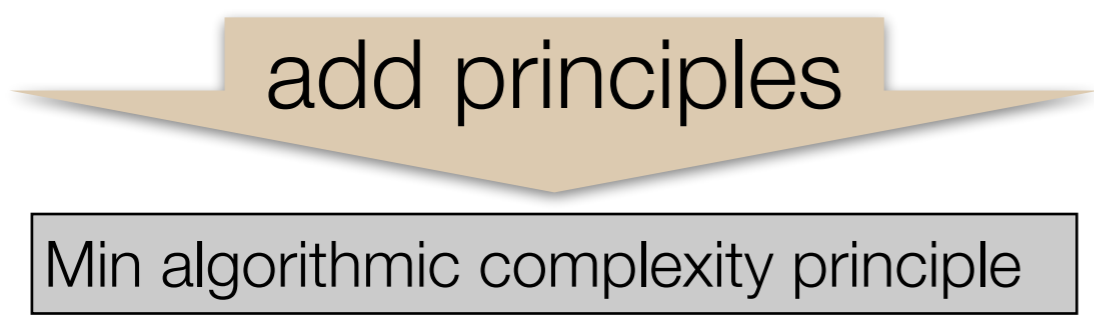


Principles for Mechanics

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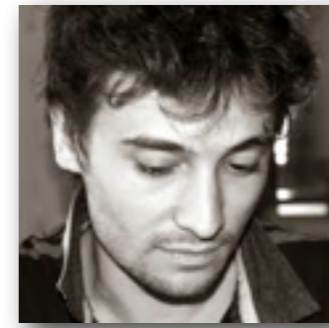
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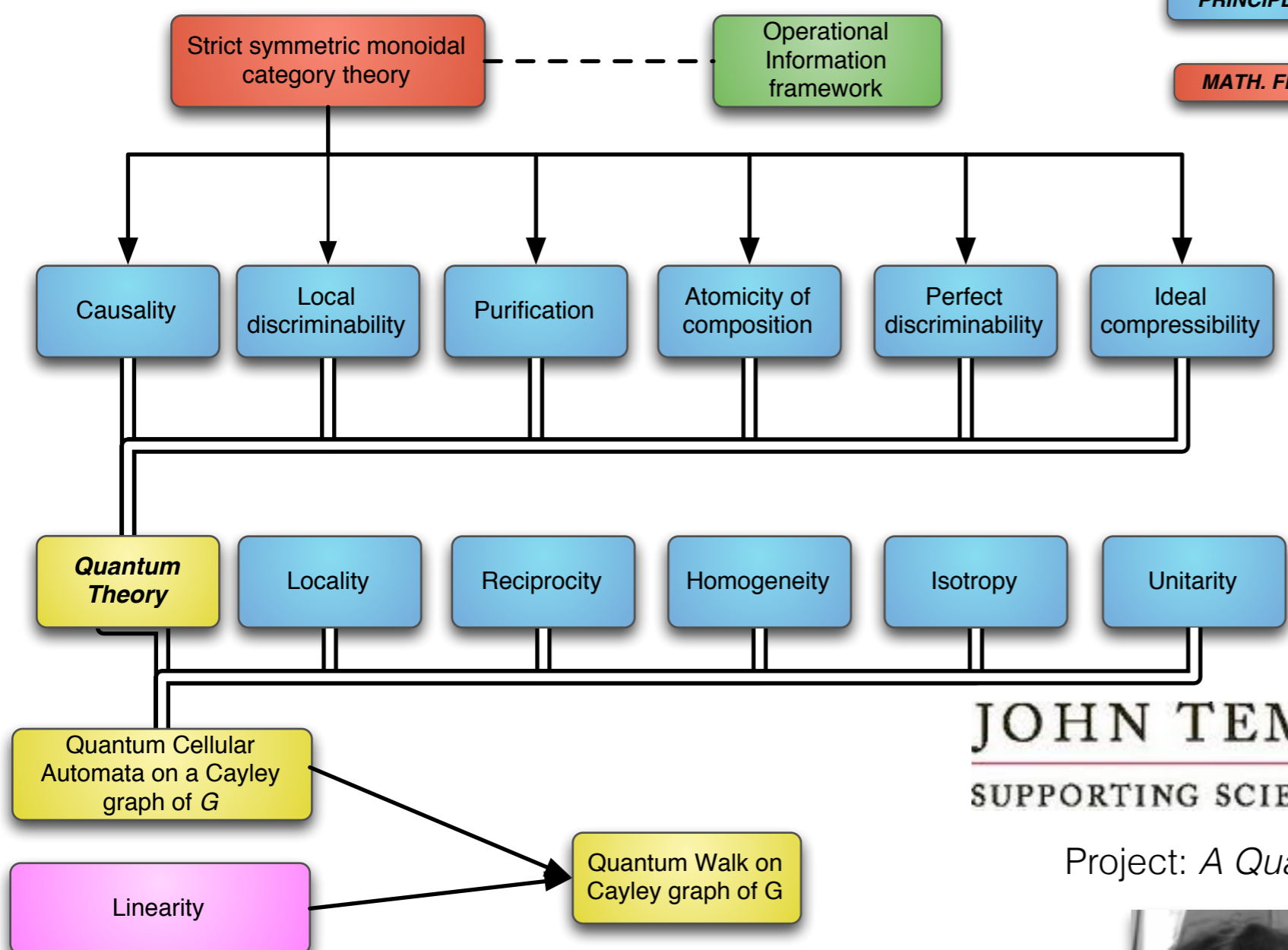
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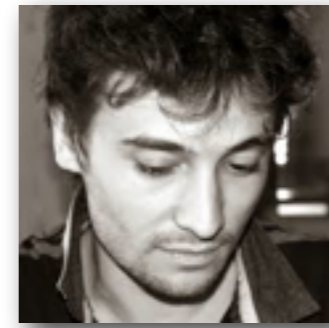
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- *Mechanics (QFT) derived in terms of countably many quantum systems in interaction*

add principles

Min algorithmic complexity principle

Quantum walk on Cayley graph

w.l.g. Hilbert space $\mathcal{H} = \bigoplus_{g \in G} \mathbb{C}^{s_g}$ $|G| \leq \aleph$, $s_g \in \mathbb{N}$

Evolution

$$\psi_g(t+1) = \sum_{g' \in S_g} A_{gg'} \psi_{g'}(t)$$

$$\sum_{g'} A_{gg'} A_{g''g'}^\dagger = \sum_{g'} A_{gg'}^\dagger A_{g''g'} = \delta_{gg''} I_{s_g}$$

- 1) Locality: S_g uniformly bounded
 2) Reciprocity: $A_{gg'} \neq 0 \implies A_{g'g} \neq 0$
 3) Homogeneity: all $g \in G$ are equivalent
- } = Quantum Walk on Cayley graph

Quantum walk on Cayley graph

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The following operator over the Hilbert space $\ell^2(G) \otimes \mathbb{C}^s$ is unitary

$$A = \sum_{h \in S} T_h \otimes A_h$$

where T is the right regular representation of G on $\ell^2(G)$ acting as

$$T_g |g'\rangle = |g'g^{-1}\rangle$$

- 1) Locality: S_g uniformly bounded
- 2) Reciprocity: $A_{gg'} \neq 0 \implies A_{g'g} \neq 0$
- 3) Homogeneity: all $g \in G$ are equivalent
- 4) Isotropy:

There exist:

- a group L of permutations of S_+ , transitive over S_+ that leaves the Cayley graph invariant
- a unitary s -dimensional (projective) representation $\{L_l\}$ of L such that:

$$A = \sum_{h \in S} T_h \otimes A_h = \sum_{h \in S} T_{lh} \otimes L_l A_h L_l^\dagger$$

PRINCIPLES

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RESTRICTIONS

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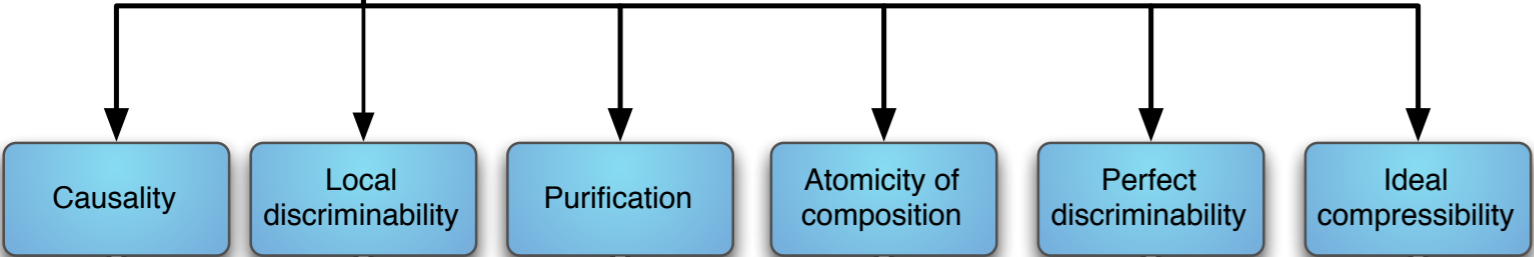
MATH. FRAMEWORK

equivalence

$A \xrightarrow{B \text{ needs } A} B$

Strict symmetric monoidal category theory

Operational Information framework



Quantum Cellular Automata on a Cayley graph of G

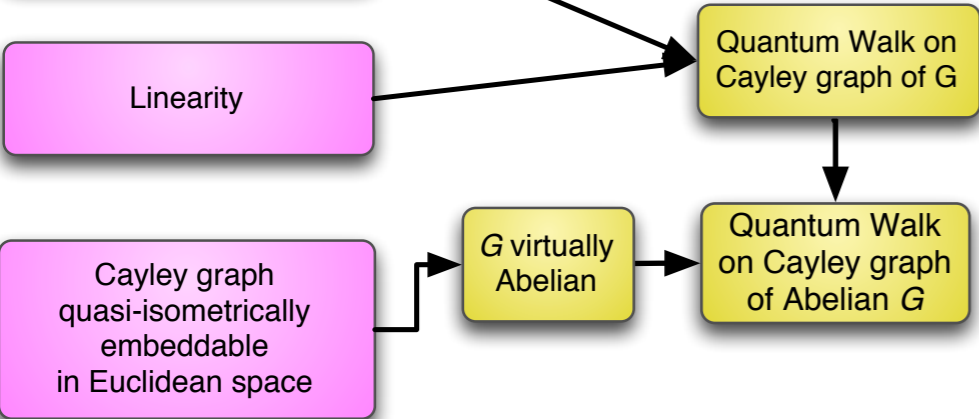
Linearity

Quantum Walk on Cayley graph of G

Cayley graph quasi-isometrically embeddable in Euclidean space

G virtually Abelian

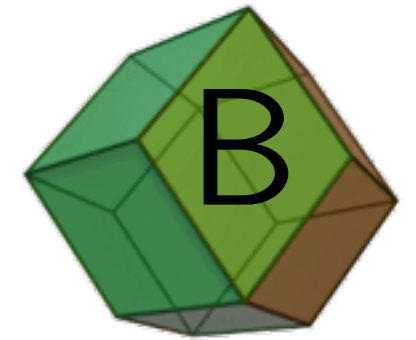
Quantum Walk on Cayley graph of Abelian G



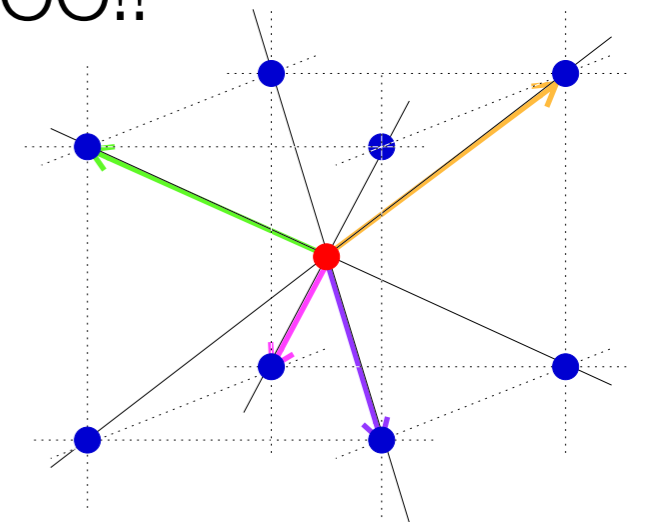
The Weyl QW

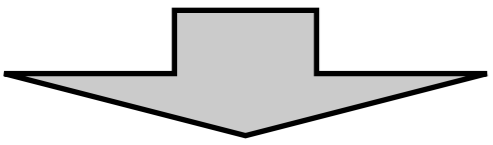
☞ Minimal dimension for nontrivial unitary A : $s=2$

Unitarity + isotropy \Rightarrow for $d=3$ the only Cayley is the BCC!!



Unitary operator:
$$A = \int_B^{\oplus} d\mathbf{k} A_{\mathbf{k}}$$




 Two QWs
 connected
 by P

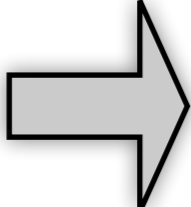

$$\begin{aligned}
 A_{\mathbf{k}}^{\pm} = & -i\sigma_x (s_x c_y c_z \pm c_x s_y s_z) \\
 & \mp i\sigma_y (c_x s_y c_z \mp s_x c_y s_z) \\
 & -i\sigma_z (c_x c_y s_z \pm s_x s_y c_z) \\
 & + I (c_x c_y c_z \mp s_x s_y s_z)
 \end{aligned}$$

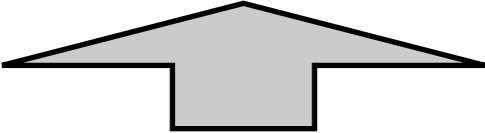
$$\begin{aligned}
 s_{\alpha} &= \sin \frac{k_{\alpha}}{\sqrt{3}} \\
 c_{\alpha} &= \cos \frac{k_{\alpha}}{\sqrt{3}}
 \end{aligned}$$

Physical interpretation: the Weyl Fermions

$$i\partial_t\psi(t) \simeq \frac{i}{2}[\psi(t+1) - \psi(t-1)] = \frac{i}{2}(A - A^\dagger)\psi(t)$$

$$\begin{aligned} \frac{i}{2}(A_{\mathbf{k}}^\pm - A_{\mathbf{k}}^{\pm\dagger}) = & + \sigma_x (s_x c_y c_z \pm c_x s_y s_z) \quad \text{“Hamiltonian”} \\ & \pm \sigma_y (c_x s_y c_z \mp s_x c_y s_z) \\ & + \sigma_z (c_x c_y s_z \pm s_x s_y c_z) \end{aligned}$$

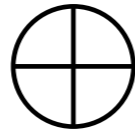
$k \ll 1$  $i\partial_t\psi = \frac{1}{\sqrt{3}}\boldsymbol{\sigma}^\pm \cdot \mathbf{k} \psi$  Weyl equation! $\boldsymbol{\sigma}^\pm := (\sigma_x, \pm\sigma_y, \sigma_z)$


Two QCAs
connected
by P

$$\begin{aligned} A_{\mathbf{k}}^\pm = & -i\sigma_x (s_x c_y c_z \pm c_x s_y s_z) \\ & \mp i\sigma_y (c_x s_y c_z \mp s_x c_y s_z) \\ & -i\sigma_z (c_x c_y s_z \pm s_x s_y c_z) \\ & + I (c_x c_y c_z \mp s_x s_y s_z) \end{aligned}$$

$$\begin{aligned} s_\alpha &= \sin \frac{k_\alpha}{\sqrt{3}} \\ c_\alpha &= \cos \frac{k_\alpha}{\sqrt{3}} \end{aligned}$$

Dirac QW



Local coupling: $A_{\mathbf{k}}$ coupled with its inverse with off-diagonal identity block matrix

$$E_{\mathbf{k}}^{\pm} = \begin{pmatrix} nA_{\mathbf{k}}^{\pm} & imI \\ imI & nA_{\mathbf{k}}^{\pm\dagger} \end{pmatrix}$$

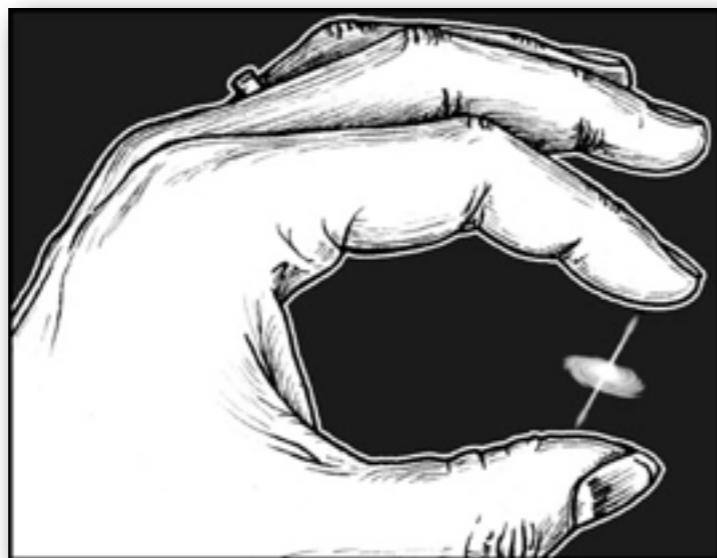
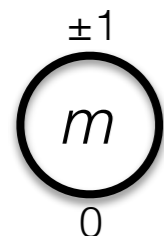
$$n^2 + m^2 = 1 \quad n, m \in \mathbb{R}$$

$E_{\mathbf{k}}^{\pm}$ CPT-connected!

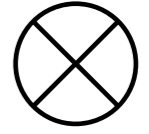
$$\omega_{\pm}^E(\mathbf{k}) = \cos^{-1} [n(c_x c_y c_z \mp s_x s_y s_z)]$$

Dirac in relativistic limit $k \ll m \ll 1$

m : mass, $m^2 \leq 1$
 n^{-1} : refraction index



Maxwell QW

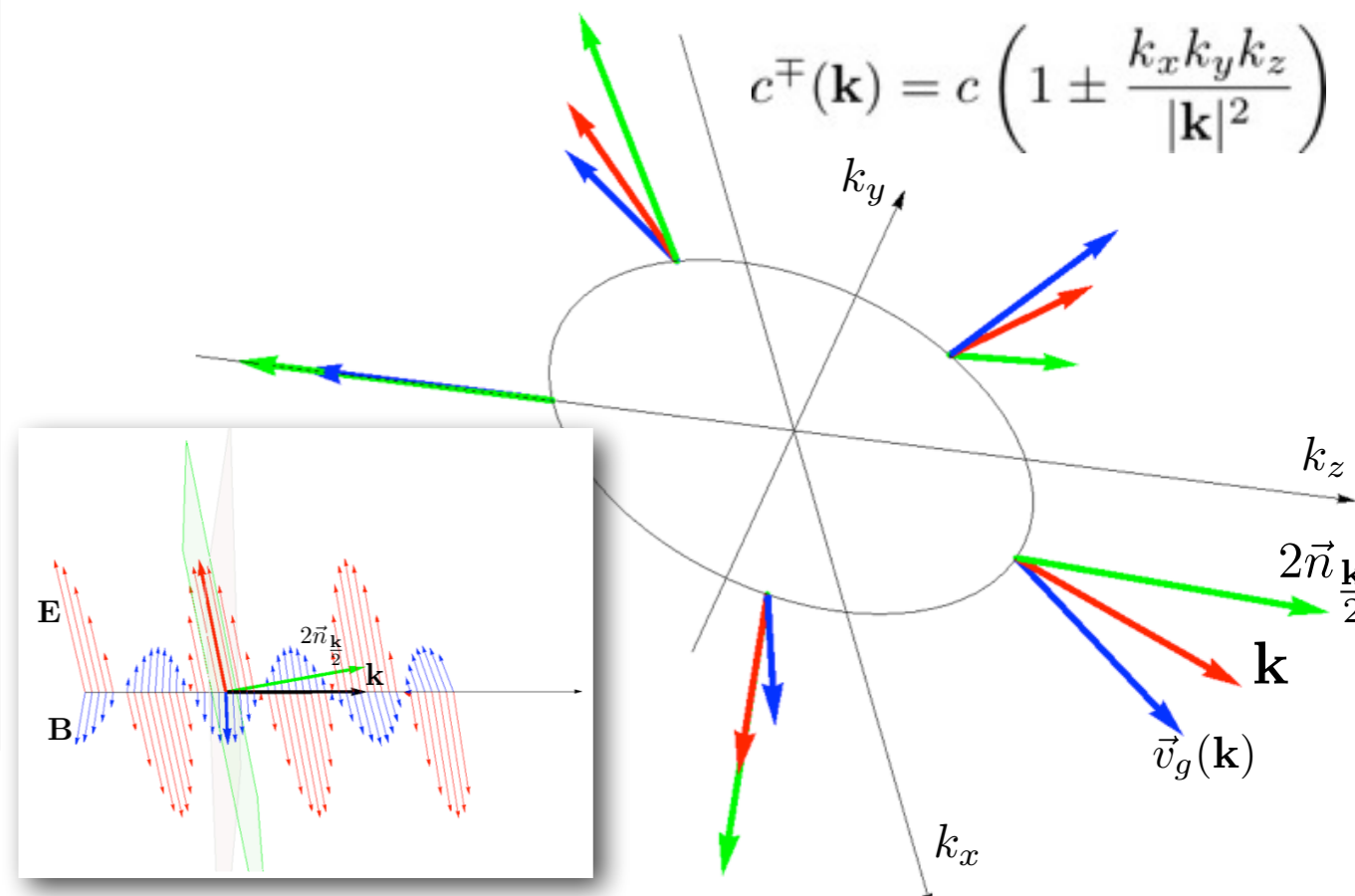


$$M_{\mathbf{k}}^{\pm} = A_{\mathbf{k}}^{\pm} \otimes A_{\mathbf{k}}^{\pm*}$$

$$F^{\mu}(\mathbf{k}) = \int \frac{d\mathbf{q}}{2\pi} f(\mathbf{q}) \tilde{\psi}(\frac{\mathbf{k}}{2} - \mathbf{q}) \sigma^{\mu} \varphi(\frac{\mathbf{k}}{2} + \mathbf{q})$$

Maxwell in relativistic limit $k \ll 1$

Boson: emergent from convolution of fermions
 (De Broglie neutrino-theory of photon)



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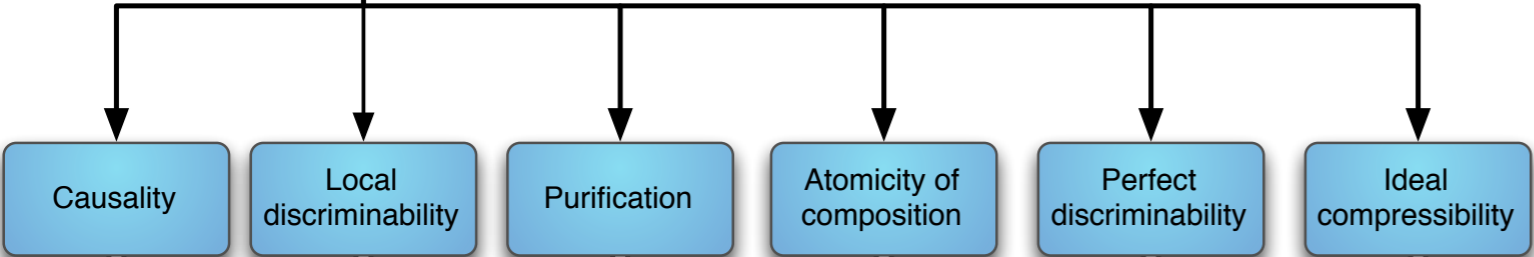
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Cayley graph quasi-isometrically embeddable in Euclidean space

G virtually Abelian

Quantum Walk on Cayley graph of Abelian G

Free Quantum Field Theory:
got it!

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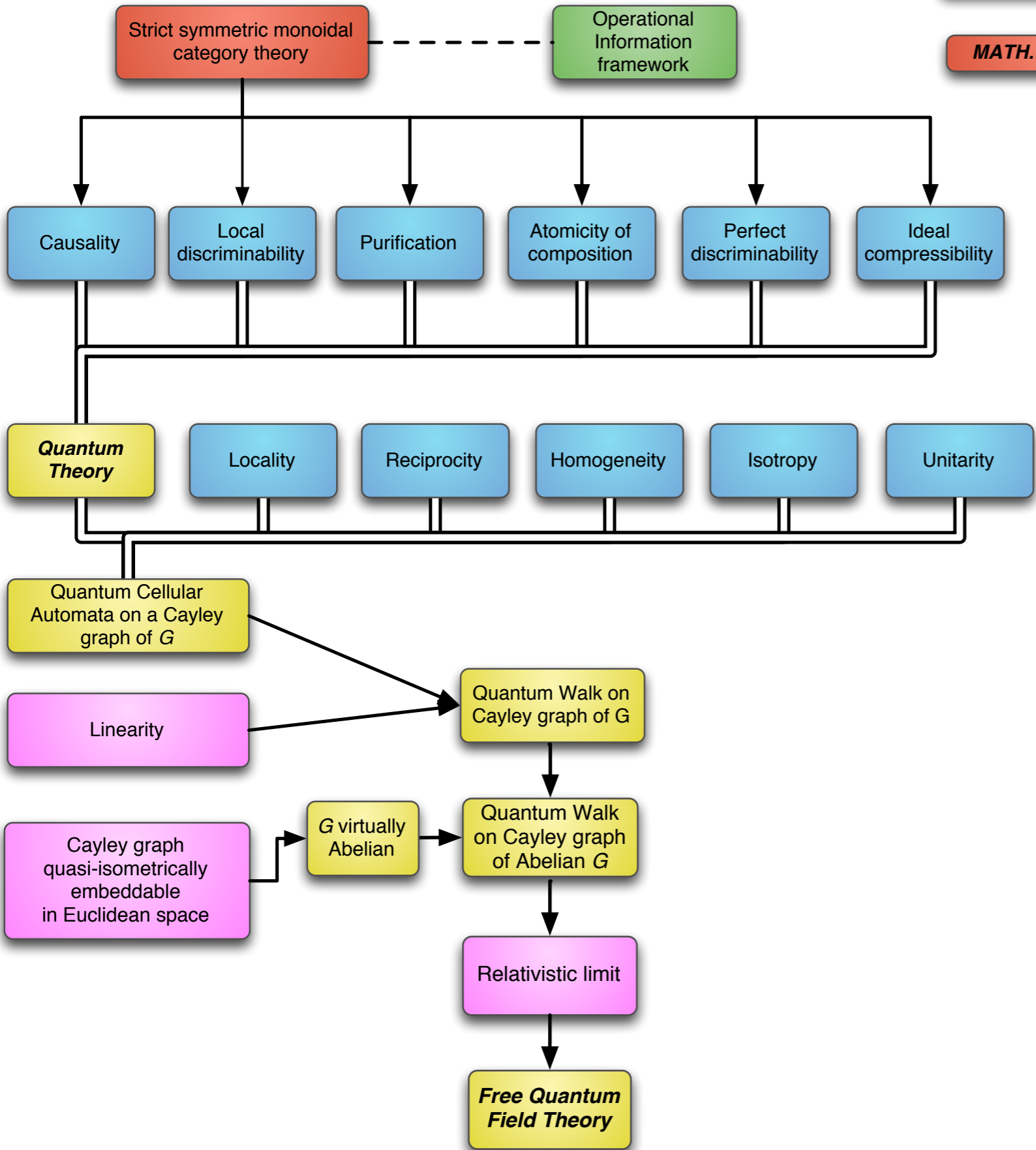
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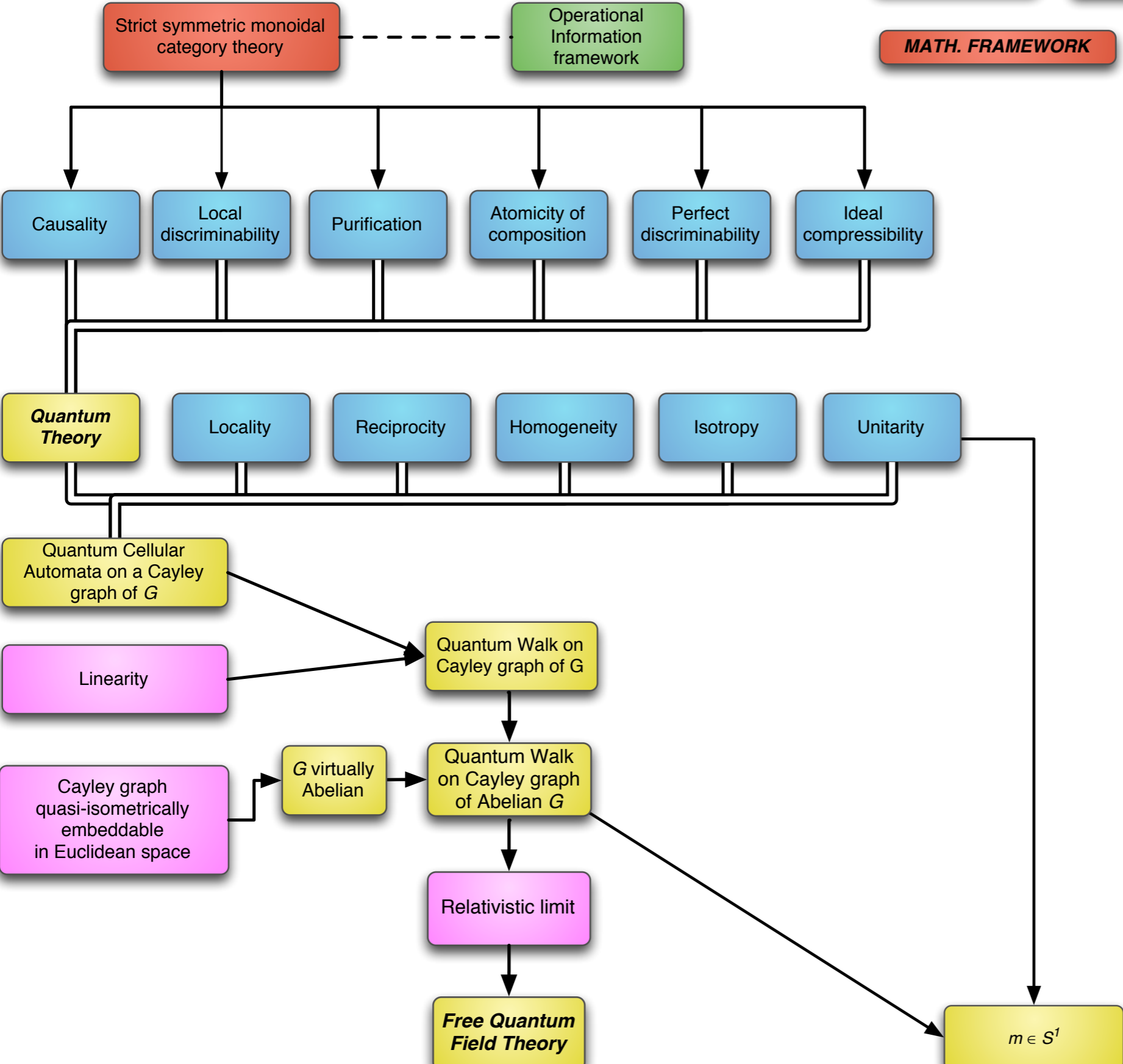
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Relativistic limit

Free Quantum Field Theory

$m \in S^1$



The theory contains its own LTM standards!

$$M \simeq \frac{1}{\sqrt{3}} \frac{\hbar k}{c(k) - c(0)}$$

$$\left\{ \begin{array}{l} c \equiv c(0) = \frac{a}{\tau} \\ \hbar = Mac \end{array} \right.$$

Case of study 1: Special Relativity recovered

- Mathematical statement:
invariance of eigenvalue equation under change of representation.
- Physical interpretation:
invariance of the physical law under change of inertial reference frame.

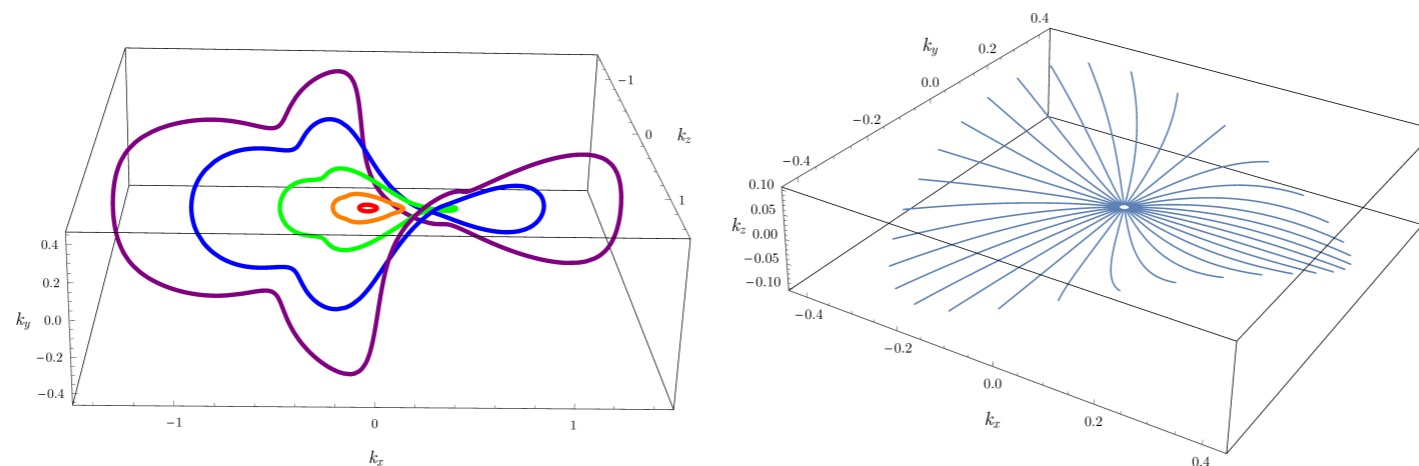


FIG. 2: The distortion effects of the Lorentz group for the discrete Planck-scale theory represented by the quantum walk in Eq. (6). Left figure: the orbit of the wavevectors $\mathbf{k} = (k_x, 0, 0)$, with $k_x \in \{.05, .2, .5, 1, 1.7\}$ under the rotation around the z axis. Right figure: the orbit of wavevectors with $|\mathbf{k}| = 0.01$ for various directions in the (k_x, k_y) plane under the boosts with β parallel to \mathbf{k} and $|\beta| \in [0, \tanh 4]$.

$m=0$

Deformed Poincaré group

- Lorentz transformations are perfectly recovered for $k \ll 1$.
- For $k \sim 1$:
 - *Double Special Relativity* (Camelia-Smolin).
 - *Relative locality* (in addition to relativity of simultaneity)

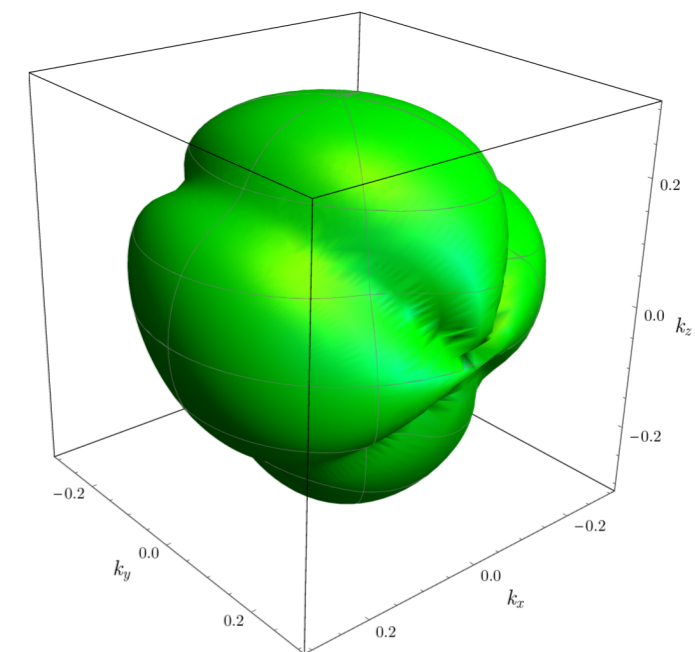
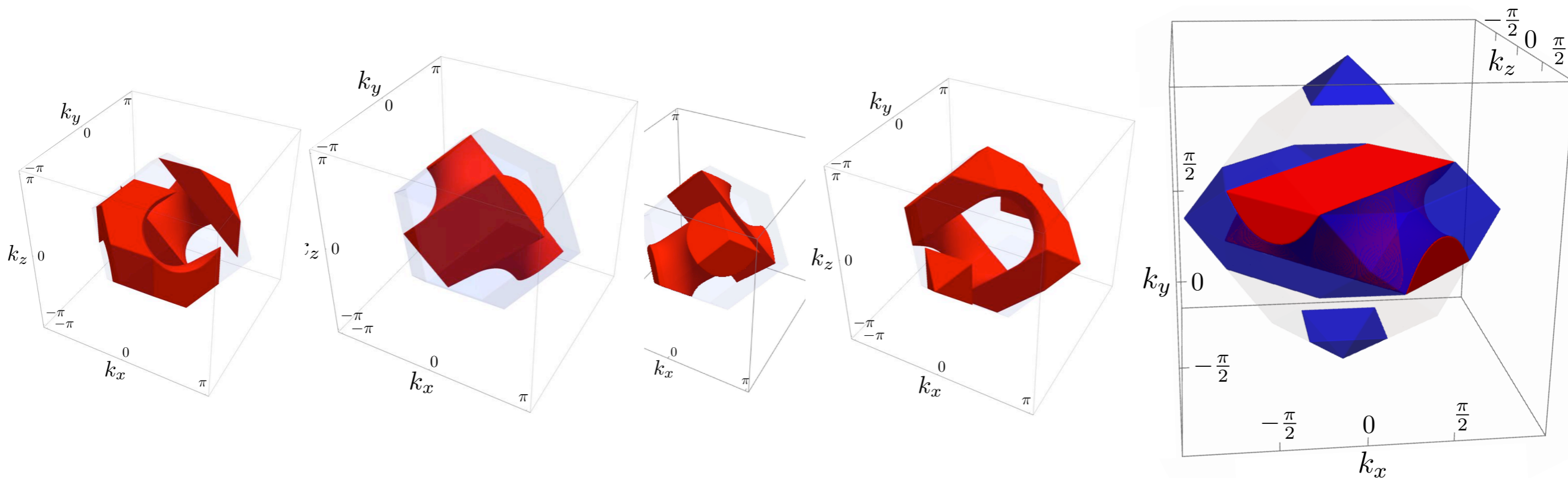


FIG. 3: The green surface represents the orbit of the wavevector $\mathbf{k} = (0.3, 0, 0)$ under the full rotation group $SO(3)$.

Case of study 2: particle notion

- Mathematical statement:
irreducible representation of deformed Poincaré group.
- Physical interpretation: particle!

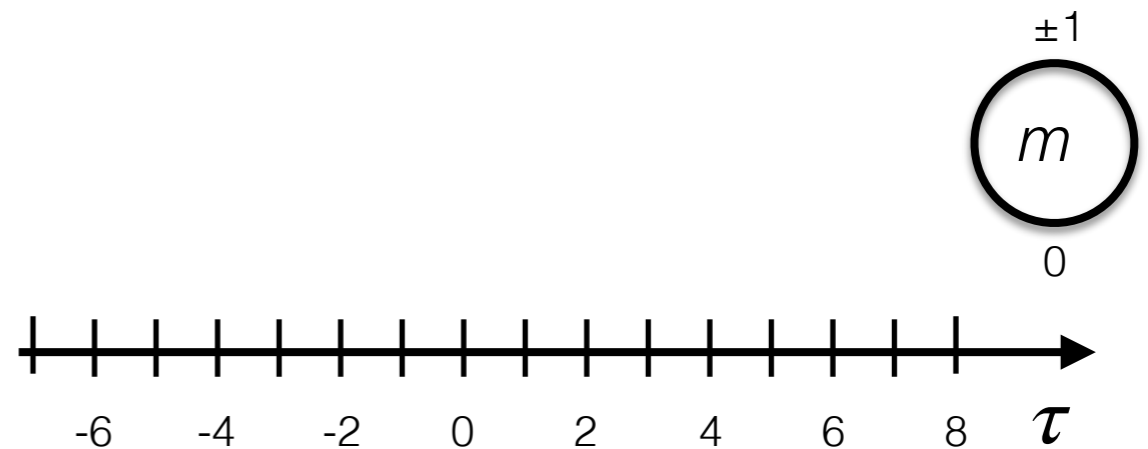


- The Brillouin zone separates into four Poincaré-invariant regions diffeomorphic to balls, corresponding to four different particles.
- $m \neq 0$ De Sitter $SO(1,4)$

Case of study 3: proper time

- Mathematical statement:
topology of the particle mass domain
- Physical interpretation: proper time is discrete!

$$H(q_\alpha, p_\alpha, \tau, m) = \sum_{\alpha} p_\alpha \dot{q}_\alpha + c^2 m \dot{\tau} - L$$



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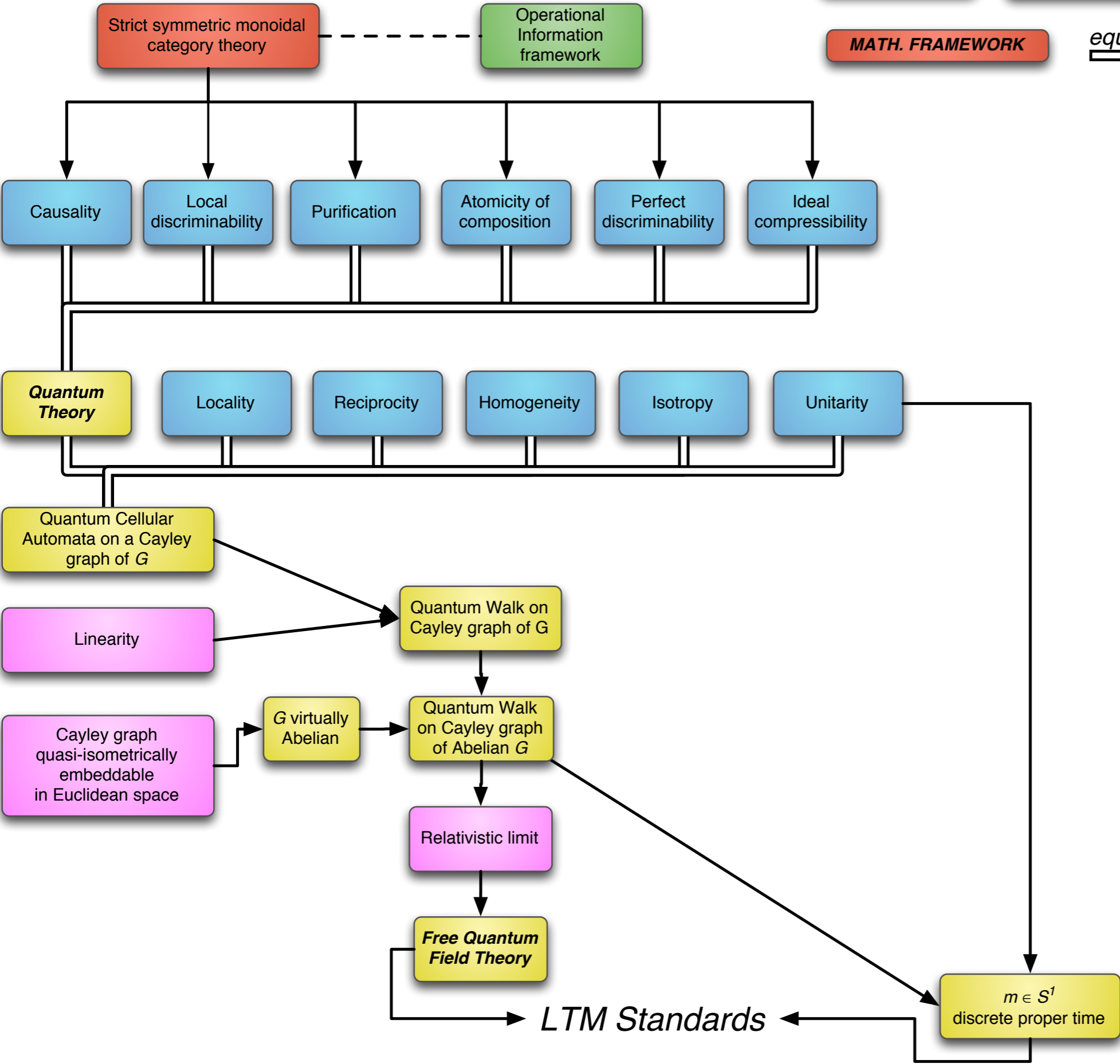
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Relativistic limit

Free Quantum Field Theory

LTM Standards

$m \in S^1$
discrete proper time



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Relativity Principle without space-time

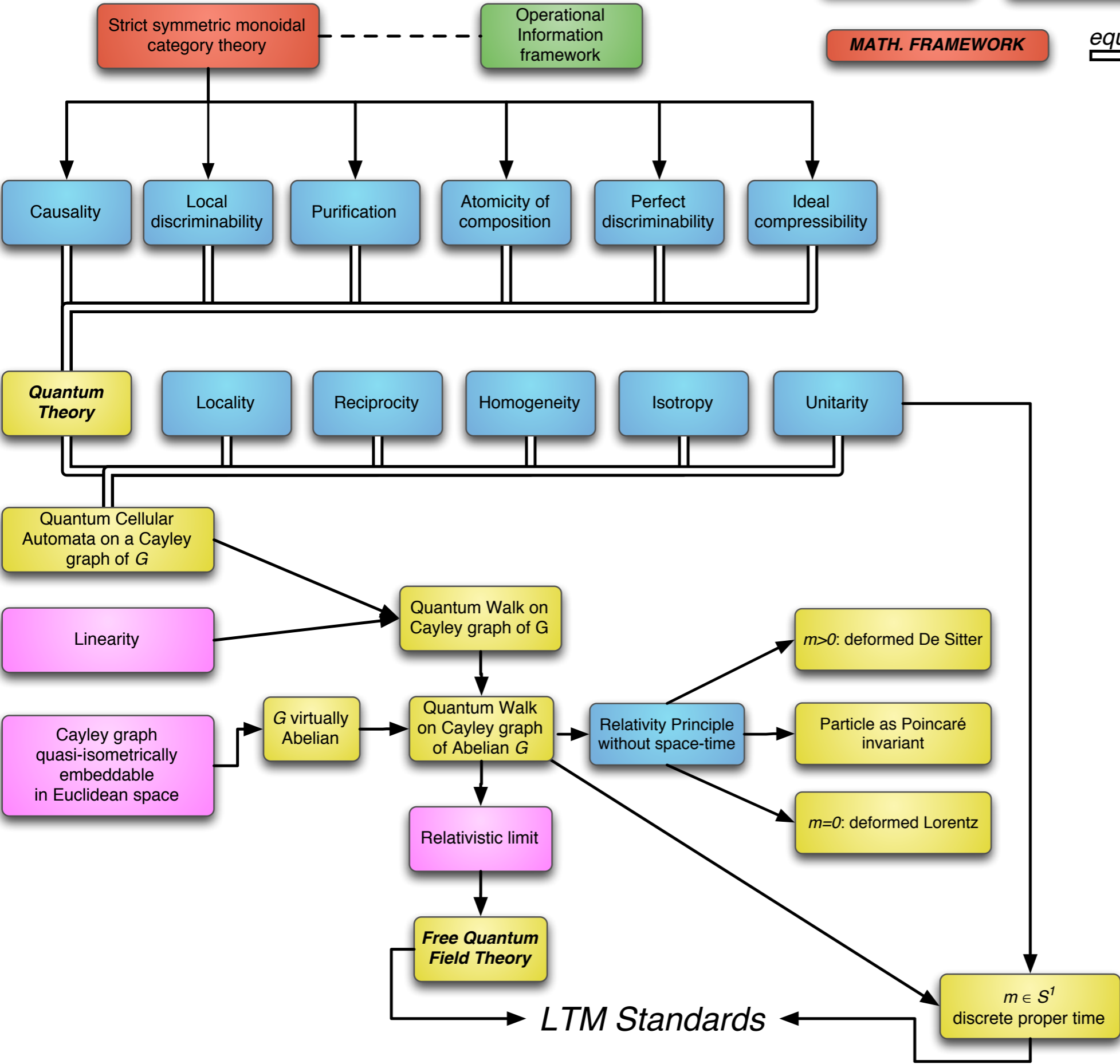
$m > 0$: deformed De Sitter

Particle as Poincaré invariant

$m = 0$: deformed Lorentz

$m \in S^1$ discrete proper time

LTM Standards



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Relativity Principle without space-time

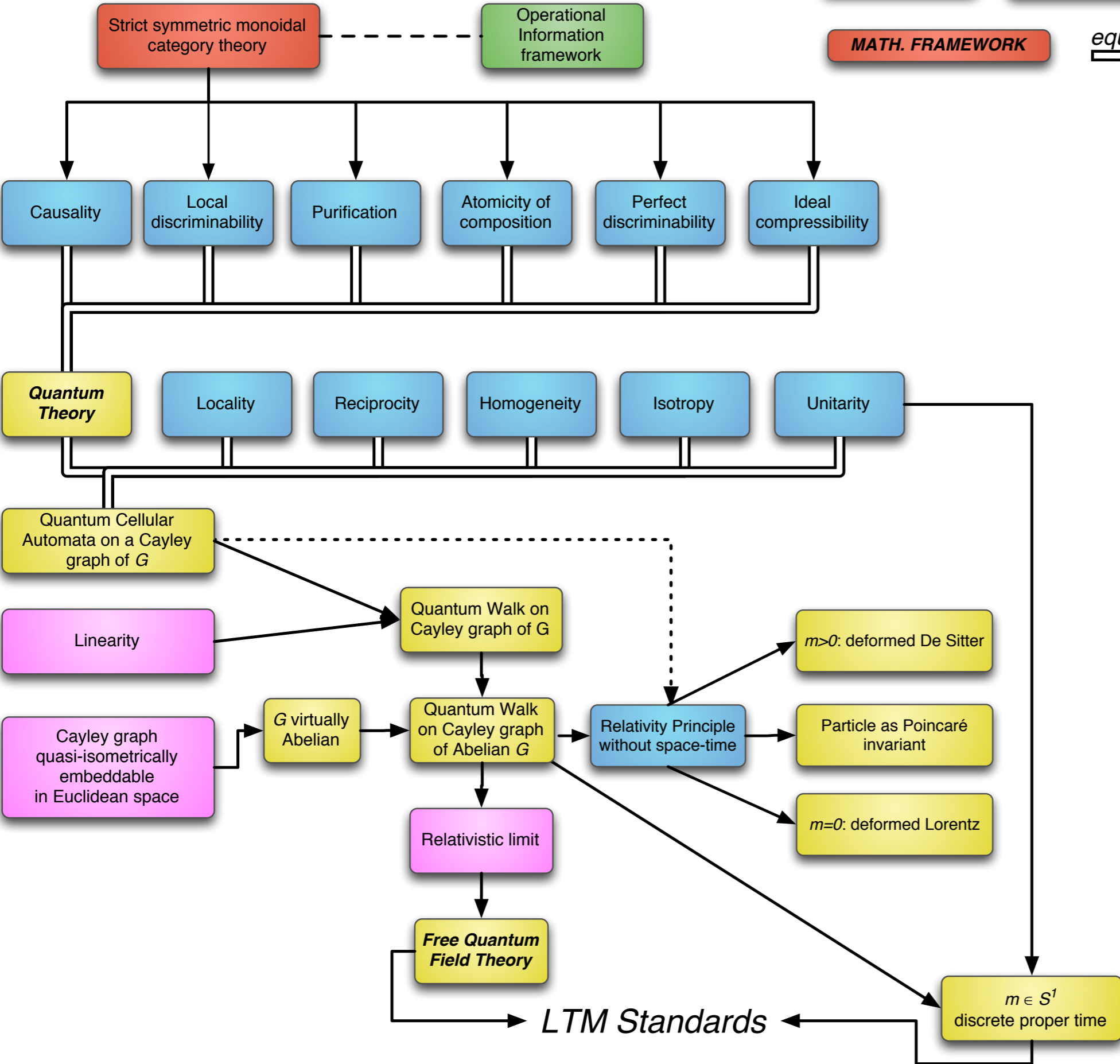
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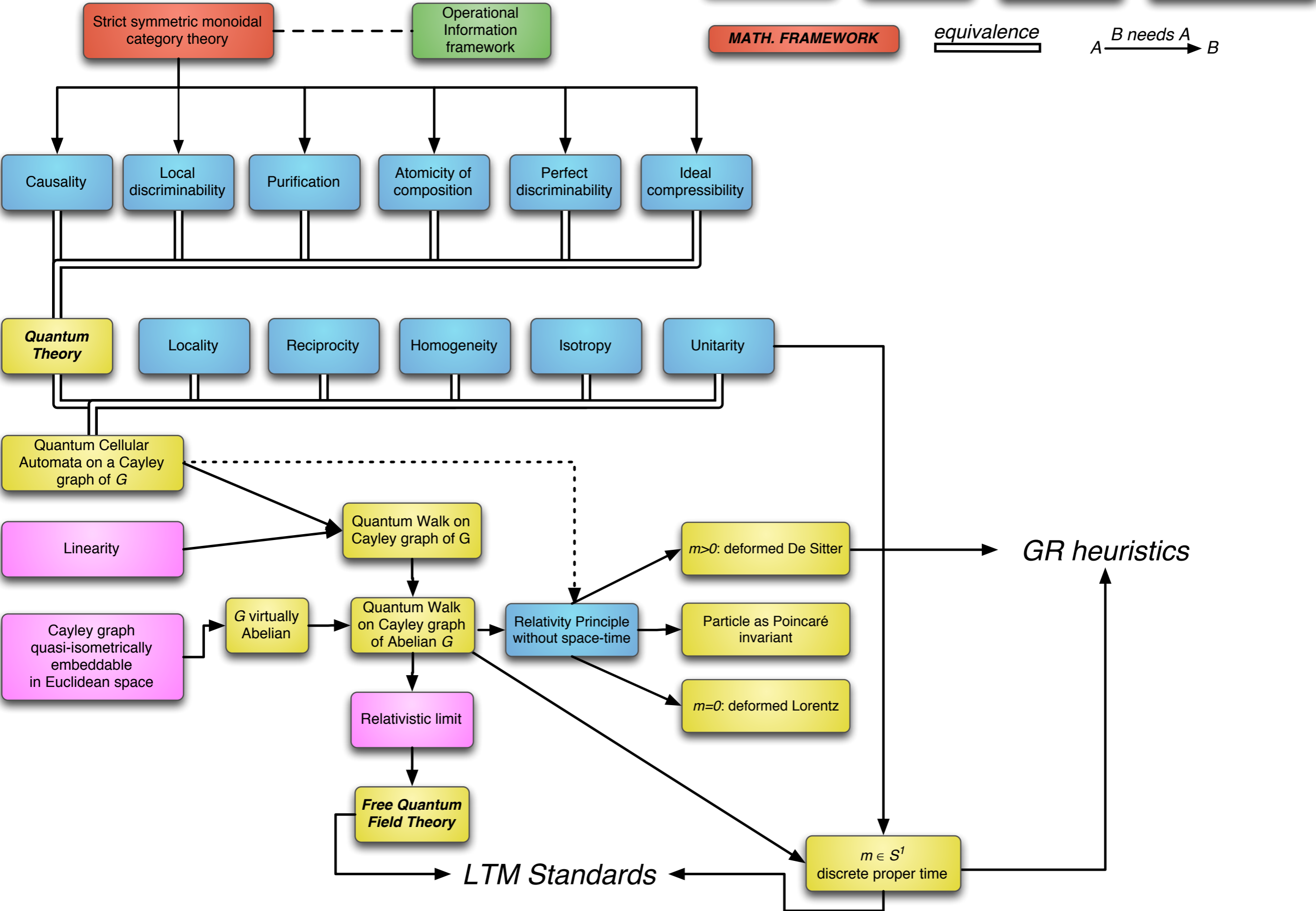
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GR heuristics

$m \in S^1$
discrete proper time

LTM Standards



A priori principles?

Conventionalism:

- Homogeneity
- Isotropy

Theory
simplicity

Homogeneity = clocks, roads ...

Symmetries

Minimization of
algorithmic complexity



Hans Reichenbach



Adolf Grünbaum

This is more or less what I wanted to say

Thank you for your attention